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DERIVATION OF THE STOCHASTIC BURGERS EQUATION WITH DIRICHLET BOUNDARY CONDITIONS FROM THE WASEP

PATRICIA GONÇALVES, NICOLAS PERKOWSKI, AND MARIELLE SIMON

ABSTRACT. We consider the weakly asymmetric simple exclusion process on the discrete space $\{1, \dots, n-1\}$ ($n \in \mathbb{N}$), in contact with stochastic reservoirs, both with density $\rho \in (0, 1)$ at the extremity points, and starting from the invariant state, namely the Bernoulli product measure of parameter ρ . Under time diffusive scaling tn^2 and for $\rho = \frac{1}{2}$, when the asymmetry parameter is taken of order $1/\sqrt{n}$, we prove that the density fluctuations at stationarity are macroscopically governed by the energy solution of the stochastic Burgers equation with Dirichlet boundary conditions, which is shown to be unique and to exhibit different boundary behavior than the Cole-Hopf solution.

1. INTRODUCTION

1.1. Context. A vast amount of physical phenomena were first described at the macroscopic scale, in terms of the classical partial differential equations (PDEs) of mathematical physics. Over the last decades the scientific community has tried to give a precise understanding of their derivation from first principles at the microscopic level in order to identify the limits of their validity. Typically, the microscopic systems are composed of a huge number of atoms and one looks at a very large time scale with respect to the typical frequency of atom vibrations. Mathematically, this corresponds to a space-time scaling limit procedure.

The macroscopic laws that can arise from microscopic systems can either be partial differential equations (PDEs) or stochastic PDEs (SPDEs) depending on whether one is looking at the convergence to the mean or at the fluctuations around that mean. Among the classical SPDEs is the Kardar-Parisi-Zhang (KPZ) equation which has been first introduced more than thirty years ago in [30] as the *universal* law describing the fluctuations of randomly growing interfaces of one-dimensional stochastic dynamics close to a stationary state (as for example, models of bacterial growth, or fire propagation). Since then, it has generated an intense research activity among the physics and mathematics community. In particular, the *weak KPZ universality conjecture* [4, 39, 40] states that the fluctuations of a large class of one-dimensional microscopic interface growth models are ruled at the macroscopic scale by solutions of the *KPZ equation*, which reads as follows: consider a time variable t and a one-dimensional space variable u , then the evolution of the height $\mathcal{Z}(t, u)$ of the randomly growing interface can be described by

$$d\mathcal{Z}(t, u) = A\Delta\mathcal{Z}(t, u) dt + B(\nabla\mathcal{Z}(t, u))^2 dt + \sqrt{C}d\mathcal{W}_t, \quad (1.1)$$

where A, B, C are thermodynamic constants depending on the model and $\partial_t \mathcal{W}$ is a space-time white noise. Note that the non-linearity $(\nabla\mathcal{Z}(t, u))^2$ makes the KPZ equation (1.1) ill-posed, essentially because the trajectory of the solution lacks space regularity (due to the presence of the white noise), and therefore the square of its distributional derivative is not defined.

One possible way to solve this equation is to consider its *Cole-Hopf transformation* Φ , which solves a stochastic heat equation with a multiplicative noise (SHE) and is related to \mathcal{Z} through a log-transformation. Since the SHE is linear, its solution can easily be constructed and it is unique, and the solution to the KPZ equation can then simply be defined as the inverse Cole-Hopf transformation of Φ . However, the solution to the SHE is too irregular to allow for a change of variable formula, and a priori there is no meaningful equation associated to its inverse Cole-Hopf transformation. Only recently Hairer has developed a meaningful notion of solution for the KPZ equation and proved existence and uniqueness of such solutions with periodic boundary conditions, see [27]. His approach uses rough path integrals to construct the nonlinearity $(\nabla \mathcal{Z}(t, u))^2$, and it has inspired the development of new technologies (regularity structures [28] and paracontrolled distributions [22]) for the so-called *singular* stochastic partial differential equations (SPDEs).

The first breakthrough towards the weak KPZ universality conjecture is due to Bertini and Giacomin: in their seminal paper [1], they show that the Cole-Hopf solution can be obtained as a scaling limit of the *weakly asymmetric exclusion process* (WASEP) (which will be defined ahead). Their approach consists in performing the Cole-Hopf transformation at the microscopic level, following [14], and then showing that this microscopic Cole-Hopf transformation solves a linear equation (similarly to what happens at the macroscopic level). Since then, this strategy has been used in more sophisticated models, see [5, 6, 7, 8], however the applicability of the microscopic Cole-Hopf transformation is limited to a very specific class of particle systems.

Another way to look at the KPZ equation is via the *stochastic Burgers equation* (SBE), which is obtained from (1.1) by taking its derivative: if $\mathcal{Y}_t = \nabla \mathcal{Z}_t$, then \mathcal{Y}_t satisfies

$$d\mathcal{Y}(t, u) = A\Delta \mathcal{Y}(t, u) dt + B\nabla(\mathcal{Y}^2(t, u)) dt + \sqrt{C}d\nabla \mathcal{W}_t, \quad (1.2)$$

which has of course the same regularity issues as the KPZ equation. Nevertheless, this formulation is well adapted to derive KPZ behavior from microscopic models. Indeed, the work initiated by Gonçalves and Jara in [17] has introduced a new tool, called *second order Boltzmann-Gibbs principle* (BGP), which makes the non-linear term $\nabla(\mathcal{Y}^2(t, u))$ of the SBE naturally emerge from the underlying microscopic dynamics. The authors have first proved the BGP for general weakly asymmetric simple exclusion processes, and shortly thereafter it has been extended to a wider class of microscopic systems, such as zero-range models [19], integrable Hamiltonian one-dimensional chains [20], non-degenerate kinetically constrained models [2], exclusion processes with defects [10], non-simple exclusion processes [18], semilinear SPDEs [24], or interacting diffusions [9]. From the BGP, it comes naturally that some suitably rescaled microscopic quantity, called *density fluctuation field* (see below for a precise meaning) subsequentially converges, as the size of the microscopic system goes to infinity, to random fields which are solutions \mathcal{Y} of a generalized martingale problem for (1.2), where the singular non-linear drift $\nabla(\mathcal{Y}^2(t, u))$ is a well-defined space-time distributional random field. Gonçalves and Jara in [17] (see also [23]) called them *energy solutions*. Recently, Gubinelli and Perkowski [25] proved uniqueness of energy solutions to (1.2) and as a significant consequence, the proof of the weak KPZ universality conjecture could be concluded for all the models mentioned above. We note that the energy solutions, compared to the methods of [1] or [28, 22], are strongly based on stationarity, and in particular the weak universality was so far only shown for stationary initial conditions or bounded entropy perturbations thereof. On the other hand, none of the models mentioned above admit a microscopic Cole-Hopf transformation,

which prevents the use of the methods of [1], and in many cases they do not have the structure of a semilinear SPDE, which means that the pathwise approach of [28, 22] does not apply either.

1.2. Purposes of this work. Our goal in this article is to go beyond the weak KPZ universality conjecture and to derive a new SPDE, namely, the KPZ equation with boundary conditions, from an interacting particle system in contact with stochastic reservoirs. Indeed, the presence of boundary conditions in evolution equations often lacks understanding from a physical and mathematical point of view. Here we intend to legitimate the choice done at the macroscopic level for the KPZ/SBE equation from the microscopic description of the system. For that purpose, we first prove two main theorems:

- (Theorem 3.3) We extend the notion of energy solutions to the stochastic Burgers equation (1.2) by adding Dirichlet boundary conditions: we set up a rigorous definition and prove existence and uniqueness of such solutions.
- (Theorem 3.17) We construct a microscopic model (inspired by the WASEP mentioned above) from which the energy solution naturally emerges as the macroscopic limit of its stationary density fluctuations.

This gives a physical justification for the Dirichlet boundary conditions in (1.2). We also introduce the notion of energy solutions to two related SPDEs: the KPZ equation with Neumann boundary conditions and the SHE with Robin boundary conditions; we prove their existence and uniqueness, and we rigorously establish the formal links between the equations discussed above. This is more subtle than expected, because the boundary conditions do not behave canonically: a formal computation suggests that the Cole-Hopf transform of the KPZ equation with Neumann boundary conditions should also satisfy Neumann boundary conditions, but we show that instead it satisfies Robin boundary conditions.

We also associate an interface growth model to our microscopic model, roughly speaking by integrating it in the space variable, and show that it converges to the energy solution of the KPZ equation, thereby giving a physical justification of the Neumann boundary conditions. In the remaining lines we go into further details and explain these results.

As for the KPZ/Burgers equation on the real line, our equation can be (and has been) solved also with the Cole-Hopf transform [5, 37] or regularity structures [15]. However, just as on the real line there are many models with boundary conditions to which our approach applies and that are not accessible with either of these methods, for example kinetically constrained exclusion processes, exclusion processes with (finite variance) long jumps, or the interacting diffusions of [9]. Moreover, since the solution cannot be evaluated pointwise the implementation of the Dirichlet boundary conditions is actually rather subtle, and it is one of the contributions of this work to sort out the links between the different formulations (which in general correspond to different processes).

1.2.1. WASEP with reservoirs. At the microscopic level, the model that we consider is the following: let a discrete system of particles evolve on the finite set $\Lambda_n = \{1, \dots, n-1\}$ of size $n \in \mathbb{N}$. For any site $x \in \Lambda_n$ there is at most one particle, and we denote by $\eta(x) \in \{0, 1\}$ its occupation variable, namely: $\eta(x) = 1$ if there is a particle at site x , and 0 otherwise. We then define a continuous-time Markov process $\eta_t = \{\eta_t(x) ; x \in \Lambda_n\}$ on the state space $\{0, 1\}^{\Lambda_n}$ using the following possible moves: for any $x \notin \{1, n-1\}$, a particle at site x can attempt to jump to its neighbouring sites $x-1$ or $x+1$, provided that they are empty. Similarly, a particle at site 1 can jump to its right neighbour 2 or it

can leave the system, and the particle at site $n - 1$ can jump to its left neighbour $n - 2$ or it can leave the system. Moreover, attached to the extremities of Λ_n there are two reservoirs of particles: one at site 0, which can send a particle to site 1 (if this site is empty), and the other one at site n , which can send a particle to site $n - 1$ (if this site is empty). Another interpretation of the boundary dynamics could be given as follows: particles can either be created (resp. annihilated) at the sites $x = 1$ or $x = n - 1$ if the site is empty (resp. occupied).

All possible moves are endowed with random Poissonian clocks, independently of each other. Along the time evolution, we launch the jump whose clock rings first, and after the jump all clocks are reset. Namely, we are given a family of independent Poisson processes, indexed by all the possible moves $x \curvearrowright y$, with intensities $\lambda(x, y)$, on the time line $[0, \infty)$. The intensities $\lambda(x, y)$ depend on the occupation variables $\eta(x), \eta(y)$ and on some small parameter $\varepsilon_n > 0$ as follows:

- if $x \in \{2, \dots, n - 2\}$, then we assume

$$\lambda(x, x + 1) = (1 + \varepsilon_n)\eta(x)(1 - \eta(x + 1)), \quad (1.3)$$

$$\lambda(x, x - 1) = \eta(x)(1 - \eta(x - 1)), \quad (1.4)$$

- while at the boundaries, for some fixed $\rho \in (0, 1)$ we set

$$\lambda(0, 1) = \rho(1 + \varepsilon_n)(1 - \eta(1)), \quad \lambda(n - 1, n) = (1 - \rho)(1 + \varepsilon_n)\eta(n - 1), \quad (1.5)$$

$$\lambda(1, 0) = (1 - \rho)\eta(1), \quad \lambda(n, n - 1) = \rho(1 - \eta(n - 1)). \quad (1.6)$$

In (1.3) and (1.5), the factor $1 + \varepsilon_n$ breaks the symmetry of the jumps: there is a non-trivial drift towards the right. But since $\varepsilon_n \rightarrow 0$, we are in the weakly asymmetric setting. Moreover, note that the product $\eta(x)(1 - \eta(x + 1))$ in (1.3) (for instance) corresponds to the *exclusion rule* explained above: for the jump $x \curvearrowright x + 1$ to have a non-zero intensity, there has to be a particle at site x , and no particle at site $x + 1$.

An invariant measure for these dynamics is the Bernoulli product measure on $\{0, 1\}^{\Lambda_n}$ of parameter ρ , which we denote by ν_ρ^n , and whose marginals satisfy:

$$\nu_\rho^n\{\eta(x) = 1\} = 1 - \nu_\rho^n\{\eta(x) = 0\} = \rho.$$

We start the Markov process η_t under the invariant measure ν_ρ^n and we look at the evolution on the diffusive time scale tn^2 , where $t \geq 0$ is the macroscopic time. The *microscopic density fluctuation field* is then defined as

$$\mathscr{Y}_t^n(\cdot) = \frac{1}{\sqrt{n}} \sum_{x=1}^{n-1} (\eta_{tn^2}(x) - \rho) \delta_{x/n}(\cdot), \quad (1.7)$$

where $\delta_{x/n}$ is the Delta dirac distribution, and therefore (\cdot) is meant to be a test function. We prove in Theorem 3.17 two main results on the large n limit of that field, assuming that the initial density is $\rho = \frac{1}{2}$ (this assumption, which aims at removing a transport phenomenon inside the system, will be explained in detail in Remark 3.21):

- if $\varepsilon_n = E/\sqrt{n}$ (for some $E > 0$), the sequence of processes \mathscr{Y}^n converges, in a suitable space of test functions, towards the unique energy solution \mathscr{Y} to (1.2) with Dirichlet boundary conditions, with parameters $A = 1$, $B = E$ and $C = \frac{1}{2}$;
- whenever $\sqrt{n}\varepsilon_n \rightarrow 0$ as $n \rightarrow \infty$, the sequence \mathscr{Y}^n converges towards the unique Ornstein-Uhlenbeck (OU) process with Dirichlet boundary conditions, which is equivalent to the unique energy solution to (1.2) with parameters $A = 1$, $B = 0$ and $C = \frac{1}{2}$.

As we will discuss in Remark 3.19 below it is easy to generalize these convergence results to initial conditions that are bounded entropy perturbations of the invariant measure.

1.2.2. The second order Boltzmann-Gibbs principle. The main ingredient that we use at the microscopic level is the BGP that we have to reprove completely in our particular setting. Indeed, this is the first time that the BGP is established in a space with boundaries.

This tool, which was first introduced in [17], permits to investigate the space-time fluctuations of the microscopic density fluctuation field, and to prove that, when properly recentered, the latter is close in the macroscopic limit to a quadratic functional of the density field (1.7). It was originally proved for general weakly asymmetric simple exclusion processes by a multi-scale argument (also given in [16]), which consists in replacing, at each step, a local function of the dynamics by another function whose support increases and whose variance decreases, and its proof used a key spectral gap inequality which is uniform in the density of particles, and is not available for many models. Later in [19] it is assumed that the models satisfy a spectral gap bound which does not need to be uniform in the density of particles and allows more general jump rates. More recently in [10, 20], and then in [2], a new proof of the BGP has permitted to extend the previous results to models which do not need to satisfy a spectral gap bound, as, for example, symmetric exclusion processes with a slow bond, and microscopic dynamics that allow degenerate exchange rates. In this paper, we adapt that strategy, which turns out to be quite robust, to our finite model with stochastic boundary reservoirs. This is the goal of Theorem 4.2.

1.2.3. Boundary behavior at the microscopic scale. As we already mentioned, the KPZ equation and SBE are closely related, and this relation can be seen also at the microscopic level. There is a natural *height function* $h(x)$ associated to the occupation variable $\eta(x)$, and in particular its increments satisfy:

$$h(x+1) - h(x) = \eta(x) - \rho.$$

Of course, this only determines h up to the addition of a constant, and we will carefully define a microscopic height process $\{h_t(x) ; t \geq 0, x = 1, \dots, n\}$, which satisfies $h_t(x+1) - h_t(x) = \eta_t(x) - \rho$ for the Markov process η_t from above, and which has nice path properties. We refer the reader to Section 2.4.2 for the rigorous definition. Similarly to (1.7), we are interested in the macroscopic limit of height fluctuations starting the evolution from ν_ρ^n and looking in the time scale tn^2 . In this case the averaged local height at site $x \in \{1, \dots, n\}$ and microscopic time tn^2 is equal to $c_n t$, where c_n is related to the initial density ρ and the strength of the asymmetry ε_n , as follows: $c_n = n^2 \varepsilon_n \rho (1 - \rho)$. Therefore, the (renormalized) *microscopic height fluctuation field* is defined as

$$\mathcal{X}_t^n(\cdot) = \frac{1}{n^{3/2}} \sum_{x=1}^n (h_{tn^2}(x) - c_n t) \delta_{x/n}(\cdot), \quad (1.8)$$

which means, formally, that $\mathcal{Y}_t^n = -\nabla \mathcal{X}_t^n$. In the same spirit as Theorem 3.17, for $\varepsilon_n = E/\sqrt{n}$ and $\rho = \frac{1}{2}$, we prove:

(Theorem 3.20) The sequence of processes \mathcal{X}^n converges, in a suitable space of distributions, towards the unique energy solution \mathcal{X} to (1.1) with Neumann boundary conditions, with the same parameters $A = 1$, $B = E$ and $C = \frac{1}{2}$.

A closely related convergence result was recently shown by Corwin and Shen [5], who consider the height function associated to a variant of the WASEP in contact with reservoirs, apply Gärtner's microscopic Cole-Hopf transform, and show that in the scaling limit the process converges to a solution of the SHE with Robin boundary conditions. Corwin and Shen study a range of parameters and do not need to know the invariant measure of their system explicitly (in general it will not be a Bernoulli product measure). However, the model we study here has parameters that are not admissible in [5], because in their formulation the parameters A and B have to be positive (see e.g. Proposition 2.7 and p.14 in [5]), while we would get $A = B = -E^2/8$ (see Proposition 3.13 below and also Remark 3.11). The extension of the Cole-Hopf approach to general parameters A, B and thus also to the model that we consider here was completed only a couple of weeks after this paper by Parekh [37]. We stress however that our methods are very different from the ones used in [5, 37] and we study the microscopic density fluctuation field without relying on the microscopic Cole-Hopf transform.

1.2.4. Uniqueness of energy solutions and boundary behavior at the macroscopic scale. The convergence proofs described above show relative compactness of the microscopic density fluctuations field (resp. the microscopic height fluctuation field) under rescaling, and that any limiting point is an energy solution to the stochastic Burgers equation with Dirichlet boundary conditions (resp. the KPZ equation with Neumann boundary conditions). To conclude the convergence, it remains to prove the uniqueness of energy solutions. We achieve this following the same strategy used in [25] for proving the uniqueness of energy solutions on the real line: We mollify an energy solution \mathcal{Y} to the SBE, \mathcal{Y}^n , find a suitable anti-derivative \mathcal{Z}^n of \mathcal{Y}^n , and let $\Phi^n = e^{\mathcal{Z}^n}$. Now we take the mollification away and show that $\Phi = \lim_n \Phi^n$ solves (a version of) the SHE. Since uniqueness holds for solutions to the SHE, $\mathcal{Y} = \nabla \log \Phi$ must be unique. But while the strategy is the same as in [25], the technical details are considerably more involved because our space is no longer translation invariant and many of the tools of [25] break down. Moreover, the dynamics of Φ^n contain a singular term that converges to Dirac deltas at the boundaries, a new effect which can be interpreted as a change of boundary conditions: Formally we would expect that $\nabla \Phi_t(0) = e^{\mathcal{Z}_t(0)} \nabla \mathcal{Z}_t(0) = 0$ because $\nabla \mathcal{Z}_t(0) = 0$ by the Neumann boundary conditions for \mathcal{Z} . However, the singular term in the dynamics changes the boundary conditions to Robin's type, $\nabla \Phi_t(0) = -D\Phi_t(0)$, $\nabla \Phi_t(1) = D\Phi_t(1)$ for a constant $D \in \mathbb{R}$ which depends on the parameters A, B, C .

In that sense \mathcal{Z} can be interpreted as a Cole-Hopf solution to the KPZ equation with inhomogeneous Neumann boundary conditions, and then the question arises which of the two formulations (inhomogeneous or homogeneous Neumann conditions) accurately describes the behavior of our stochastic process at the boundary. While of course the main difficulty with the KPZ equation is that its solutions are non-differentiable and in particular we cannot evaluate $\nabla \mathcal{Z}_t(0)$ pointwise, we show in Proposition 3.14 that after averaging a bit in time there are canonical interpretations for $\nabla \mathcal{Z}_t(0)$ and $\nabla \mathcal{Z}_t(1)$, and both indeed vanish. We also show in Proposition 3.15 that the formal change of the boundary conditions from Neumann to Robin in the exponential transformation $\Phi = e^{\mathcal{Z}}$ is reflected in the “pointwise” boundary behavior of Φ (again after averaging in time).

This should be compared with the recent work of Gerencsér and Hairer [15] who show that the classical solution \mathcal{Z}^ε to the KPZ equation,

$$d\mathcal{Z}_t^\varepsilon = \Delta \mathcal{Z}_t^\varepsilon dt + ((\nabla \mathcal{Z}_t^\varepsilon)^2 - c^\varepsilon) dt + d\mathcal{W}_t^\varepsilon, \quad (1.9)$$

with Neumann boundary condition $\nabla \mathcal{Z}_t^\varepsilon(0) = \nabla \mathcal{Z}_t^\varepsilon(1) = 0$ and where \mathcal{W}^ε is a mollification of \mathcal{W} and $\{c^\varepsilon\}_{\varepsilon>0}$ is a sequence of diverging constants, may converge to different limits satisfying different boundary conditions as $\varepsilon \rightarrow 0$, depending on which mollifier was used for \mathcal{W}^ε . But if the noise is only mollified in space and white in time, then the limit is always the same and it agrees with the Cole-Hopf solution with formal Neumann boundary condition $\nabla \mathcal{Z}_t(0) = \nabla \mathcal{Z}_t(1) = 0$. Since our results show that the Cole-Hopf solution with formal boundary condition $\nabla \mathcal{Z}_t(0) = -\nabla \mathcal{Z}_t(1) = -E^2/4$ satisfies $\nabla \mathcal{Z}_t(0) = \nabla \mathcal{Z}_t(1) = 0$ in the “physical” sense, this suggests that in order to obtain the correct limit it is not only necessary to subtract a large diverging constant c^ε in (1.9), but additionally one should perform a boundary renormalization. Indeed, under boundary conditions the solution \mathcal{Z}^ε to (1.9) is not spatially homogeneous, and therefore there is no reason to renormalize it by subtracting a constant c^ε and instead the renormalization might also be spatially inhomogeneous.

1.3. Outline of the paper: In Section 2 we give the precise definition of our microscopic dynamics and its invariant measures, we also introduce all the spaces of test functions where the microscopic fluctuation fields, namely the density and the height, will be defined and we give the proper definition of these fields. Section 3 contains all the rigorous definitions of solutions to the SPDEs that we obtain, namely the OU/SBE with Dirichlet boundary conditions, the KPZ equation with Neumann boundary conditions and the SBE with linear Robin boundary conditions, and we explain how these equations are linked. In Section 4 we prove the convergence of the microscopic fields to the solutions of the respective SPDEs, namely Theorems 3.17 and 3.20. In Section 5 we prove the second order Boltzmann-Gibbs principle, which is the main technical result that we need at the microscopic level in order to recognize the macroscopic limit of the density fluctuation field as an energy solution to the SBE. Finally, in Section 6 we give the proof of the uniqueness of solutions to the aforementioned SPDEs. The appendices contain some important aside results that are needed along the paper, but to facilitate the reading flow we removed them from the main body of the text. In particular, Appendix E sketches how one could prove that the microscopic Cole-Hopf transformation of the microscopic density fluctuation field converges to the SHE, and in particular we show that already at the microscopic level the Cole-Hopf transformation changes the boundary conditions from Neumann to Robin’s type.

2. MODEL AND DEFINITIONS

2.1. The microscopic dynamics: WASEP in contact with stochastic reservoirs. For n a positive integer, we define $\Lambda_n = \{1, \dots, n-1\}$ and $\Omega_n = \{0, 1\}^{\Lambda_n}$ as the state space of a Markov process $\{\eta_t^n; t \geq 0\}$, whose dynamics is entirely encoded into its infinitesimal generator, denoted below by \mathcal{L}_n . More precisely, our process belongs to the family of well-known *weakly asymmetric simple exclusion processes*. Here we consider that the strength of asymmetry is ruled by a parameter $\gamma \geq \frac{1}{2}$, and we put the system of particles in contact with two stochastic reservoirs, whose rates of injection/removal of particles from the bulk depend on a parameter $\rho \in (0, 1)$, which is fixed. To keep notation simple we omit the dependence on ρ in all the quantities that we define ahead and in order to facilitate future computations, we write the generator as

$$\mathcal{L}_n = \mathcal{L}_n^{\text{bulk}} + \mathcal{L}_n^{\text{bnd}},$$

where $\mathcal{L}_n^{\text{bulk}}$ and $\mathcal{L}_n^{\text{bnd}}$ given below encode respectively the dynamics on the bulk, and on the left/right boundaries. For any $E > 0$ and $\gamma \geq \frac{1}{2}$ we define $\mathcal{L}_n^{\text{bulk}}$ and $\mathcal{L}_n^{\text{bnd}}$ acting

on functions $f : \Omega_n \rightarrow \mathbb{R}$ as follows: first,

$$(\mathcal{L}_n^{\text{bulk}} f)(\eta) = \sum_{x=1}^{n-2} \left\{ \left(1 + \frac{E}{n^\gamma} \right) \eta(x)(1 - \eta(x+1)) + \eta(x+1)(1 - \eta(x)) \right\} (\nabla_{x,x+1} f)(\eta),$$

where $(\nabla_{x,x+1} f)(\eta) = f(\sigma^{x,x+1} \eta) - f(\eta)$ and $\sigma^{x,x+1} \eta$ is the configuration obtained from η after exchanging $\eta(x)$ with $\eta(x+1)$, namely

$$(\sigma^{x,x+1} \eta)(y) = \begin{cases} \eta(x), & \text{if } y = x+1, \\ \eta(x+1), & \text{if } y = x, \\ \eta(y), & \text{if } y \notin \{x, x+1\}. \end{cases}$$

Second, the generator of the dynamics at the boundaries is given by

$$\begin{aligned} (\mathcal{L}_n^{\text{bnd}} f)(\eta) &= \left\{ \left(1 + \frac{E}{n^\gamma} \right) \rho(1 - \eta(1)) + \eta(1)(1 - \rho) \right\} (f(\sigma^1 \eta) - f(\eta)) \\ &\quad + \left\{ \left(1 + \frac{E}{n^\gamma} \right) \eta(n-1)(1 - \rho) + \rho(1 - \eta(n-1)) \right\} (f(\sigma^{n-1} \eta) - f(\eta)) \end{aligned}$$

where

$$(\sigma^x \eta)(y) = \begin{cases} 1 - \eta(y), & \text{if } y = x, \\ \eta(y), & \text{if } y \neq x. \end{cases}$$

From now on to simplify notation we denote the rates that appear above by $r_{x,x+1}(\eta)$: precisely, at any $x \in \{0, \dots, n-1\}$, and assuming the convention $\eta(0) = \eta(n) = \rho$, they are given by

$$r_{x,x+1}(\eta) = \left(1 + \frac{E}{n^\gamma} \right) \eta(x)(1 - \eta(x+1)) + \eta(x+1)(1 - \eta(x)). \quad (2.1)$$

We refer to Figure 1 for an illustration of the dynamics.

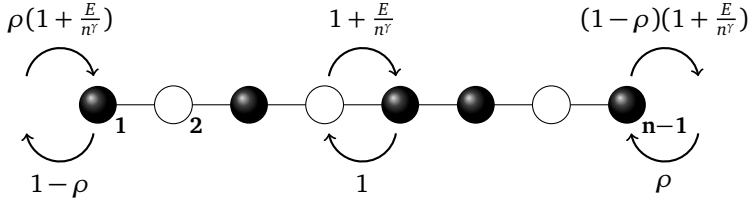


FIGURE 1. Illustration of the jump rates. The leftmost and rightmost rates are the entrance/exiting rates.

Throughout the paper, a time horizon line $T > 0$ is fixed. We are interested in the evolution of this exclusion process in the diffusive time scale tn^2 , thus we denote by $\{\eta_{tn^2}^n; t \in [0, T]\}$ the Markov process on Ω_n associated to the accelerated generator $n^2 \mathcal{L}_n$. The path space of càdlàg trajectories with values in Ω_n is denoted by $\mathcal{D}([0, T], \Omega_n)$. For any initial probability measure μ on Ω_n , we denote by \mathbb{P}_μ the probability measure on $\mathcal{D}([0, T], \Omega_n)$ induced by μ and the Markov process $\{\eta_{tn^2}^n; t \in [0, T]\}$. The corresponding expectation is denoted by \mathbb{E}_μ .

Notations: Throughout the paper, for any measurable space (U, ν) we denote by $L^2(\nu)$ or $L^2(U)$ the usual L^2 -space with norm $\|\cdot\|_{L^2(\nu)}$ and scalar product $\langle \cdot, \cdot \rangle_{L^2(\nu)}$. Whenever the integration variable u may be not clear to the reader, we enlighten it by denoting $L_u^2(\nu)$ or $L_u^2(U)$. We also write $f \lesssim g$ or $g \gtrsim f$ if there exists a constant $C > 0$,

independent of the parameters involved in f and g , such that $f \leq Cg$. If $f \lesssim g$ and $f \gtrsim g$, we also write $f \simeq g$. Finally we denote by $\mathbb{N} = \{1, 2, \dots\}$ the set of positive integers and by $\mathbb{N}_0 = \{0\} \cup \mathbb{N}$ the set of non-negative integers.

2.2. Invariant measures and Dirichlet form. For $\lambda \in (0, 1)$, let ν_λ^n be the *Bernoulli product measure* on Ω_n with density λ : its marginals satisfy

$$\nu_\lambda^n\{\eta : \eta(x) = 1\} = 1 - \nu_\lambda^n\{\eta : \eta(x) = 0\} = \lambda, \quad \text{for each } x \in \Lambda_n.$$

When $\lambda = \rho$, the measure ν_ρ^n is invariant but not reversible with respect to the evolution of $\{\eta_t^n; t \geq 0\}$. To ensure the invariance, it is enough to check that $\int \mathcal{L}_n f(\eta) \nu_\rho^n(d\eta) = 0$, for any function $f : \Omega_n \rightarrow \mathbb{R}$. This is a long albeit elementary computation, which is omitted here.

In what follows we consider that the initial measure of the Markov process $\{\eta_{tn^2}^n; t \in [0, T]\}$ is the invariant measure ν_ρ^n . For short, we denote the probability measure $\mathbb{P}_{\nu_\rho^n}$ by \mathbb{P}_ρ and the corresponding expectation by \mathbb{E}_ρ .

The *Dirichlet form* \mathfrak{D}_n is the functional acting on functions $f : \Omega_n \rightarrow \mathbb{R}$ which belong to $L^2(\nu_\rho^n)$ as:

$$\mathfrak{D}_n(f) = \int f(\eta)(-\mathcal{L}_n f)(\eta) \nu_\rho^n(d\eta). \quad (2.2)$$

Invoking a general result [31, Appendix 1, Prop. 10.1] we can rewrite \mathfrak{D}_n as

$$\mathfrak{D}_n(f) = \mathfrak{D}_n^{\text{bulk}}(f) + \mathfrak{D}_n^{\text{bnd}}(f), \quad (2.3)$$

where

$$\begin{aligned} \mathfrak{D}_n^{\text{bulk}}(f) &= \sum_{x=1}^{n-2} \int \left(1 + \frac{E}{n^\gamma}\right) (\eta(x) - \eta(x+1))^2 (\nabla_{x,x+1} f(\eta))^2 \nu_\rho^n(d\eta) \\ \mathfrak{D}_n^{\text{bnd}}(f) &= \int r_{0,1}(\eta) (f(\sigma^1 \eta) - f(\eta))^2 \nu_\rho^n(d\eta) \\ &\quad + \int r_{n-1,n}(\eta) (f(\sigma^{n-1} \eta) - f(\eta))^2 \nu_\rho^n(d\eta). \end{aligned}$$

2.3. The spaces of test functions. Let $\mathcal{C}^\infty([0, 1])$ be the space of real valued functions $\varphi : [0, 1] \rightarrow \mathbb{R}$ such that φ is continuous in $[0, 1]$ as well as all its derivatives. We denote by $d^k \varphi$ the k -th derivative of φ , and for $k = 1$ (resp. $k = 2$) we simply denote it by $\nabla \varphi$ (resp. $\Delta \varphi$).

Before defining the *fluctuation fields* associated to our process, we first need to introduce the suitable space for test functions for each one of the fields that we will analyze. First of all, let \mathcal{S}_{Dir} and \mathcal{S}_{Neu} be the following spaces of functions:

$$\mathcal{S}_{\text{Dir}} = \left\{ \varphi \in \mathcal{C}^\infty([0, 1]); d^{2k} \varphi(0) = d^{2k} \varphi(1) = 0, \text{ for any } k \in \mathbb{N}_0 \right\} \quad (2.4)$$

$$\mathcal{S}_{\text{Neu}} = \left\{ \varphi \in \mathcal{C}^\infty([0, 1]); d^{2k+1} \varphi(0) = d^{2k+1} \varphi(1) = 0, \text{ for any } k \in \mathbb{N}_0 \right\}, \quad (2.5)$$

both equipped with the family of seminorms $\left\{ \sup_{u \in [0, 1]} |d^k \varphi(u)| \right\}_{k \in \mathbb{N}_0}$. Then \mathcal{S}_{Dir} and \mathcal{S}_{Neu} are Fréchet spaces, and we write $\mathcal{S}'_{\text{Dir}}$ and $\mathcal{S}'_{\text{Neu}}$ for their topological duals.

Now, let $-\Delta$ be the closure of the Laplacian operator $-\Delta : \mathcal{S}_{\text{Dir}} \rightarrow L^2([0, 1])$ as an unbounded operator in the Hilbert space $L^2([0, 1])$. It is a positive and self-adjoint operator whose eigenvalues and eigenfunctions are given for any integer $m \geq 1$, respectively by $\lambda_m = (m\pi)^2$ and $e_m(u) = \sqrt{2} \sin(m\pi u)$. We note that the set $\{e_m; m \geq 1\}$

forms an orthonormal basis of $L^2([0, 1])$. We denote by $\{T_t^{\text{Dir}} ; t \geq 0\}$ the semi-group associated to Δ .

For $k \in \mathbb{N}$, let us define $\mathcal{H}_{\text{Dir}}^k = \text{Dom}((-\Delta)^{k/2})$, endowed with the inner product

$$(\varphi, \psi)_k = \int_0^1 (-\Delta)^{k/2} \varphi(u) (-\Delta)^{k/2} \psi(u) du.$$

By the spectral theorem for self-adjoint operators, $\mathcal{H}_{\text{Dir}}^k$ equals

$$\mathcal{H}_{\text{Dir}}^k = \left\{ \varphi \in L^2([0, 1]) ; \|\varphi\|_k^2 < \infty \right\},$$

where $\|\varphi\|_k^2 = (\varphi, \varphi)_k$ and for $\varphi, \psi \in L^2([0, 1])$ we also have

$$(\varphi, \psi)_k = \sum_{m=1}^{+\infty} \left\{ (m\pi)^{2k} \langle \varphi, e_m \rangle_{L^2([0,1])} \langle \psi, e_m \rangle_{L^2([0,1])} \right\}.$$

Moreover, if $\mathcal{H}_{\text{Dir}}^{-k}$ denotes the topological dual space of $\mathcal{H}_{\text{Dir}}^k$, then

$$\mathcal{H}_{\text{Dir}}^{-k} = \left\{ \mathcal{Y} \in \mathcal{S}'_{\text{Dir}} ; \|\mathcal{Y}\|_{-k}^2 < \infty \right\},$$

where $\|\mathcal{Y}\|_{-k}^2 = (\mathcal{Y}, \mathcal{Y})_{-k}$ and the inner product $(\cdot, \cdot)_{-k}$ is defined as

$$(\mathcal{Y}, \mathcal{Z})_{-k} = \sum_{m=1}^{+\infty} \left\{ (m\pi)^{-2k} \mathcal{Y}(e_m) \mathcal{Z}(e_m) \right\}, \quad (2.6)$$

with $\mathcal{Y}(e_m)$ denoting the action of the distribution \mathcal{Y} on e_m .

Finally, we define Δ_{Neu} (resp. Δ_{Dir}) as the Neumann (resp. Dirichlet) Laplacian acting on $\varphi \in \mathcal{S}_{\text{Neu}}$ and $\mathcal{Y} \in \mathcal{S}'_{\text{Neu}}$ (resp. $\varphi \in \mathcal{S}_{\text{Dir}}$ and $\mathcal{Y} \in \mathcal{S}'_{\text{Dir}}$) as follows:

$$\Delta_{\text{Neu}} \mathcal{Y}(\varphi) = \mathcal{Y}(\Delta \varphi), \quad \text{resp.} \quad \Delta_{\text{Dir}} \mathcal{Y}(\varphi) = \mathcal{Y}(\Delta \varphi).$$

Let $\mathcal{H}_{\text{Neu}}^{-k}$ be the topological dual space of $\mathcal{H}_{\text{Neu}}^k$, both defined similarly to $\mathcal{H}_{\text{Dir}}^k$ and $\mathcal{H}_{\text{Dir}}^{-k}$ but replacing the basis $e_m(u) = \sqrt{2} \sin(m\pi u)$ by $\tilde{e}_m(u) = \sqrt{2} \cos(m\pi u)$. In the next sections we will also need to consider one last operator, denoted by ∇_{Dir} and defined as follows: given $k > 0$, $\mathcal{Y} \in \mathcal{S}'_{\text{Neu}}$ and $\varphi \in \mathcal{S}_{\text{Dir}}$ we set

$$\nabla_{\text{Dir}} \mathcal{Y}(\varphi) = -\mathcal{Y}(\nabla \varphi). \quad (2.7)$$

2.4. Fluctuation fields. Now we introduce all the microscopic fluctuation fields for which we will prove convergence to some infinite dimensional stochastic partial differential equations (SPDEs). In the following, for any space \mathcal{S} of distributions, we denote by $\mathcal{D}([0, T], \mathcal{S})$ (resp. $\mathcal{C}([0, T], \mathcal{S})$) the set of càdlàg (resp. continuous) trajectories taking values in \mathcal{S} and endowed with the Skorohod topology.

2.4.1. Density fluctuation field. Since the particle system is stationary, the microscopic average $\mathbb{E}_\rho[\eta_{tn^2}^n(x)]$ is constantly equal to ρ . We are therefore looking at the fluctuations of the microscopic configurations around their mean, namely:

Definition 2.1. For any $t > 0$, let \mathcal{Y}_t^n be the density fluctuation field which acts on functions $\varphi \in \mathcal{S}_{\text{Dir}}$ as

$$\mathcal{Y}_t^n(\varphi) = \frac{1}{\sqrt{n}} \sum_{x=1}^{n-1} \varphi\left(\frac{x}{n}\right) (\eta_{tn^2}^n(x) - \rho).$$

In the following, we will see that the right space one has to consider in order to prove the convergence of the field \mathcal{Y}_t^n is the space $\mathcal{H}_{\text{Dir}}^{-k}$, with $k > \frac{5}{2}$. Note that, for any $t > 0$ and any $k > \frac{5}{2}$, \mathcal{Y}_t^n is indeed an element of $\mathcal{H}_{\text{Dir}}^{-k}$ and $\mathcal{Y}^n \in \mathcal{D}([0, T], \mathcal{H}_{\text{Dir}}^{-k})$.

2.4.2. Height function. Alternatively, instead of working with the density fluctuation field of the particle system, we may also consider the height function associated to it. Roughly speaking, the height function integrates the density fluctuation field in the space variable, i.e. it describes a curve h from $\{1, \dots, n\}$ to \mathbb{R} which satisfies $h(x+1) = h(x) + (\eta(x) - \rho)$. This suggests the definition $h(x) = \sum_{y=1}^{x-1} (\eta(y) - \rho)$, which however has the disadvantage that if a new particle enters at site 1, then $h(x)$ increases by 1 for all x . Therefore, we set

$$h_t^n(x) = h_t^n(1) + \sum_{y=1}^{x-1} (\eta_t^n(y) - \rho), \quad \text{for any } x \in \{1, \dots, n\}, \quad (2.8)$$

where, by definition, $h_0^n(1) = 0$ and

- $h_t^n(1)$ increases by 1 whenever a particle at site 1 leaves the system to the left,
- $h_t^n(1)$ decreases by 1 whenever a new particle enters the system at the site 1.

In other words $h_t^n(1)$ is exactly the net number of particles removed from the system at the left boundary until time tn^2 .

The weak asymmetry in the particle system causes the height function to slowly decrease because $E > 0$ (it would grow if $E < 0$), and at the first order the decrease is of order $c_n t$ for

$$c_n = -n^{2-\gamma} \rho(1-\rho)E. \quad (2.9)$$

For instance, with our stationary dynamics, one can easily see (see (4.26) for the case $x = 1$ but note that for the other values of x we also have a similar expression) that

$$\mathbb{E}_\rho[h_t^n(x)] = \mathbb{E}_\rho[h_0^n(x)] - \frac{E}{n^\gamma} \rho(1-\rho)t = c_n t, \quad \text{for any } x \in \{1, \dots, n\}.$$

In the case $\gamma = \frac{1}{2}$ (where we will see KPZ behavior) note that $c_n = -n^{3/2} \rho(1-\rho)E$. So while on the microscopic scale the decrease is negligible, it becomes important on time scales of order n^2 . We want to investigate the fluctuations around the flat interface $c_n t$, namely:

Definition 2.2. For any $t > 0$, let \mathcal{Z}_t^n be the height fluctuation field which acts on functions $\varphi \in \mathcal{S}_{\text{Neu}}$ as

$$\mathcal{Z}_t^n(\varphi) = \frac{1}{n^{3/2}} \sum_{x=1}^n \varphi\left(\frac{x}{n}\right) (h_{tn^2}^n(x) - c_n t).$$

Remark 2.3. We will see, as expected, that \mathcal{Z}_t^n and \mathcal{Y}_t^n are related. Note that: if $\varphi \in \mathcal{S}_{\text{Dir}}$, then $\nabla \varphi \in \mathcal{S}_{\text{Neu}}$ and, a simple computation, based on a summation by parts, shows that $\mathcal{Z}_t^n(\tilde{\nabla}_n \varphi)$ can be written as $-\mathcal{Y}_t^n(\varphi)$ plus the boundary terms:

$$-\frac{1}{\sqrt{n}} \varphi\left(\frac{1}{n}\right) (h_{tn^2}^n(1) - c_n t) + \frac{1}{\sqrt{n}} \varphi\left(\frac{n}{n}\right) (h_{tn^2}^n(n) - c_n t).$$

Above $\tilde{\nabla}_n \varphi$ is essentially a discrete gradient, see (4.27) below. Note that since $\varphi \in \mathcal{S}_{\text{Dir}}$, the last expression vanishes in $L^2(\mathbb{P}_\rho)$ as $n \rightarrow \infty$, as a consequence of Lemma 4.5 given ahead.

Below, we will prove that the convergence of the field \mathcal{Z}_t^n is in the space $\mathcal{H}_{\text{Neu}}^{-k}$, with $k > \frac{5}{2}$ and we note that $\mathcal{Z}^n \in \mathcal{D}([0, T], \mathcal{H}_{\text{Neu}}^{-k})$.

3. SOLUTIONS TO NON-LINEAR SPDES AND STATEMENT OF THE RESULTS

In this section, we first define properly the notion of solutions for four stochastic partial differential equations, all with boundary conditions (Sections 3.1, 3.2, 3.3 and 3.4), and we connect them to each other, especially through their boundary conditions (Section 3.5). These SPDEs are going to describe the macroscopic behavior of the fluctuation fields of our system, for different values of the parameter γ ruling the strength of the asymmetry: the precise statements are then given in Section 3.6. For the sake of clarity, the proofs will be postponed to further sections.

3.1. Ornstein-Uhlenbeck process with Dirichlet boundary condition. Let us start with the *generalized Ornstein-Uhlenbeck process* described by the formal SPDE

$$d\mathcal{Y}_t = A\Delta_{\text{Dir}} \mathcal{Y}_t dt + \sqrt{D} \nabla_{\text{Dir}} d\mathcal{W}_t, \quad (3.1)$$

with $A, D > 0$ and where $\{\mathcal{W}_t ; t \in [0, T]\}$ is a standard $\mathcal{S}'_{\text{Neu}}$ -valued Brownian motion, with covariance

$$\mathbb{E}[\mathcal{W}_s(\varphi)\mathcal{W}_t(\psi)] = (s \wedge t) \langle \varphi, \psi \rangle_{L^2([0,1])}. \quad (3.2)$$

Since $\{\mathcal{W}_t ; t \geq 0\}$ is $\mathcal{S}'_{\text{Neu}}$ -valued, then $\nabla_{\text{Dir}} \mathcal{W}_t$ is a well defined object in $\mathcal{S}'_{\text{Dir}}$.

The following proposition aims at defining in a unique way the stochastic process solution to (3.1):

Proposition 3.1. *There exists a unique (in law) stochastic process $\{\mathcal{Y}_t ; t \in [0, T]\}$, with trajectories in $\mathcal{C}([0, T], \mathcal{S}'_{\text{Dir}})$, called Ornstein-Uhlenbeck process solution of (3.1), such that:*

- (1) *the process $\{\mathcal{Y}_t ; t \in [0, T]\}$ is “stationary”, in the following sense: for any $t \in [0, T]$, the random variable \mathcal{Y}_t is a $\mathcal{S}'_{\text{Dir}}$ -valued space white noise with covariance given on $\varphi, \psi \in \mathcal{S}_{\text{Dir}}$ by*

$$\mathbb{E}[\mathcal{Y}_t(\varphi)\mathcal{Y}_t(\psi)] = \frac{D}{2A} \langle \varphi, \psi \rangle_{L^2([0,1])};$$

- (2) *for any $\varphi \in \mathcal{S}_{\text{Dir}}$,*

$$\mathcal{M}_t(\varphi) = \mathcal{Y}_t(\varphi) - \mathcal{Y}_0(\varphi) - A \int_0^t \mathcal{Y}_s(\Delta \varphi) ds$$

is a continuous martingale with respect to the natural filtration associated to \mathcal{Y} , namely

$$\mathcal{F}_t = \sigma(\mathcal{Y}_s(\varphi) ; s \leq t, \varphi \in \mathcal{S}_{\text{Dir}}), \quad (3.3)$$

of quadratic variation

$$\langle \mathcal{M}(\varphi) \rangle_t = tD \|\nabla \varphi\|_{L^2([0,1])}^2.$$

Moreover, for every function $\varphi \in \mathcal{S}_{\text{Dir}}$, the stochastic process $\{\mathcal{Y}_t(\varphi) ; t \in [0, T]\}$ is Gaussian, and the distribution of $\mathcal{Y}_t(\varphi)$ conditionally to $\{\mathcal{Y}_u ; u < s\}$, is normal of mean $\mathcal{Y}_s(T_{A(t-s)}^{\text{Dir}} \varphi)$ and variance

$$D \int_0^{t-s} \|\nabla(T_{Ar}^{\text{Dir}} \varphi)\|_{L^2([0,1])}^2 dr.$$

Proof of Proposition 3.1. The proof of this fact is standard, and we refer the interested reader to [31, 33]. \square

3.2. Stochastic Burgers equation with Dirichlet boundary condition. Now we define the notion of *stationary energy solutions* for the *stochastic Burgers equation with Dirichlet boundary condition*, written as

$$d\mathcal{Y}_t = A \Delta_{\text{Dir}} \mathcal{Y}_t dt + \bar{E} \nabla_{\text{Dir}} (\mathcal{Y}_t^2) dt + \sqrt{D} \nabla_{\text{Dir}} d\mathcal{W}_t, \quad (3.4)$$

with $A, D > 0$ and $\bar{E} \in \mathbb{R}$, and where $\{\mathcal{W}_t ; t \in [0, T]\}$ is a $\mathcal{S}'_{\text{Neu}}$ -valued Brownian motion with covariance (3.2). Note that $\nabla_{\text{Dir}} (\mathcal{Y}_t^2)$ so far has not a precise meaning but this will be rigorously given below in Theorem 3.3. More precisely, we are going to adapt to our case (i.e. adding boundary conditions) the notion of energy solutions as given for the first time in [17, 23].

Definition 3.2. For $u \in [0, 1]$ and $\varepsilon > 0$, let $\iota_\varepsilon(u) : [0, 1] \rightarrow \mathbb{R}$ be the approximation of the identity defined as

$$\iota_\varepsilon(u)(v) := \begin{cases} \varepsilon^{-1} \mathbf{1}_{]u, u+\varepsilon]}(v) & \text{if } u \in [0, 1-2\varepsilon), \\ \varepsilon^{-1} \mathbf{1}_{[u-\varepsilon, u[}(v) & \text{if } u \in [1-2\varepsilon, 1]. \end{cases}$$

The following theorem, which we prove in this paper, aims at defining uniquely the stochastic process solution to (3.4):

Theorem 3.3. *There exists a unique (in law) stochastic process $\{\mathcal{Y}_t ; t \in [0, T]\}$ with trajectories in $\mathcal{C}([0, T], \mathcal{S}'_{\text{Dir}})$, called stationary energy solution of (3.4), such that:*

- (1) *the process $\{\mathcal{Y}_t ; t \in [0, T]\}$ is “stationary”, in the following sense: for any $t \in [0, T]$, the random variable \mathcal{Y}_t is a $\mathcal{S}'_{\text{Dir}}$ -valued space white noise with covariance given on $\varphi, \psi \in \mathcal{S}_{\text{Dir}}$ by*

$$\mathbb{E}[\mathcal{Y}_t(\varphi)\mathcal{Y}_t(\psi)] = \frac{D}{2A} \langle \varphi, \psi \rangle_{L^2([0,1])};$$

- (2) *the process $\{\mathcal{Y}_t ; t \in [0, T]\}$ satisfies the following energy estimate: there exists $\kappa > 0$ such that for any $\varphi \in \mathcal{S}_{\text{Dir}}$, any $\delta, \varepsilon \in (0, 1)$ with $\delta < \varepsilon$, and any $s, t \in [0, T]$ with $s < t$,*

$$\mathbb{E}[(\mathcal{A}_{s,t}^\varepsilon(\varphi) - \mathcal{A}_{s,t}^\delta(\varphi))^2] \leq \kappa(t-s)\varepsilon \|\nabla \varphi\|_{L^2([0,1])}^2, \quad (3.5)$$

where

$$\mathcal{A}_{s,t}^\varepsilon(\varphi) = - \int_s^t \int_0^1 (\mathcal{Y}_r(\iota_\varepsilon(u)))^2 \nabla \varphi(u) du dr; \quad (3.6)$$

- (3) *for any $\varphi \in \mathcal{S}_{\text{Dir}}$ and any $t \in [0, T]$, the process*

$$\mathcal{M}_t(\varphi) = \mathcal{Y}_t(\varphi) - \mathcal{Y}_0(\varphi) - A \int_0^t \mathcal{Y}_s(\Delta \varphi) ds - \bar{E} \mathcal{A}_t(\varphi)$$

is a continuous martingale with respect to the natural filtration (3.3) associated to \mathcal{Y} , of quadratic variation

$$\langle \mathcal{M}(\varphi) \rangle_t = tD \|\nabla \varphi\|_{L^2([0,1])}^2,$$

where the process $\{\mathcal{A}_t ; t \in [0, T]\}$ is obtained through the following limit, which holds in the L^2 -sense:

$$\mathcal{A}_t(\varphi) = \lim_{\varepsilon \rightarrow 0} \mathcal{A}_{0,t}^\varepsilon(\varphi); \quad (3.7)$$

- (4) *the reversed process $\{\mathcal{Y}_{T-t} ; t \in [0, T]\}$ satisfies item (3) with \bar{E} replaced by $-\bar{E}$.*

Proof of Theorem 3.3. The proof is somewhat lengthy and we give it in Section 6.3. \square

Remark 3.4. Note that there is a way to make sense of $\mathcal{Y}_s(\iota_\varepsilon(u))$ (which is a priori not obvious since $\iota_\varepsilon(u)$ is not a test function), as explained in [17, Remark 4].

Remark 3.5. Note that the definition of the Ornstein-Uhlenbeck process \mathcal{Y} given in Proposition 3.1 and of the energy solution to the stochastic Burgers equation (3.4) when $\bar{E} = 0$, are equivalent. Indeed, the only part which is not obvious is that the Ornstein-Uhlenbeck process \mathcal{Y} satisfies conditions (2) and (4) in Theorem 3.3. But both of these will follow from our convergence result Theorem 3.17 with $\bar{E} = 0$.

3.3. KPZ equation with Neumann boundary condition. Here we define the notion of solution for the KPZ equation with Neumann boundary condition, which is formally given by

$$d\mathcal{Z}_t = A \Delta_{\text{Neu}} \mathcal{Z}_t dt + \bar{E} (\nabla \mathcal{Z}_t)^{\circ 2} dt + \sqrt{D} d\mathcal{W}_t, \quad (3.8)$$

with $A, D > 0$ and $\bar{E} \in \mathbb{R}$, and where $\{\mathcal{W}_t ; t \in [0, T]\}$ is a standard $\mathcal{S}'_{\text{Neu}}$ -valued Brownian motion with covariance (3.2) and $(\nabla \mathcal{Z}_t)^{\circ 2}$ denotes a *renormalized square* which will have a precise meaning in Theorem 3.7 below.

Definition 3.6. For $\varepsilon > 0$ and $\mathcal{Z} \in \mathcal{C}([0, 1])$, let us define

$$\nabla_\varepsilon \mathcal{Z}(u) = \begin{cases} \varepsilon^{-1}(\mathcal{Z}(u + \varepsilon) - \mathcal{Z}(u)), & \text{if } u \in [0, 1 - 2\varepsilon], \\ \varepsilon^{-1}(\mathcal{Z}(u) - \mathcal{Z}(u - \varepsilon)), & \text{if } u \in [1 - 2\varepsilon, 1]. \end{cases} \quad (3.9)$$

The following theorem aims at defining uniquely the stochastic process solution to (3.8):

Theorem 3.7. Let Z be a random variable with values in $\mathcal{C}([0, 1])$, such that $\nabla_{\text{Dir}} Z$ is a white noise with variance $D/(2A)$ and such that $\sup_{u \in [0, 1]} \mathbb{E}[e^{2Z(u)}] < \infty$.

There exists a unique (in law) stochastic process $\{\mathcal{Z}_t ; t \in [0, T]\}$ with trajectories in $\mathcal{C}([0, T], \mathcal{S}'_{\text{Neu}})$, called almost stationary energy solution of (3.8), such that:

- (1) $\text{law}(\mathcal{Z}_0) = \text{law}(Z)$;
- (2) there exists a stationary energy solution $\{\mathcal{Y}_t ; t \in [0, T]\}$ to the stochastic Burgers equation (3.4) – with trajectories in $\mathcal{C}([0, T], \mathcal{S}'_{\text{Dir}})$ – such that $\nabla_{\text{Dir}} \mathcal{Z} = \mathcal{Y}$;
- (3) for any $\varphi \in \mathcal{S}_{\text{Neu}}$ and any $t \in [0, T]$, the process

$$\mathcal{N}_t(\varphi) = \mathcal{Z}_t(\varphi) - \mathcal{Z}_0(\varphi) - A \int_0^t \mathcal{Z}_s(\Delta \varphi) ds - \bar{E} \mathcal{B}_t(\varphi)$$

is a continuous martingale with respect to the natural filtration associated to \mathcal{Z} , defined as in (3.3) but with $\varphi \in \mathcal{S}_{\text{Neu}}$ and with \mathcal{Z}_s in the place of \mathcal{Y}_s , of quadratic variation

$$\langle \mathcal{N}(\varphi) \rangle_t = tD \|\varphi\|_{L^2([0, 1])}^2,$$

where the process $\{\mathcal{B}_t ; t \in [0, T]\}$ is obtained through the following limit, which holds in the L^2 -sense:

$$\mathcal{B}_t(\varphi) = \lim_{\varepsilon \rightarrow 0} \int_0^t \int_0^1 \left\{ (\nabla_\varepsilon \mathcal{Z}_s(u))^2 - \frac{D}{2A\varepsilon} \right\} \varphi(u) du ds \quad (3.10)$$

where $\nabla_\varepsilon \mathcal{Z}_s(u)$ has been defined in (3.9).

Proof of Theorem 3.7. The proof of that result is also contained in Section 6.3. □

Remark 3.8. Note that we did not require \mathcal{Z} to be a continuous function in u , so it is not obvious that $\nabla_\varepsilon \mathcal{Z}(u)$ as defined in (3.9) is well defined. But of course $\nabla_\varepsilon \mathcal{Z}(u) = \mathcal{Y}(\iota_\varepsilon(u))$, and as discussed before the right hand side can be made sense of with the arguments of [17, Remark 4].

Remark 3.9. By “almost stationary” we mean that the derivative of \mathcal{Z} is stationary: It is a (multiple of the) white noise for all times, so \mathcal{Z} is a Brownian motion for all times. But the starting point $\mathcal{Z}_t(0)$ of this Brownian motion will change with time.

3.4. Stochastic heat equation with Robin boundary conditions. Finally we define the notion of solutions to the *stochastic heat equation with Robin boundary condition*, written as

$$d\Phi_t = A \Delta_{\text{Rob}} \Phi_t dt + \sqrt{D} \Phi_t d\mathcal{W}_t, \quad (3.11)$$

with $A, D > 0$ and $\{\mathcal{W}_t; t \in [0, T]\}$ is a standard $\mathcal{S}'_{\text{Neu}}$ -valued Brownian motion with covariance (3.2), and where we want to impose (formally):

$$\nabla \Phi_t(0) = \alpha \Phi_t(0), \quad \nabla \Phi_t(1) = \beta \Phi_t(1), \quad (3.12)$$

with $\alpha, \beta \in \mathbb{R}$. To see how we should implement these boundary conditions, consider $f \in \mathcal{C}^2([0, 1])$ and $\varphi \in \mathcal{S}'_{\text{Neu}}$, and assume that f satisfies the Robin boundary conditions $\nabla f(0) = \alpha f(0)$, and $\nabla f(1) = \beta f(1)$. Integrating by parts twice and using the Neumann boundary conditions for φ and the Robin boundary conditions for f , we obtain

$$\int_0^1 \Delta f(u) \varphi(u) du = \int_0^1 f(u) \Delta \varphi(u) du - \alpha f(0) \varphi(0) + \beta f(1) \varphi(1).$$

This suggests to define the operator

$$\Delta_{\text{Rob}}: \mathcal{S}'_{\text{Neu}} \cap \mathcal{C}([0, 1]) \rightarrow \mathcal{S}'_{\text{Neu}}, \quad \Delta_{\text{Rob}} f(\varphi) = f(\Delta \varphi) - \alpha f(0) \varphi(0) + \beta f(1) \varphi(1),$$

and in principle this leads to a definition of solutions to (3.11). But for technical reasons we do not want to a priori assume our solution to be continuous in (t, u) , and this means we could change the values of $\Phi_t(0)$ and $\Phi_t(1)$ without changing $\Phi_t(\varphi)$, so we cannot hope to get uniqueness without further assumptions. Let us introduce a suitable class of processes for which the boundary term is well defined: we write $\Phi \in \mathcal{L}^2_{\mathcal{C}}([0, T])$ if Φ takes values in the Borel measurable functions on $[0, T] \times [0, 1]$, if

$$\sup_{(u,t) \in [0,1] \times [0,T]} \mathbb{E}[\Phi_t(u)^2] < \infty,$$

and if for all $t > 0$ the process $u \mapsto \Phi_t(u)$ is continuous in $L^2(\mathbb{P})$, where \mathbb{P} is the law of the process Φ . and above \mathbb{E} is the expectation w.r.t. to \mathbb{P} . For $\Phi \in \mathcal{L}^2_{\mathcal{C}}([0, T])$ we cannot change the value of $\Phi_t(0)$ without changing Φ_t in an environment of 0, and this would also change $\Phi_t(\varphi)$ for some test functions φ .

Proposition 3.10. Let F be a random variable with values in the Borel measurable functions on $[0, 1]$, such that $\sup_{u \in [0,1]} \mathbb{E}[F(u)^2] < \infty$.

Then there exists at most one (law of a) process $\{\Phi_t; t \in [0, T]\} \in \mathcal{L}^2_{\mathcal{C}}([0, T])$, called weak solution to (3.11) with boundary condition (3.12), such that:

- (1) $\text{law}(\Phi_0) = \text{law}(F)$;

(2) for any $\varphi \in \mathcal{S}_{\text{Neu}}$ and any $t \in [0, T]$, the process

$$\mathfrak{M}_t(\varphi) = \Phi_t(\varphi) - \Phi_0(\varphi) - A \int_0^t \Phi_s(\Delta \varphi) ds - A \int_0^t (-\alpha \Phi_s(0)\varphi(0) + \beta \Phi_s(1)\varphi(1)) ds \quad (3.13)$$

is a continuous martingale of quadratic variation

$$\langle \mathfrak{M}(\varphi) \rangle_t = D \int_0^t \int_0^1 |\Phi_s(u)\varphi(u)|^2 du ds.$$

Proof. The proof of this proposition is given in Appendix D. \square

Remark 3.11. In [5, Proposition 2.7] the authors show uniqueness for (the mild solution of) the same equation, but they require $\alpha > 0$ and $\beta < 0$, while here we will need to take $\alpha = -E^2/2$ and $\beta = E^2/2$ (where $E > 0$ is the parameter ruling the weak asymmetry). A couple of weeks after this paper was displayed, Parekh in [37] extended the approach of [5] to general parameters α, β . As the probabilistic representation given in [13, 36] shows, the key difference is that our boundary conditions correspond to sources at the boundary, while the case of [5] corresponds to sinks. Such sink terms make the Robin Laplacian negative and in particular its spectrum is negative, while for the Robin Laplacian with a source term some eigenvalues may be positive. In particular, the method used in Section 4.1 of [5] for proving heat kernel estimates breaks down for our choice of parameters. Theorem 3.4 of [36] gives L^1 and L^∞ bounds for the heat kernel that are sufficient for our purposes, but it is not quite obvious whether the heat kernel is in $\mathcal{C}^{1,2}((0, \infty) \times [0, 1])$ and satisfies the Robin boundary condition in the strong sense, which is what we would need here. While we expect this to be true, we avoid the problem by formulating the equation slightly differently: we do not work with the Robin heat kernel but we use the Neumann heat kernel instead and deal with the resulting boundary corrections by hand.

Remark 3.12. It would be possible to weaken the assumptions on the initial conditions (say to allow for a Dirac delta initial condition), and of course we would also be able to prove existence and not just uniqueness. Since given the computations in Appendix D these results follow from standard arguments as in [38, 41] and we do not need them here, we choose not to include the proofs.

We conclude this section by making the link between the KPZ equation (3.8) and the stochastic heat equation (3.11), with their respective boundary conditions. This link is, as expected, done through the Cole-Hopf transformation, although the boundary conditions are linked in a more complicated way than one might guess naively:

Proposition 3.13. Let \mathcal{X} be the almost stationary energy solution of the KPZ equation (3.8) with Neumann boundary condition as defined in Theorem 3.7. Then, for any $t \in [0, T]$ we have

$$\mathcal{X}_t = \frac{A}{E} \log \Phi_t + \frac{D^2(\bar{E})^3}{48A^4} t,$$

where Φ solves the stochastic heat equation

$$d\Phi_t = A\Delta_{\text{Rob}}\Phi_t dt + \frac{\sqrt{D}\bar{E}}{A}\Phi_t d\mathcal{N}_t$$

with the Robin boundary condition

$$\nabla \Phi_t(0) = -\frac{D(\bar{E})^2}{4A^3}\Phi_t(0), \quad \nabla \Phi_t(1) = \frac{D(\bar{E})^2}{4A^3}\Phi_t(1),$$

and where above \mathcal{N}_t has been defined in item (3) of Theorem 3.7.

Proof of Proposition 3.13. This proof is contained in Section 6 (see in particular Subsection 6.3). \square

3.5. Boundary behavior for singular SPDEs. The stochastic Burgers equation is an important example of a *singular SPDE*, a class of equations for which tremendous progress was made in the last years [28, 22]. The vast majority of all papers on singular SPDEs treats domains without boundaries (mostly the torus, sometimes Euclidean space), and only quite recently some articles appeared that deal with boundaries [5, 15].

But in some cases it is not quite obvious how to formulate the boundary conditions. For example, the solution to the stochastic Burgers equation is almost surely distribution valued and therefore we cannot evaluate it at the boundary. In Theorem 3.3 we proposed a formulation for the notion of solutions to the stochastic Burgers equation with Dirichlet boundary condition $\mathcal{Y}_t(0) \equiv \mathcal{Y}_t(1) \equiv 0$ (for any $t \in [0, T]$), that seems natural to us. But as we saw in Proposition 3.13, we have $\mathcal{Y} = \nabla_{\text{Dir}}(\frac{A}{E} \log \Phi)$, where Φ solves the stochastic heat equation with Robin boundary condition $\nabla \Phi_t(0) \equiv -\frac{D(\bar{E})^2}{4A^3} \Phi_t(0)$ and $\nabla \Phi_t(1) \equiv \frac{D(\bar{E})^2}{4A^3} \Phi_t(1)$.

Moreover, Gerencsér and Hairer observe in [15, Theorem 1.6] that the classical solution \mathcal{Y}^ε to the stochastic Burgers equation without renormalization,

$$d\mathcal{Y}_t^\varepsilon = \Delta_{\text{Dir}} \mathcal{Y}_t^\varepsilon dt + \nabla_{\text{Dir}}((\mathcal{Y}_t^\varepsilon)^2) dt + \nabla_{\text{Dir}} d\mathcal{W}_t^\varepsilon, \quad (3.14)$$

with Dirichlet boundary condition $\mathcal{Y}_t^\varepsilon(0) \equiv \mathcal{Y}_t^\varepsilon(1) \equiv 0$ and where \mathcal{W}^ε is a mollification of \mathcal{W} , may converge to different limits satisfying different boundary conditions as $\varepsilon \rightarrow 0$, depending on which mollifier was used for \mathcal{W}^ε . But if the noise is only mollified in space and white in time, then the limit is always the same and it agrees with the Cole-Hopf solution of [5].

So it is not obvious whether there is a “canonical” way of formulating singular SPDEs with boundary conditions, and in the case of the stochastic Burgers equation with Dirichlet boundary conditions it may seem that the Cole-Hopf solution is the most canonical solution. But below, in Proposition 3.14, we will see that our solution \mathcal{Y} indeed satisfies $\mathcal{Y}_t(0) \equiv \mathcal{Y}_t(1) \equiv 0$, as long as we do not try to evaluate $\mathcal{Y}_t(0)$ at a fixed time but allow for a bit of averaging in time instead. We also show, in Lemma 3.15, that the (approximate) Cole-Hopf transformation Ψ of \mathcal{Y} satisfies the Robin boundary condition $\nabla \Psi_t(0) \equiv -\frac{D(\bar{E})^2}{4A^3} \Psi_t(0)$ and $\nabla \Psi_t(1) \equiv \frac{D(\bar{E})^2}{4A^3} \Psi_t(1)$ after averaging in time. This sheds some light on the actual boundary behavior of solutions to singular SPDEs and indicates that our formulation of the equation is maybe more natural than the Cole-Hopf formulation of [5]. Although the approximation result in (3.14) looks natural at first sight, note that *not* renormalizing $\nabla_{\text{Dir}}((\mathcal{Y}_t^\varepsilon)^2)$ means that we renormalize $(\mathcal{Y}_t^\varepsilon)^2$ with a constant which is killed by ∇_{Dir} and therefore does not appear in the equation. But of course \mathcal{Y}^ε is not spatially homogeneous, so there is no reason why the renormalization should be spatially constant. And as our results show, taking it constant actually leads to an unnatural boundary behavior in the limit.

Proposition 3.14. *Let \mathcal{Y} be the solution to the stochastic Burgers equation (3.4) as defined in Theorem 3.3, and let $\{\rho_\varepsilon\}_{\varepsilon>0} \subset L^2([0, 1])$ be a sequence of positive functions that converges weakly to the Dirac delta at 0 (respectively 1), in the sense that, for any function $f \in \mathcal{C}([0, 1])$, $\lim_{\varepsilon \rightarrow 0} \int_0^1 \rho_\varepsilon(u) f(u) du = f(0)$ (respectively $f(1)$).*

Then we have for all $p \in [1, \infty)$

$$\lim_{\varepsilon \rightarrow 0} \mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t \mathcal{Y}_s(\rho_\varepsilon) ds \right|^p \right] = 0.$$

Proof of Proposition 3.14. We prove this proposition in Appendix B. \square

Consider now, for any $u, v \in [0, 1]$

$$\Theta_u(v) = \mathbf{1}_{[0, u]}(v) + v - 1.$$

This is an integral kernel which will be used in Section 6.3 to map the energy solution \mathcal{Y} to an approximate solution of the stochastic heat equation¹. We set $\mathcal{Z}_t(u) = \mathcal{Y}_t(\Theta_u)$, so that $\nabla_{\text{Dir}} \mathcal{Z} = \mathcal{Y}$, and then $\Psi_t(u) = \exp(\frac{\bar{E}}{A} \mathcal{Z}_t(u))$. Then one could compute formally

$$\nabla \Psi_t(0) = \Psi_t(0) \frac{\bar{E}}{A} \nabla \mathcal{Z}_t(0) = \Psi_t(0) \frac{\bar{E}}{A} \mathcal{Y}_t(0) = 0.$$

But we will show now that this formal computation is incorrect and the boundary condition for Ψ is not of Neumann type (i.e. $\nabla \Psi_t(0) \equiv \nabla \Psi_t(1) \equiv 0$), but rather of Robin type, more precisely $\nabla \Psi_t(0) \equiv -\frac{D(\bar{E})^2}{4A^3} \Psi_t(0)$ and $\nabla \Psi_t(1) \equiv \frac{D(\bar{E})^2}{4A^3} \Psi_t(1)$.

Proposition 3.15. *Let \mathcal{Y} be the solution to the stochastic Burgers equation (3.4) as defined in Theorem 3.3 and let $\Psi_t(u) = \exp(\frac{\bar{E}}{A} \mathcal{Y}_t(\Theta_u))$.*

Then, we have for all $t \in [0, T]$

$$\lim_{\varepsilon \rightarrow 0} \mathbb{E} \left[\left| \int_0^t \left(\frac{\Psi_s(\varepsilon) - \Psi_s(0)}{\varepsilon} + \frac{D(\bar{E})^2}{4A^3} \Psi_s(0) \right) ds \right|^2 \right] = 0$$

and

$$\lim_{\varepsilon \rightarrow 0} \mathbb{E} \left[\left| \int_0^t \left(\frac{\Psi_s(1) - \Psi_s(1 - \varepsilon)}{\varepsilon} - \frac{D(\bar{E})^2}{4A^3} \Psi_s(1) \right) ds \right|^2 \right] = 0.$$

Proof of Proposition 3.15. The proof will be exposed in Appendix B. \square

Remark 3.16. *When we naively apply the chain rule in the formal derivation “ $\nabla \mathcal{Z}_t(1) = 1$ and therefore $\nabla \exp(\frac{\bar{E}}{A} \mathcal{Z}_t)(1) = \exp(\frac{\bar{E}}{A} \mathcal{Z}_t(1)) \frac{\bar{E}}{A} \nabla \mathcal{Z}_t(1) = 0$ ”, we implicitly assume that $\mathcal{Z}_t(\cdot)$ allows for a Stratonovich type calculus. But in the space variable $\mathcal{Z}_t(\cdot)$ is $\sqrt{D/(2A)}$ times a Brownian motion, so maybe it is more natural to apply Itô’s formula (as Martin Hairer kindly pointed out to us), with which we get*

$$\begin{aligned} d_x \exp \left(\frac{\bar{E}}{A} \mathcal{Z}_t(x) \right) \Big|_{x=1} &= d_x \exp \left(\frac{\bar{E}}{A} \mathcal{Z}_t(x) \right) \frac{\bar{E}}{A} d_x \mathcal{Z}_t(x) \Big|_{x=1} \\ &\quad + \frac{1}{2} \exp \left(\frac{\bar{E}}{A} \mathcal{Z}_t(x) \right) \frac{(\bar{E})^2}{A^2} \frac{D}{2A} dx \Big|_{x=1} \\ &= 0 + \exp \left(\frac{\bar{E}}{A} \mathcal{Z}_t(1) \right) \frac{D(\bar{E})^2}{4A^3} dx \Big|_{x=1}, \end{aligned}$$

which is exactly the shift to Robin boundary conditions that we see in Proposition 3.13. At the left boundary the Robin condition has a negative sign, which means that here we rather have to apply the backward Itô’s formula to get the correct transformation of the boundary conditions. See Remark 6.15 for the precise point where the correction shows up in our computations.

¹In the notation of Section 6.3 we have $\Theta_u^\varepsilon(v) = \langle \Theta_u, p_\varepsilon^{\text{Dir}}(v, \cdot) \rangle$.

3.6. Statement of the convergence theorems. From now and for the rest of the paper we assume $\rho = \frac{1}{2}$. We are now ready to state all the convergence results for the different fields that we defined previously in Section 2.4:

Theorem 3.17 (Convergence of the density field). *Fix $T > 0$, $k > \frac{5}{2}$ and $\rho = \frac{1}{2}$.*

Then, the sequence of processes $\{\mathcal{Y}_t^n ; t \in [0, T]\}_{n \in \mathbb{N}}$ converges in distribution in $\mathcal{D}([0, T], \mathcal{H}_{\text{Dir}}^{-k})$ as $n \rightarrow \infty$ towards:

- (1) *if $\gamma > \frac{1}{2}$, the Ornstein-Uhlenbeck process solution of (3.1) with Dirichlet boundary conditions (as defined in Proposition 3.1) with $A = 1$ and $D = \frac{1}{2}$;*
- (2) *if $\gamma = \frac{1}{2}$, the unique stationary energy solution of the stochastic Burgers equation (3.4) with Dirichlet boundary conditions (as defined in Theorem 3.3) with $A = 1$, $\bar{E} = E$ and $D = \frac{1}{2}$.*

Remark 3.18. *We note that the previous theorem, when the strength of the asymmetry is taken with $\gamma = 1$, has been already proved in [21] by using a different strategy, namely, considering the microscopic Cole-Hopf transformation – which avoids the derivation of a Boltzmann-Gibbs principle – but in a wide scenario since there the initial measure is quite general.*

Remark 3.19. *It is possible to weaken the assumptions on the initial conditions and to start the dynamics in bounded entropy perturbations of the invariant measure. More precisely, we could assume that \mathcal{Y}_0^n converges weakly in $\mathcal{H}_{\text{Dir}}^{-k}$ to a random variable \mathcal{Y}_0 and that the relative entropy of $\tilde{\gamma}^n = \text{law}(\eta_0^n)$ with respect to $\nu_{1/2}^n$ is (asymptotically) bounded in n :*

$$\limsup_{n \rightarrow \infty} H(\tilde{\gamma}^n | \nu_{1/2}^n) := \limsup_{n \rightarrow \infty} \int \log \frac{d\tilde{\gamma}^n}{d\nu_{1/2}^n}(\eta) d\tilde{\gamma}^n(\eta) < \infty.$$

Indeed, the methods of [19] (essentially the entropy inequality) apply here as well, and under this assumption they allow us to show the convergence of the sequence of processes $\{\mathcal{Y}_t^n ; t \in [0, T]\}_{n \in \mathbb{N}}$ to a non-stationary energy solution, which however has finite entropy with respect to a stationary energy solution (as a measure on the Borel sets of $\mathcal{C}([0, T], \mathcal{H}_{\text{Dir}}^{-k})$). The uniqueness for such bounded entropy perturbations is established in [26] on the real line, but the proof in that paper easily carries over to our situation.

Theorem 3.20 (Convergence of the height field). *Fix $T > 0$, $k > \frac{5}{2}$ and $\rho = \frac{1}{2}$.*

Then, the sequence of processes $\{\mathcal{Z}_t^n ; t \in [0, T]\}_{n \in \mathbb{N}}$ converges in distribution in $\mathcal{D}([0, T], \mathcal{H}_{\text{Neu}}^{-k})$ as $n \rightarrow \infty$ towards:

- (1) *if $\gamma > \frac{1}{2}$, the unique almost stationary energy solution of the KPZ equation (3.8) with Neumann boundary conditions (as defined in Theorem 3.7) with $A = 1$, $\bar{E} = 0$ and $D = \frac{1}{2}$;*
- (2) *if $\gamma = \frac{1}{2}$, the unique almost stationary energy solution of the KPZ equation (3.8) with Neumann boundary conditions (as defined in Theorem 3.7) with $A = 1$, $\bar{E} = E$ and $D = \frac{1}{2}$.*

The strategy of the proof of Theorem 3.17 is quite well known and has been widely used in the past literature. Let us recall here the main steps:

- (1) first, prove that the sequence of probability measures $\{\mathcal{Q}_n\}_{n \in \mathbb{N}}$, where \mathcal{Q}_n is induced by the density fluctuation field \mathcal{Y}^n and the initial measure ν_ρ^n , is tight.

Note that \mathcal{Q}_n is a measure on the Skorokhod space $\mathcal{D}([0, T], \mathcal{H}_{\text{Dir}}^{-k})$. This is the purpose of Section 4.3 below;

- (2) second, write down the approximate martingale problem satisfied by \mathcal{Y}_t^n in the large n limit, and prove that it coincides with the martingale characterization of the solutions of the SPDEs given in Theorem 3.17. The closure of the martingale problem is explained in Section 4.1. In the case $\gamma = \frac{1}{2}$, we need to prove an additional important tool, the so-called *second order Boltzmann-Gibbs principle*, which is stated in Section 4.2 below and proved in Section 5.

The strategy of the proof of Theorem 3.20 is completely similar to the one described above.

Remark 3.21. *Most of our arguments work for any $\rho \in (0, 1)$. However, the restriction $\rho = \frac{1}{2}$ is not just for convenience of notation: otherwise we would pick up an additional diverging transport term in the martingale decomposition for $\mathcal{Y}^n(\varphi)$, roughly speaking $E(1 - 2\rho)n^{3/2-\gamma}\mathcal{Y}^n(\nabla\varphi)$. In the periodic case or if the underlying lattice is \mathbb{Z} we can kill that term by observing our system in a moving frame, see [17] for instance, but of course this does not work in finite volume with boundary conditions. Therefore, we need to assume either $\gamma \geq \frac{3}{2}$ or $\rho = \frac{1}{2}$. Since we are mostly interested in the case $\gamma = \frac{1}{2}$, we take $\rho = \frac{1}{2}$.*

From now on and up to the end of the paper we will mainly assume $\gamma = \frac{1}{2}$ but we will point out the differences with respect to the case $\gamma > \frac{1}{2}$. We also essentially focus on the convergence of \mathcal{Y}^n , since the convergence of \mathcal{Z}^n will follow by very similar arguments.

4. PROOF OF THEOREM 3.17 AND THEOREM 3.20

We start by giving all the details on the proof of Theorem 3.17, and at the end of this section we present the only necessary steps which need to be adapted for the proof of Theorem 3.20. They mainly concern the control of boundary terms for the height fluctuation field.

The section is split in the following way. In Section 4.1 we write down the martingale decomposition which is associated to the density fluctuation field. In Section 4.2 we state the second order Boltzmann-Gibbs principle, whose proof will be given later in Section 5. This principle is needed to control the term in the martingale decomposition which gives rise in the regime $\gamma = \frac{1}{2}$ to the Burgers non-linearity. In Section 4.3 we prove tightness of the density fluctuation field \mathcal{Y}^n , and finally in Section 4.4 we characterize the limit point as a solution to the corresponding SPDE. In Section 4.5 we give the martingale decomposition for the field \mathcal{Z}^n , and we present the estimate that is needed in order to control extra terms that appear at the boundary.

For the sake of clarity from now on we denote $\eta_{tn^2} := \eta_{tn^2}^n$.

4.1. Martingale decomposition for the density fluctuation field \mathcal{Y} . Fix a test function $\varphi \in \mathcal{S}_{\text{Dir}}$ so that $\varphi(0) = \varphi(1) = 0$. By Dynkin's formula, we know that

$$\mathcal{M}_t^n(\varphi) = \mathcal{Y}_t^n(\varphi) - \mathcal{Y}_0^n(\varphi) - \int_0^t n^2 \mathcal{L}_n \mathcal{Y}_s^n(\varphi) ds \quad (4.1)$$

and

$$(\mathcal{M}_t^n(\varphi))^2 - \int_0^t n^2 \mathcal{L}_n (\mathcal{Y}_s^n(\varphi)^2) - 2\mathcal{Y}_s^n(\varphi) n^2 \mathcal{L}_n \mathcal{Y}_s^n(\varphi) ds \quad (4.2)$$

are martingales. The computations from Appendix A.1 show that the integral term in the first martingale (4.1) rewrites as

$$\int_0^t n^2 \mathcal{L}_n \mathcal{Y}_s^n(\varphi) ds = \left(1 + \frac{E}{2n^\gamma}\right) \mathcal{J}_t^n(\varphi) + E \mathcal{A}_t^n(\varphi), \quad (4.3)$$

where

$$\mathcal{J}_t^n(\varphi) = \int_0^t \mathcal{Y}_s^n(\Delta_n \varphi) ds, \quad (4.4)$$

$$\mathcal{A}_t^n(\varphi) = -\frac{\sqrt{n}}{n^\gamma} \int_0^t \sum_{x=1}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \bar{\eta}_{sn^2}(x) \bar{\eta}_{sn^2}(x+1) ds, \quad (4.5)$$

and where, above, $\nabla_n^+ \varphi$ and $\Delta_n \varphi$ are the two functions that approximate on the discrete line the gradient and Laplacian of φ , respectively. They are defined for $x \in \Lambda_n$ by:

$$\begin{aligned} \nabla_n^+ \varphi\left(\frac{x}{n}\right) &= n\left(\varphi\left(\frac{x+1}{n}\right) - \varphi\left(\frac{x}{n}\right)\right), \\ \Delta_n \varphi\left(\frac{x}{n}\right) &= n^2\left(\varphi\left(\frac{x+1}{n}\right) - 2\varphi\left(\frac{x}{n}\right) + \varphi\left(\frac{x-1}{n}\right)\right). \end{aligned}$$

Moreover, in (4.5) we have used a short notation for the centered variable defined as: $\bar{\eta}(x) = \eta(x) - \rho$ for any $x \in \Lambda_n$. From (4.3) we get the identity

$$\mathcal{M}_t^n(\varphi) = \mathcal{Y}_t^n(\varphi) - \mathcal{Y}_0^n(\varphi) - \left(1 + \frac{E}{2n^\gamma}\right) \mathcal{J}_t^n(\varphi) - E \mathcal{A}_t^n(\varphi). \quad (4.6)$$

It is quite easy to see that in the macroscopic limit, the integral term \mathcal{J}_t^n shall correspond to the diffusive macroscopic term $\Delta_{\text{Dir}} \mathcal{Y}_t$. Moreover, when $\gamma = \frac{1}{2}$, \mathcal{A}_t^n shall give rise to the non-linear term in the stochastic Burgers equation, as explained in the next lines, and it will disappear when $\gamma > \frac{1}{2}$.

We also note that since $\varphi \in \mathcal{S}_{\text{Dir}}$, a simple computation shows that the integral term in the second martingale (4.2) rewrites as

$$\begin{aligned} & \int_0^t \frac{1}{n} \sum_{x=1}^{n-2} (\eta_{sn^2}(x) - \eta_{sn^2}(x+1))^2 \left(\nabla_n \varphi\left(\frac{x}{n}\right)\right)^2 ds \\ & + \frac{E}{n^\gamma} \int_0^t \frac{1}{n} \sum_{x=1}^{n-2} \eta_{sn^2}(x) (1 - \eta_{sn^2}(x+1)) \left(\nabla_n \varphi\left(\frac{x}{n}\right)\right)^2 ds \\ & + \int_0^t \frac{1}{n} r_{0,1}(\eta_{sn^2}) \left(\nabla_n \varphi\left(\frac{0}{n}\right)\right)^2 + \frac{1}{n} r_{n-1,n}(\eta_{sn^2}) \left(\nabla_n \varphi\left(\frac{n-1}{n}\right)\right)^2 ds. \end{aligned} \quad (4.7)$$

4.2. Case $\gamma = \frac{1}{2}$: second order Boltzmann-Gibbs principle. In this section we state another important result of this work, which is essential to the proof of Theorem 3.17, since we will be able to treat the term $\mathcal{A}_t^n(\varphi)$ given in (4.5). We focus on the case $\gamma = \frac{1}{2}$, but ahead we make some comments on the case $\gamma > \frac{1}{2}$. Before proceeding, we need to introduce some notations.

Definition 4.1. For any $x \in \Lambda_n$ and $\ell_1 \in \mathbb{N}$ that satisfy $x + \ell_1 \in \Lambda_n$ (resp. $\ell_2 \in \mathbb{N}$ that satisfy $x - \ell_2 \in \Lambda_n$), we denote by $\overrightarrow{\eta}^{\ell_1}(x)$ (resp. $\overleftarrow{\eta}^{\ell_2}(x)$) the average centered configuration on a box of size ℓ_1 (resp. ℓ_2) situated to the right (resp. left) of the site

$x \in \Lambda_n$:

$$\overrightarrow{\eta}^{\ell_1}(x) = \frac{1}{\ell_1} \sum_{z=x+1}^{x+\ell_1} \overline{\eta}(z) \quad (\text{resp.} \quad \overleftarrow{\eta}^{\ell_2}(x) = \frac{1}{\ell_2} \sum_{z=x-\ell_2}^{x-1} \overline{\eta}(z)).$$

For any measurable function $v : \Lambda_n \rightarrow \mathbb{R}$, let us define $\|v\|_{2,n}^2 = \frac{1}{n} \sum_{x \in \Lambda_n} v^2(x)$. From now on and up to the end, $C > 0$ is a constant that does not depend on $t > 0$, nor on $n, \ell \in \mathbb{N}$, and that may change from line to line. Finally, we denote the *static compressibility of the system* by $\chi(\rho) = \int (\eta(1) - \rho)^2 v_\rho^n(d\eta) = \rho(1 - \rho)$.

Theorem 4.2 (Second order Boltzmann-Gibbs principle). *Let $v : \Lambda_n \rightarrow \mathbb{R}$ be a measurable function. There exists $C > 0$ such that, for any $n, \ell \in \mathbb{N}$ such that $\ell < \frac{n}{4}$ and any $t > 0$,*

$$\mathbb{E}_\rho \left[\left(\int_0^t \sum_{x=1}^{n-2} v(x) [\overline{\eta}_{sn^2}(x) \overline{\eta}_{sn^2}(x+1) - \mathcal{Q}(x, \ell, \eta_{sn^2})] ds \right)^2 \right] \leq C t \left(\frac{\ell}{n} + \frac{tn}{\ell^2} \right) \|v\|_{2,n}^2, \quad (4.8)$$

where

$$\mathcal{Q}(x, \ell, \eta) = \begin{cases} (\overrightarrow{\eta}^\ell(x))^2 - \frac{\chi(\rho)}{\ell}, & \text{if } x \in \{1, \dots, n-2\ell-1\}; \\ (\overleftarrow{\eta}^\ell(x))^2 - \frac{\chi(\rho)}{\ell}, & \text{if } x \in \{n-2\ell, \dots, n-2\}. \end{cases} \quad (4.9)$$

Remark 4.3. Notice that the assumption $\ell < \frac{n}{4}$ ensures that one of the two conditions in (4.9) is always satisfied and $\overrightarrow{\eta}^\ell(x)$ and $\overleftarrow{\eta}^\ell(x)$ are always well defined. Moreover, the function $\mathcal{Q}(x, \ell, \cdot)$ is centered, namely $\int \mathcal{Q}(x, \ell, \eta) v_\rho^n(d\eta) = 0$.

Proof of Theorem 4.2. The proof is quite involved and is postponed to Section 5. \square

Now, let us apply Theorem 4.2 when $\ell = \varepsilon n$ (which actually means $\lfloor \varepsilon n \rfloor$ with some abuse of notation): this choice makes the right hand side of (4.8) vanish when we let first $n \rightarrow \infty$ and then $\varepsilon \rightarrow 0$.

As a consequence, when $\gamma = \frac{1}{2}$, $\mathcal{A}_t^n(\varphi)$ is well approximated in $L^2(\mathbb{P}_\rho)$ by the time integral of the following quantity:

$$-E \sum_{x=1}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \mathcal{Q}(x, \varepsilon n, \eta_{sn^2}) \quad (4.10)$$

$$\begin{aligned} &= -E \sum_{x=1}^{n-1-2\varepsilon n} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \left[\frac{1}{\varepsilon n} \sum_{y=x+1}^{x+\varepsilon n} \overline{\eta}_{sn^2}(y) \right]^2 \\ &\quad - E \sum_{x=n-2\varepsilon n}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \left[\frac{1}{\varepsilon n} \sum_{y=x-\varepsilon n}^{x-1} \overline{\eta}_{sn^2}(y) \right]^2 + E \frac{\chi(\rho)}{\varepsilon n} n \left[\varphi\left(\frac{n-1}{n}\right) - \varphi\left(\frac{1}{n}\right) \right] \\ &= -\frac{E}{n} \sum_{x=1}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \left[\frac{1}{\sqrt{n}} \sum_{y=1}^{n-1} \overline{\eta}_{sn^2}(y) \iota_\varepsilon\left(\frac{x}{n}\right)\left(\frac{y}{n}\right) \right]^2 + E \frac{\chi(\rho)}{\varepsilon} \left[\varphi\left(\frac{n-1}{n}\right) - \varphi\left(\frac{1}{n}\right) \right] \\ &= -\frac{E}{n} \sum_{x=1}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \left[\mathcal{Q}_s^n\left(\iota_\varepsilon\left(\frac{x}{n}\right)\right) \right]^2 + E \frac{\chi(\rho)}{\varepsilon} \left[\varphi\left(\frac{n-1}{n}\right) - \varphi\left(\frac{1}{n}\right) \right]. \end{aligned} \quad (4.11)$$

Since $\varphi(0) = \varphi(1) = 0$, the second term of (4.11) is of order $1/n$ and therefore vanishes when we let $n \rightarrow \infty$. We also note that for the case $\gamma > \frac{1}{2}$ the previous term (4.10)

has a factor \sqrt{n}/n^γ in front of it and for that reason it vanishes as $n \rightarrow \infty$. Finally, the computation above (and more precisely the first term of (4.11)) motivates the definition of energy solutions as defined in Theorem 3.3, that is the definition of the macroscopic field \mathcal{A}_t as given in (3.7). Indeed, putting all these considerations together in the case $\gamma = \frac{1}{2}$, we see that (4.6) rewrites as

$$\begin{aligned} \mathcal{M}_t^n(\varphi) &= \mathcal{Y}_t^n(\varphi) - \mathcal{Y}_0^n(\varphi) - \int_0^t \mathcal{Y}_s^n(\Delta_n \varphi) ds \\ &\quad - E \int_0^t \frac{1}{n} \sum_{x=1}^{n-2} \nabla_n^+ \varphi\left(\frac{x}{n}\right) \left[\mathcal{Y}_s^n\left(\iota_\varepsilon\left(\frac{x}{n}\right)\right) \right]^2 ds + o_t^n(1), \end{aligned} \quad (4.12)$$

where $o_t^n(1)$ is deterministic and satisfies $\sup_{t \in [0, T]} |o_t^n(1)| \rightarrow 0$ as $n \rightarrow \infty$. This computation will be useful to characterize limit points of the density fluctuation field (see Section 4.4 below). Before that, let us prove tightness.

4.3. Tightness of the density fluctuation field. In this section, for the sake of completeness we show tightness of the sequence $\{\mathcal{Y}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}$, following closely [31]. The main difference is the presence of the extra term \mathcal{A}_t^n in the martingale decomposition. Tightness is a consequence of the following lemma.

Lemma 4.4. *For $k > \frac{5}{2}$,*

$$\begin{aligned} (1) \quad & \lim_{A \rightarrow \infty} \limsup_{n \rightarrow \infty} \mathbb{P}_\rho \left[\sup_{t \in [0, T]} \|\mathcal{Y}_t^n\|_{-k} > A \right] = 0 \\ (2) \quad & \lim_{\delta \rightarrow 0} \limsup_{n \rightarrow \infty} \mathbb{P}_\rho \left[\omega_\delta(\mathcal{Y}^n) \geq \varepsilon \right] = 0 \end{aligned}$$

for every $\varepsilon > 0$, where for $\delta > 0$ we define

$$\omega_\delta(\mathcal{Y}^n) = \sup_{\substack{|s-t| < \delta \\ s, t \in [0, T]}} \|\mathcal{Y}_t^n - \mathcal{Y}_s^n\|_{-k}$$

and $\|\cdot\|_{-k}$ has been defined in (2.6).

Proof of Lemma 4.4. We split the proof of this lemma into two steps. To prove (1), from Markov's inequality it is enough to notice that

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} \|\mathcal{Y}_t^n\|_{-k}^2 \right] \leq \sum_{m \geq 1} (m\pi)^{-2k} \mathbb{E}_\rho \left[\sup_{t \in [0, T]} (\mathcal{Y}_t^n(e_m))^2 \right]. \quad (4.13)$$

Now we compute the expectation at the right hand side of (4.13) using the martingale decomposition (4.6) for $\varphi = e_m$ which makes sense since $e_m \in \mathcal{S}_{\text{Dir}}$. First, we have by independence under ν_ρ^n ,

$$\mathbb{E}_\rho \left[|\mathcal{Y}_0^n(e_m)|^2 \right] \leq \frac{1}{4n} \sum_{x=1}^{n-1} \left(e_m\left(\frac{x}{n}\right) \right)^2 \leq \frac{1}{2}. \quad (4.14)$$

Now, from the Cauchy-Schwarz inequality, invariance of ν_ρ^n and independence, we get that

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} |\mathcal{A}_t^n(e_m)|^2 \right] \leq \frac{T^2}{4n} \sum_{x=1}^{n-1} \left(\Delta_n e_m\left(\frac{x}{n}\right) \right)^2 \leq \frac{\pi^4 T^2}{2} m^4. \quad (4.15)$$

The martingale term can be easily estimated by Doob's inequality as

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} |\mathcal{M}_t^n(e_m)|^2 \right] \lesssim \mathbb{E}_\rho \left[\int_0^T \frac{1}{n} \sum_{x=1}^{n-1} \left(\nabla_n^+ e_m \left(\frac{x}{n} \right) \right)^2 ds \right] \lesssim T m^2.$$

Finally, in order to estimate the remaining term $\mathcal{A}_t^n(e_m)$, we sum and subtract to $\bar{\eta}(x)\bar{\eta}(x+1)$ the term $\mathcal{Q}(x, \ell, \eta)$, and from the elementary inequality $(x+y)^2 \leq 2x^2 + 2y^2$, it remains to bound the following two expectations:

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} \left(\int_0^t \sum_{x=1}^{n-2} \nabla_n^+ e_m \left(\frac{x}{n} \right) (\bar{\eta}_{sn^2}(x) \bar{\eta}_{sn^2}(x+1) - \mathcal{Q}(x, \ell, \eta_{sn^2})) ds \right)^2 \right] \quad (4.16)$$

and

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} \left(\int_0^t \sum_{x=1}^{n-2} \nabla_n^+ e_m \left(\frac{x}{n} \right) \mathcal{Q}(x, \ell, \eta_{sn^2}) ds \right)^2 \right]. \quad (4.17)$$

From the Cauchy-Schwarz inequality, the expectation (4.17) is bounded from above by $T^2 m^2 n / \ell$ and by choosing $\ell = Cn$, we can bound it from above by $T^2 m^2$. The remaining expectation (4.16) can be bounded by $C(T)m^2$, from [3, Lemma 4.3] and following the same steps as in the proof of Theorem 4.2 (given in Section 5). Collecting all the previous computations we get that (4.13) is bounded from above by $C(T) \sum_{m \geq 1} m^{4-2k}$, which is finite as long as $2k - 4 > 1$.

Now we prove (2). For that purpose, at first we notice that from the previous computations we have: for $k > \frac{5}{2}$

$$\lim_{j \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \mathbb{E}_\rho \left[\sup_{t \in [0, T]} \sum_{|m| \geq j} \left(\mathcal{Y}_t^n(e_m) \right)^2 (m\pi)^{-2k} \right] = 0.$$

Therefore, to conclude the proof we just have to show that for any $j \in \mathbb{N}$ and $\varepsilon > 0$,

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\sup_{\substack{|s-t| < \delta \\ s, t \in [0, T]}} \sum_{|m| \leq j} \left(\mathcal{Y}_t^n(e_m) - \mathcal{Y}_s^n(e_m) \right)^2 (m\pi)^{-2k} > \varepsilon \right] = 0. \quad (4.18)$$

In fact we prove that for every $m \geq 1$ and $\varepsilon > 0$

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\sup_{\substack{|s-t| < \delta \\ s, t \in [0, T]}} \left| \mathcal{Y}_t^n(e_m) - \mathcal{Y}_s^n(e_m) \right| > \varepsilon \right] = 0,$$

from which (4.18) follows. Now, as before, we recall (4.6) so that the previous result is accomplished if we derive the same result for each term in the martingale decomposition of $\mathcal{Y}_t^n(e_m)$. We start by the most demanding one, which is the term that involves the martingales, more precisely, we show that for every $m \geq 1$ and $\varepsilon > 0$

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\sup_{\substack{|s-t| < \delta \\ s, t \in [0, T]}} \left| \mathcal{M}_t^n(e_m) - \mathcal{M}_s^n(e_m) \right| > \varepsilon \right] = 0. \quad (4.19)$$

Since it is easy to see that

$$\limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\sup_{t \in [0, T]} \left| \mathcal{M}_t^n(e_m) - \mathcal{M}_{t-}^n(e_m) \right| > \varepsilon \right] = 0,$$

the claim (4.19) becomes a consequence of the following fact:

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\omega'_\delta(\mathcal{M}_t^n(e_m)) > \varepsilon \right] = 0,$$

where $\omega'_\delta(\mathcal{M}_t^n(e_m))$ is the modified modulus of continuity defined by

$$\omega'_\delta(\mathcal{M}_t^n(e_m)) = \inf_{\{t_i\}} \max_{0 \leq i \leq r} \sup_{t_i \leq s < t \leq t_{i+1}} |\mathcal{M}_t^n(e_m) - \mathcal{M}_s^n(e_m)|,$$

the infimum being taken over all partitions of $[0, T]$ such that $0 = t_0 < t_1 < \dots < t_r = T$ with $t_{i+1} - t_i > \delta$. By the Aldous criterion, see for example [31, Proposition 4.1.6], it is enough to show that:

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \sup_{\tau \in \mathfrak{T}_T} \mathbb{P}_\rho \left[|\mathcal{M}_{\tau+\tilde{\tau}}^n(e_m) - \mathcal{M}_\tau^n(e_m)| > \varepsilon \right] = 0$$

for every $\varepsilon > 0$, where \mathfrak{T}_T denotes the family of all stopping times bounded by T . Using Tchebychev's inequality together with the optional stopping Theorem, the last probability is bounded from above by

$$\varepsilon^{-2} \mathbb{E}_\rho \left[(\mathcal{M}_{\tau+\tilde{\tau}}^n(e_m))^2 - (\mathcal{M}_\tau^n(e_m))^2 \right],$$

which, by definition of the quadratic variation of the martingale can be bounded from above by $\varepsilon^{-2} C m^2 \delta$, and vanishes as $\delta \rightarrow 0$.

Now, we compute the remaining term that involves $\mathcal{J}_t^n(e_m)$. We have to show that for every $\varepsilon > 0$

$$\lim_{\delta \rightarrow 0} \limsup_{n \rightarrow +\infty} \mathbb{P}_\rho \left[\sup_{\substack{|s-t| < \delta \\ s, t \in [0, T]}} \left| \int_s^t \mathcal{J}_r^n(\Delta_n e_m) dr \right| > \varepsilon \right] = 0$$

By Tchebychev's inequality and Cauchy-Schwarz's inequality, we obtain that the last probability is bounded by $T \delta m^4 / \varepsilon^2$, which vanishes as $\delta \rightarrow 0$. For the last term involving $\mathcal{A}_t^n(e_m)$ we can repeat the computations that we did above: we sum and subtract the term $\mathcal{Q}(x, \ell, \eta)$ as in (4.16) and (4.17), then we chose $\ell = Cn$, $C > 0$, and we prove that each contribution is of order $T \delta m^2 / \varepsilon^2$, and therefore it goes to zero as $\delta \rightarrow 0$, which finishes the proof. \square

4.4. Characterization of limiting points. In this section we prove that any limit point of the tight sequence $\{\mathcal{Y}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}$ concentrates on stationary energy solutions of (3.4) as defined in Theorem 3.3. Up to extraction, one can consider that the four sequences

$$\{\mathcal{Y}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}; \{\mathcal{M}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}; \{\mathcal{J}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}; \{\mathcal{A}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}$$

converge as $n \rightarrow \infty$ to

$$\{\mathcal{Y}_t; t \in [0, T]\}; \{\mathcal{M}_t; t \in [0, T]\}; \{\mathcal{J}_t; t \in [0, T]\}; \{\mathcal{A}_t; t \in [0, T]\}$$

respectively. First, one can repeat the argument taken from [17, Section 5.3] to prove that the limit point $\{\mathcal{Y}_t; t \in [0, T]\}$ has continuous trajectories and it is stationary in the sense of Theorem 3.3 (see item (1)). The characterization will be complete if we prove that this limit process also satisfies the remaining three items of Theorem 3.3. This is what we explain briefly in the next paragraphs, since the argument is now standard and is given for example in [10, 17, 19].

4.4.1. *Proof of item (2).* We give a few proof elements for the sake of completeness. First note that $n\mathcal{Q}(x, \varepsilon n, \eta_{sn^2}) = \mathcal{Y}_s^n(\iota_\varepsilon(\frac{x}{n}))^2 - \frac{\chi(\rho)}{\varepsilon}$. Since \mathcal{Y}^n converges to \mathcal{Y} and since we can approximate the function $\iota_\varepsilon(\frac{x}{n})$ by proper functions in the space \mathcal{S}_{Dir} , we get

$$\mathcal{A}_{s,t}^\varepsilon(\varphi) = \lim_{n \rightarrow \infty} \int_s^t \sum_{x=1}^{n-2} \nabla_n \varphi(\frac{x}{n}) \mathcal{Q}(x, \varepsilon n, \eta_{rn^2}) dr.$$

Now, rewriting the Boltzmann-Gibbs principle stated in Theorem 4.2, in terms of $\mathcal{A}_t^n(\varphi)$, we see that there exists $C > 0$ such that, for any $\varphi \in \mathcal{S}_{\text{Dir}}$,

$$\mathbb{E}_\rho \left[\left(\mathcal{A}_t^n(\varphi) - \mathcal{A}_s^n(\varphi) - \int_s^t \sum_{x=1}^{n-2} \nabla_n \varphi(\frac{x}{n}) \mathcal{Q}(x, \varepsilon n, \eta_{rn^2}) dr \right)^2 \right] \leq C(t-s)\varepsilon \|\nabla \varphi\|_{2,n}^2. \quad (4.20)$$

The last claim is proved as follows: in $\mathcal{A}_t^n(\varphi)$ given in (4.5) we sum and subtract $\nabla_n \varphi(\frac{x}{n}) \mathcal{Q}(x, \varepsilon n, \eta_{sn^2})$ inside the sum, and we use a standard convexity inequality in order to treat two terms separately. The first one is handled using the Boltzmann-Gibbs principle, the second one is estimated thanks to the computation (4.11). Then, since L^2 -bounds are stable under the convergence in law, from the previous estimate we conclude that

$$\mathbb{E}_\rho \left[\left(\mathcal{A}_t(\varphi) - \mathcal{A}_s(\varphi) - \mathcal{A}_{s,t}^\varepsilon(\varphi) \right)^2 \right] \leq C(t-s)\varepsilon \|\nabla \varphi\|_{L^2([0,1])}^2. \quad (4.21)$$

Finally, the energy estimate (3.5) is a trivial consequence of (4.21), since it follows from adding and subtracting the quantity $(\mathcal{A}_t(\varphi) - \mathcal{A}_s(\varphi))$ inside the square.

4.4.2. *Proof of item (3).* This point is now a straightforward consequence of the martingale decomposition given in (4.12) and in (4.7), in which one can pass to the limit $n \rightarrow \infty$, together with the previous paragraph.

4.4.3. *Proof of item (4).* This last property can be obtained easily by considering the reversed dynamics with the adjoint of the infinitesimal generator \mathcal{L}_n^* with respect to the Bernoulli product measure ν_ρ^n and repeating the same exact arguments as we did above.

4.5. **Sketch of the proof of Theorem 3.20.** As mentioned previously, the proof of Theorem 3.20 is essentially the same as for \mathcal{Y}^n . Let us give here some hints and follow the sketch from the previous paragraphs.

First, let us note that the analogue of the martingale decomposition (4.3) contains also one boundary term. Indeed, fix a test function $\varphi \in \mathcal{S}_{\text{Neu}}$, and let $n^2 \mathcal{L}_n^\otimes$ denote the generator of the joint process $\{\{\eta_{tn^2}(x), h_{tn^2}^n(1)\}_{x \in \Lambda_n} ; t \geq 0\}$. This generator acts on functions $f : \Omega_n \times \mathbb{Z} \rightarrow \mathbb{R}$ as follows:

$$\begin{aligned} \mathcal{L}_n^\otimes f(\eta, h) &= \sum_{x=1}^{n-2} r_{x,x+1}(\eta) \left(f(\sigma^{x,x+1} \eta, h) - f(\eta, h) \right) \\ &\quad + \eta(1)(1-\rho) \left(f(\sigma^1 \eta, h+1) - f(\eta, h) \right) \\ &\quad + \left(1 + \frac{E}{n^\gamma} \right) \rho(1-\eta(1)) \left(f(\sigma^1 \eta, h-1) - f(\eta, h) \right) \\ &\quad + \left\{ \left(1 + \frac{E}{n^\gamma} \right) (1-\rho)\eta(n-1) + \rho(1-\eta(n-1)) \right\} \left(f(\sigma^{n-1} \eta, h) - f(\eta, h) \right) \end{aligned} \quad (4.22)$$

where $r_{x,x+1}$ has been already defined in (2.1). From the computations in Appendix A.2 we get that

$$\mathcal{N}_t^n(\varphi) = \mathcal{X}_t^n(\varphi) - \mathcal{X}_0^n(\varphi) - \left(1 + \frac{E}{2n^\gamma}\right) \int_0^t \mathcal{X}_s^n(\tilde{\Delta}_n \varphi) ds - E \mathcal{B}_t^n(\varphi) + \mathcal{R}_t^n(\varphi) + o_t^n(1), \quad (4.23)$$

is a martingale. Above the terms \mathcal{R}_t^n and \mathcal{B}_t^n are given by:

$$\begin{aligned} \mathcal{R}_t^n(\varphi) &= \sqrt{n} \int_0^t \left[\left(\varphi\left(\frac{1}{n}\right) - \varphi(0) \right) (h_{sn^2}^n(1) - c_n s) \right. \\ &\quad \left. + \left(\varphi(1) - \varphi\left(\frac{n-1}{n}\right) \right) (h_{sn^2}^n(n) - c_n s) \right] ds \\ \mathcal{B}_t^n(\varphi) &= \frac{\sqrt{n}}{n^\gamma} \int_0^t \sum_{x=2}^{n-1} \varphi\left(\frac{x}{n}\right) \nabla^- h_{sn^2}^n(x) \nabla^+ h_{sn^2}^n(x) ds, \end{aligned} \quad (4.24)$$

where we set $\nabla^- h(x) = h(x) - h(x-1) = \bar{\eta}(x-1)$ and $\nabla^+ h(x) = h(x+1) - h(x) = \bar{\eta}(x)$. Moreover, in (4.23), $o_t^n(1)$ is a deterministic sequence of real numbers that vanishes as $n \rightarrow \infty$, uniformly in $t \in [0, T]$, and we also have used the notation $\tilde{\Delta}_n \varphi$ to denote the following approximation of the Laplacian:

$$\tilde{\Delta}_n \varphi\left(\frac{x}{n}\right) = \begin{cases} \Delta_n \varphi\left(\frac{x}{n}\right) & \text{if } x \in \{1, \dots, n-1\}, \\ 2n^2 \left(\varphi\left(\frac{n-1}{n}\right) - \varphi\left(\frac{n}{n}\right) \right) & \text{if } x = n. \end{cases} \quad (4.25)$$

Note that, if $\varphi \in \mathcal{S}_{\text{Neu}}$ and therefore satisfies $\nabla \varphi(1) = 0$, then $\tilde{\Delta}_n \varphi$ is indeed an approximation of the usual Laplacian as $n \rightarrow \infty$.

Let us start with \mathcal{B}_t^n . Note that we have for any test function φ ,

$$\mathcal{B}_t^n(\nabla_n^- \varphi) = -\mathcal{A}_t^n(\varphi),$$

where ∇_n^- is defined similarly to ∇_n^+ except that the discrete gradient is shifted, namely: $\nabla_n^- \varphi\left(\frac{x}{n}\right) = n(\varphi\left(\frac{x}{n}\right) - \varphi\left(\frac{x-1}{n}\right))$. As a consequence, this term can be treated as \mathcal{A}_t^n , using the Boltzmann-Gibbs principle (Theorem 4.2): it gives rise to the KPZ non-linearity as soon as $\gamma = \frac{1}{2}$ and vanishes when $\gamma > \frac{1}{2}$.

Next, in (4.23), the term \mathcal{R}_t^n (which does not depend on γ) comes from boundary effects, but does not contribute to the limit if $\varphi \in \mathcal{S}_{\text{Neu}}$, as a consequence of the following lemma.

Lemma 4.5. *For any $\rho \in (0, 1)$ we have*

$$\lim_{n \rightarrow \infty} \left\{ \mathbb{E}_\rho \left[\sup_{t \in [0, T]} |n^{-3/2} (h_{tn^2}^n(1) - c_n t)|^2 \right] + \mathbb{E}_\rho \left[\sup_{t \in [0, T]} |n^{-3/2} (h_{tn^2}^n(n) - c_n t)|^2 \right] \right\} = 0,$$

and in particular, for any $\varphi \in \mathcal{S}_{\text{Neu}}$, the term $\mathcal{R}_t^n(\varphi)$ defined in (4.24) converges to 0 in $L^2(\mathbb{P}_\rho)$, locally uniformly in time.

Proof of Lemma 4.5. Since $h^n(1)$ increases by 1 whenever a particle leaves the system and decreases by 1 whenever a particle enters, one can easily write

$$\begin{aligned} h_{tn^2}^n(1) &= h_0^n(1) + \int_0^{tn^2} \left\{ \frac{1}{2} \eta_s(1) - \frac{1}{2} \left(1 + \frac{E}{n^\gamma} \right) (1 - \eta_s(1)) \right\} ds + M_{tn^2}^n \\ &= \int_0^{tn^2} \left\{ \left(1 + \frac{E}{2n^\gamma} \right) \bar{\eta}_s(1) + n^{-2} c_n \right\} ds + M_{tn^2}^n, \end{aligned} \quad (4.26)$$

where $c_n = -n^{2-\gamma}E/4$ has been defined in (2.9) and M^n is a martingale with predictable quadratic variation

$$\langle M^n \rangle_t = \frac{1}{2} \left(\left(1 + \frac{E}{n^\gamma} \right) - \frac{E}{n^\gamma} \eta_t^n(1) \right).$$

The Burkholder-Davis-Gundy inequality implies for all $p \geq 1$ the following: there exists $C > 0$ such that

$$\begin{aligned} \mathbb{E}_\rho \left[\sup_{t \in [0, Tn^2]} |n^{-3/2} M_t^n|^p \right] &\lesssim n^{-3p/2} \mathbb{E}_\rho \left[\left\{ \int_0^{Tn^2} \frac{1}{2} \left(\left(1 + \frac{E}{n^\gamma} \right) - \frac{E}{n^\gamma} \eta_t^n(1) \right) dt \right\}^{p/2} \right] \\ &\quad + n^{-3p/2} \mathbb{E}_\rho \left[\left| \sup_{t \in [0, Tn^2]} \Delta_t M^n \right|^p \right], \end{aligned}$$

where $\Delta_t M^n$ is the jump of M^n at time t and therefore bounded by 1. In the integrand we can bound $-\eta_t^n(1)$ from above by 0 and therefore $n^{-3/2} M_{tn^2}^n$ vanishes in the limit. From (4.26) we get

$$n^{-3/2} (h_{tn^2}^n(1) - c_n t) = \int_0^t \sqrt{n} \left(1 + \frac{E}{2n^\gamma} \right) \bar{\eta}_{sn^2}(1) ds + n^{-3/2} M_{tn^2}^n,$$

and therefore we are left with bounding $\sqrt{n} \int_0^t \bar{\eta}_{sn^2}(1) ds$. With the Kipnis-Varadhan Lemma given for instance in [31, Proposition A.1.6.1] we estimate

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} \left| \int_0^t \bar{\eta}_{sn^2}(1) ds \right|^2 \right] \lesssim T \sup_{f \in L^2(\nu_\rho^n)} \left\{ 2 \int \bar{\eta}(1) f(\eta) \nu_\rho^n(d\eta) - n^2 \mathfrak{D}_n(f) \right\}$$

where $\mathfrak{D}_n(f)$ is the Dirichlet form defined in (2.2). From the decomposition (2.3) we easily obtain

$$\mathfrak{D}_n(f) \geq (\rho \wedge (1 - \rho)) \int (f(\sigma^1 \eta) - f(\eta))^2 \nu_\rho^n(d\eta).$$

Moreover, as in [12, Lemma 3] we can use Young's inequality to get

$$\begin{aligned} 2 \int \bar{\eta}(1) f(\eta) \nu_\rho^n(d\eta) &= \int \bar{\eta}(1) (f(\eta) - f(\sigma^1 \eta)) \nu_\rho^n(d\eta) \\ &\leq \varepsilon \int (\bar{\eta}(1))^2 \nu_\rho^n(d\eta) + \frac{1}{4\varepsilon} \int (f(\eta) - f(\sigma^1 \eta))^2 \nu_\rho^n(d\eta), \end{aligned}$$

for any $\varepsilon > 0$, that we choose now such that $(4\varepsilon)^{-1} = (\rho \wedge (1 - \rho)) n^2$. And therefore we obtain

$$2 \int \bar{\eta}(1) f(\eta) \nu_\rho^n(d\eta) - n^2 \mathfrak{D}_n(f) \lesssim \frac{1}{n^2} \int (\bar{\eta}(1))^2 \nu_\rho^n(d\eta) \simeq \frac{1}{n^2},$$

which leads to

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} \left| \sqrt{n} \int_0^t \bar{\eta}_{sn^2}(1) ds \right|^2 \right] \lesssim \frac{T}{n}.$$

The bound for $h^n(n)$ is shown with the same arguments. To conclude the proof, take $\varphi \in \mathcal{S}_{\text{Neu}}$, so that $\nabla \varphi(0) = \nabla \varphi(1) = 0$. In that case

$$n^2 \left(\varphi\left(\frac{1}{n}\right) - \varphi(0) \right) \xrightarrow{n \rightarrow \infty} \Delta \varphi(0), \quad \text{and} \quad n^2 \left(\varphi(1) - \varphi\left(\frac{n-1}{n}\right) \right) \xrightarrow{n \rightarrow \infty} \Delta \varphi(1),$$

and therefore

$$\mathbb{E}_\rho \left[\sup_{t \in [0, T]} |\mathcal{R}_t^n(\varphi)|^2 \right] \lesssim \frac{T}{n}.$$

□

There are two remaining steps, the first one being tightness of the sequence $\{\mathcal{Z}_t^n; t \in [0, T]\}_{n \in \mathbb{N}}$. We let the reader repeat the same proof as in Lemma 4.4, noting that since the height fluctuation field is now defined in $\mathcal{H}_{\text{Neu}}^{-k}$, the basis that one has to use is \tilde{e}_m given in Section 2.3. The arguments remain unchanged, we only note that the restriction $k > \frac{5}{2}$ comes from the analog of (4.15).

Finally, for the characterization of limit points we essentially use the relation between \mathcal{Z}_t^n and \mathcal{Y}_t^n , which reads: for any $\varphi \in \mathcal{S}_{\text{Dir}}$,

$$\mathcal{Z}_t^n(\tilde{\nabla}_n \varphi) = -\mathcal{Y}_t^n(\varphi) - \frac{1}{\sqrt{n}} \varphi\left(\frac{1}{n}\right) (h_{tn^2}^n(1) - c_n t) + \frac{1}{\sqrt{n}} \varphi\left(\frac{n}{n}\right) (h_{tn^2}^n(n) - c_n t), \quad (4.27)$$

where $\tilde{\nabla}_n \varphi\left(\frac{x}{n}\right) = \mathbf{1}_{\{1, \dots, n-1\}}(x) \nabla_n^+ \varphi\left(\frac{x}{n}\right)$, in particular $\tilde{\nabla}_n \varphi\left(\frac{n}{n}\right) = 0$. Since $\varphi \in \mathcal{S}_{\text{Dir}}$, Lemma 4.5 implies that the last two terms in (4.27) vanish in $L^2(\mathbb{P}_\rho)$ as $n \rightarrow \infty$ uniformly in $t \in [0, T]$. If \mathcal{Z} is the limit point of \mathcal{Z}^n then passing to the limit in (4.27), we get: for any $\varphi \in \mathcal{S}_{\text{Dir}}$,

$$\nabla_{\text{Dir}} \mathcal{Z}_t(\varphi) = -\mathcal{Z}_t(\nabla \varphi) = \mathcal{Y}_t(\varphi),$$

where we used the definition of ∇_{Dir} given in (2.7). From this, we deduce item (1) of Theorem 3.7. The last item (3) is obtained similarly combining (4.27) with the two martingale decompositions (4.6) and (4.23). For the quadratic variation we observe that by Dynkin's formula

$$\left(\mathcal{N}_t^n(\varphi)\right)^2 - \int_0^t n^2 \mathcal{L}_n^\otimes(\mathcal{Z}_s^n(\varphi)^2) - 2\mathcal{Z}_s^n(\varphi) n^2 \mathcal{L}_n^\otimes \mathcal{Z}_s^n(\varphi) ds$$

is a martingale. Here we used that the drift $-c_n t$ gives rise to a first order differential operator \mathcal{G}^n which satisfies Leibniz's rule and therefore the difference $\mathcal{G}^n(\mathcal{Z}_s^n(\varphi)^2) - 2\mathcal{Z}_s^n(\varphi) \mathcal{G}^n \mathcal{Z}_s^n(\varphi)$ vanishes. A simple computation shows that last integral term can be rewritten as

$$\begin{aligned} & \int_0^t \frac{1}{n} \sum_{x=2}^{n-2} (\eta_{sn^2}(x) - \eta_{sn^2}(x+1))^2 \left(\varphi\left(\frac{x}{n}\right)\right)^2 ds \\ & + \frac{E}{n^\gamma} \int_0^t \frac{1}{n} \sum_{x=2}^{n-2} \eta_{sn^2}(x) (1 - \eta_{sn^2}(x+1)) \left(\varphi\left(\frac{x}{n}\right)\right)^2 ds \\ & + \int_0^t \frac{1}{n} r_{0,1}(\eta_{sn^2}) \left(\varphi\left(\frac{1}{n}\right)\right)^2 + \frac{1}{n} r_{n-1,n}(\eta_{sn^2}) \left(\varphi\left(\frac{n-1}{n}\right)\right)^2 ds. \end{aligned}$$

By taking expectation w.r.t. ν_ρ^n and sending $n \rightarrow \infty$ we conclude item (3).

5. PROOF OF THE SECOND ORDER BOLTZMANN-GIBBS PRINCIPLE

In this section we prove the Boltzmann-Gibbs principle stated in Theorem 4.2. The proof uses similar arguments as in the original ones [20, 2, 10]. The main difference is the separation of the lattice sites into two sets: $\{1, \dots, n-2\ell-1\}$ and $\{n-2\ell, \dots, n-1\}$, which are treated separately. This technical novelty is necessary in order to take into account the presence of fixed endpoints.

Let us illustrate how the proof of this principle works: let us choose a site x which is not too close to the right boundary, in the sense that there are at least 2ℓ sites between x and $n-1$, then we can replace the local function $\bar{\eta}(x)\bar{\eta}(x+1)$ by the square of the average to its right $(\bar{\eta}^\ell(x))^2$ (see Figure 2). The main reason to keep at least 2ℓ

sites between x and $n-1$ is because the proof makes use of the sites situated between $x+\ell+1$ and $x+2\ell$, as explained in Section 5.1 below. Otherwise, when $x+2\ell > n-1$, we replace the same local function by the square of the average to its left $(\overleftarrow{\eta}^\ell(x))^2$ (see Figure 3).

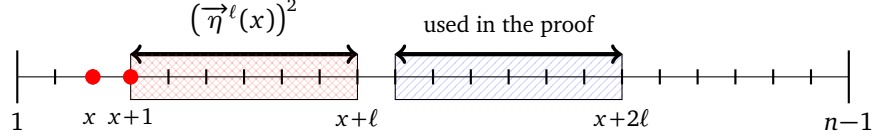


FIGURE 2. Replacement for the local function $\overrightarrow{\eta}(x)\overrightarrow{\eta}(x+1)$ when $x+2\ell \leq n-1$.

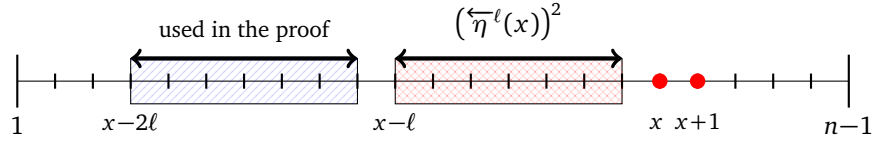


FIGURE 3. Replacement for the local function $\overrightarrow{\eta}(x)\overrightarrow{\eta}(x+1)$ when $x+2\ell > n-1$.

Before going into the proof details, let us introduce some notations: for a function $g : \Omega_n \rightarrow \mathbb{R}$ we denote by $\|g\|_2$ its $L^2(\nu_\rho^n)$ -norm:

$$\|g\|_2^2 = \int_{\Omega_n} g^2(\eta) \nu_\rho^n(d\eta).$$

In the following, $C = C(\rho)$ denotes a constant that does not depend on n nor on t nor on the sizes of all boxes involved, and that may change from line to line. We fix once and for all a measurable function $v : \Lambda_n \rightarrow \mathbb{R}$, for which $\|v\|_{2,n}^2 < \infty$. We say that a function $g : \{0, 1\}^{\mathbb{Z}} \rightarrow \mathbb{R}$ has its support (denoted below by $\text{Supp}(g)$) included in some subset $\Lambda \subset \mathbb{Z}$ if g depends on the configuration η only through the variables $\{\eta(x) ; x \in \Lambda\}$. We denote by τ_x the usual shift operator, that acts on functions $g : \Omega_n \rightarrow \mathbb{R}$ as follows: $\tau_x g(\eta) = g(\tau_x \eta)$, which is well defined for any x such that $\text{Supp}(g) \subset \{1, \dots, n-1-x\}$. To keep the presentation as clear as possible, we define two quantities that are needed in due course:

Definition 5.1. Let $m \in \Lambda_n$ be an integer such that $m < \frac{n}{2}$, and let

- $g^\rightarrow : \{0, 1\}^{\mathbb{Z}} \rightarrow \mathbb{R}$ be a function whose support is included in $\{0, \dots, m\}$,
- $g^\leftarrow : \{0, 1\}^{\mathbb{Z}} \rightarrow \mathbb{R}$ be a function whose support is included in $\{-m, \dots, 1\}$.

Let us define

$$\mathfrak{J}_{t,n}^{\text{left}}(g^\rightarrow) = \mathbb{E}_\rho \left[\left(\int_0^t \sum_{x=1}^{n-2\ell-1} v(x) \tau_x g^\rightarrow(\eta_{sn^2}) ds \right)^2 \right],$$

$$\mathfrak{I}_{t,n}^{\text{right}}(g^{\leftarrow}) = \mathbb{E}_{\rho} \left[\left(\int_0^t \sum_{x=n-2\ell}^{n-2} v(x) \tau_x g^{\leftarrow}(\eta_{sn^2}) ds \right)^2 \right].$$

With this definition, (4.8) follows from showing that for any $n, \ell \in \mathbb{N}$ such that $\ell < \frac{n}{4}$, and any $t > 0$,

$$\mathfrak{I}_{t,n}^{\text{left}}(g_{\ell}^{\rightarrow}) \leq C t \left(\frac{\ell}{n} + \frac{tn}{\ell^2} \right) \|v\|_{2,n}^2 \quad (5.1)$$

$$\mathfrak{I}_{t,n}^{\text{right}}(g_{\ell}^{\leftarrow}) \leq C t \left(\frac{\ell}{n} + \frac{tn}{\ell^2} \right) \|v\|_{2,n}^2 \quad (5.2)$$

where the two local functions g_{ℓ}^{\rightarrow} and g_{ℓ}^{\leftarrow} are given by

$$g_{\ell}^{\rightarrow}(\eta) = \bar{\eta}(0)\bar{\eta}(1) - (\bar{\eta}^{\ell}(0))^2 + \frac{\chi(\rho)}{\ell}, \quad (5.3)$$

$$g_{\ell}^{\leftarrow}(\eta) = \bar{\eta}(0)\bar{\eta}(1) - (\bar{\eta}^{\ell}(0))^2 + \frac{\chi(\rho)}{\ell}. \quad (5.4)$$

5.1. Strategy of the proof. We prove (5.1) and (5.2) separately. For both of them, we need to decompose g_{ℓ}^{\rightarrow} and g_{ℓ}^{\leftarrow} as sums of several local functions, for which the estimates are simpler. With a small abuse of language, we say that, at each step, we *replace* a local function with another one. More precisely, let $\ell_0 \leq \ell$ and assume first (to simplify) that $\ell = 2^M \ell_0$ for some integer $M \in \mathbb{N}$. Denote $\ell_k = 2^k \ell_0$ for $k \in \{0, \dots, M\}$. One can easily check the decomposition

$$g_{\ell}^{\rightarrow}(\eta) = \bar{\eta}(0) [\bar{\eta}(1) - \bar{\eta}^{\ell_0}(\ell_0)] \quad (5.5)$$

$$+ \bar{\eta}^{\ell_0}(\ell_0) [\bar{\eta}(0) - \bar{\eta}^{\ell_0}(0)] \quad (5.6)$$

$$+ \sum_{k=0}^{M-1} \bar{\eta}^{\ell_k}(0) [\bar{\eta}^{\ell_k}(\ell_k) - \bar{\eta}^{2\ell_k}(2\ell_k)] \quad (5.7)$$

$$+ \sum_{k=0}^{M-1} \bar{\eta}^{2\ell_k}(2\ell_k) [\bar{\eta}^{\ell_k}(0) - \bar{\eta}^{2\ell_k}(0)] \quad (5.8)$$

$$+ \bar{\eta}^{\ell}(0) [\bar{\eta}^{\ell}(\ell) - \bar{\eta}(\ell+1)] \quad (5.9)$$

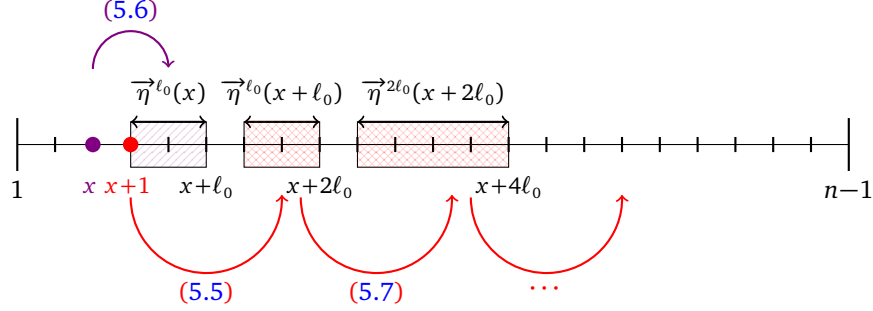
$$+ \bar{\eta}^{\ell}(0) [\bar{\eta}(\ell+1) - \bar{\eta}(0)] \quad (5.10)$$

$$+ \bar{\eta}^{\ell}(0) [\bar{\eta}(0) - \bar{\eta}^{\ell}(0)] + \frac{\chi(\rho)}{\ell}. \quad (5.11)$$

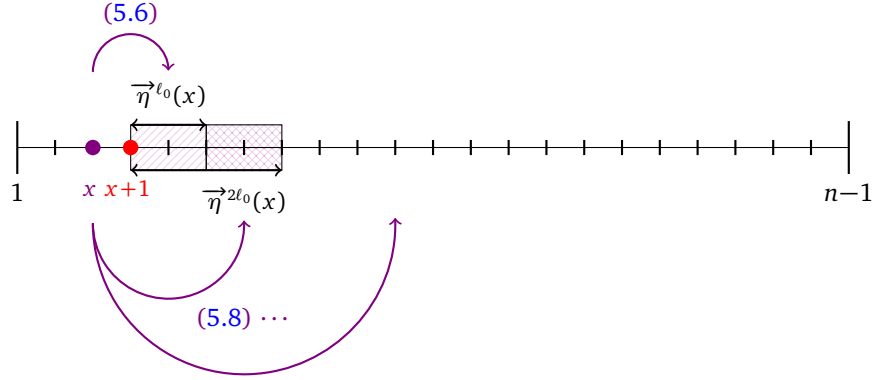
For example, in (5.5) we say that we *replace* $\bar{\eta}(1)$ by $\bar{\eta}^{\ell_0}(\ell_0)$, while $\bar{\eta}(0)$ is considered as *fixed*. Seven terms appear from (5.5) to (5.11). Let us denote them by order of appearance as follows:

$$g_I^{\rightarrow}(\eta), \quad g_{II}^{\rightarrow}(\eta), \quad g_{III}^{\rightarrow}(\eta), \quad g_{IV}^{\rightarrow}(\eta), \quad g_V^{\rightarrow}(\eta), \quad g_{VI}^{\rightarrow}(\eta), \quad g_{VII}^{\rightarrow}(\eta).$$

The decomposition above can naturally be written for $\tau_x g_{\ell}^{\rightarrow}$ ($x \in \Lambda_n$) by translating any term. Let us now illustrate the first steps of the decomposition: in Figure 4 below, we use the arrows as symbols for the replacements we perform, and we illustrate the consecutive replacements from (5.5) to (5.7), the latter corresponding to $\bar{\eta}^{\ell_k}(x + \ell_k) \mapsto \bar{\eta}^{2\ell_k}(x + 2\ell_k)$.

FIGURE 4. Illustration of steps (5.5)–(5.7) when $x \leq n - 2\ell - 1$.

Simultaneously, one can see in (5.8) that $\vec{\eta}^{\ell_k}(x)$ is replaced with $\vec{\eta}^{2\ell_k}(x)$ as we illustrate now in Figure 5:

FIGURE 5. Illustration of step (5.8): successive replacements when $x \leq n - 2\ell - 1$.

The role of the pre-factors $\vec{\eta}^{\ell_k}(x)$ in (5.7) and $\vec{\eta}^{2\ell_k}(x+2\ell_k)$ in (5.8) can be roughly understood as follows: these local functions have a variance of order $(\ell_k)^{-1}$ under ν_ρ^n , which compensates the price to pay when one tries to replace $\vec{\eta}^{\ell_k}(x+\ell_k)$ by $\vec{\eta}^{2\ell_k}(x+2\ell_k)$ and $\vec{\eta}^{\ell_k}(x)$ by $\vec{\eta}^{2\ell_k}(x)$. More precisely, this compensation is optimal if the support of the pre-factor does not intersect the set of sites which are used in the replacement: for example, the support of $\vec{\eta}^{\ell_k}(x)$ is $\{x+1, \dots, x+\ell_k\}$ and it does not intersect $\{x+\ell_k+1, \dots, x+4\ell_k\}$, which corresponds to the sites used in the replacement $\vec{\eta}^{\ell_k}(x+\ell_k) \mapsto \vec{\eta}^{2\ell_k}(x+2\ell_k)$, see (5.7). More details are given below.

The decomposition which works for g_ℓ^\leftarrow is very similar and there is no difficulty to find it out, following closely (5.5)–(5.11).

Let us go back to our goal estimate (5.1). From the standard convexity inequality $(a_1 + \dots + a_p)^2 \leq p(a_1^2 + \dots + a_p^2)$, one can see that (5.1) follows from seven independent estimates. More precisely, it is enough to prove that

$$\mathfrak{J}_{t,n}^{\text{left}}(g_w^\rightarrow) \leq Ct \left(\frac{\ell}{n} + \frac{tn}{\ell^2} \right) \|v\|_{2,n}^2, \quad \text{for any } w \in \{\text{I, II, III, IV, V, VI, VII}\}. \quad (5.12)$$

There is one further step in two particular cases: for $w = \text{III}$ and $w = \text{IV}$, we also use Minkowski's inequality, so that we can bound as follows:

$$\mathfrak{I}_{t,n}^{\text{left}}(g_{\text{III}}^{\rightarrow}) \leq C \left(\sum_{k=0}^{M-1} \left\{ \mathbb{E}_{\rho} \left[\left(\int_0^t \sum_{x=1}^{n-2\ell-1} v(x) \overrightarrow{\eta}_{sn^2}^{\ell_k}(x) \times \left[\overrightarrow{\eta}_{sn^2}^{\ell_k}(x + \ell_k) - \overrightarrow{\eta}_{sn^2}^{2\ell_k}(x + 2\ell_k) \right] ds \right)^2 \right] \right\}^{1/2} \right)^2 \quad (5.13)$$

$$\mathfrak{I}_{t,n}^{\text{left}}(g_{\text{IV}}^{\rightarrow}) \leq C \left(\sum_{k=0}^{M-1} \left\{ \mathbb{E}_{\rho} \left[\left(\int_0^t \sum_{x=1}^{n-2\ell-1} v(x) \overrightarrow{\eta}_{sn^2}^{2\ell_k}(x + 2\ell_k) \times \left[\overrightarrow{\eta}_{sn^2}^{\ell_k}(x) - \overrightarrow{\eta}_{sn^2}^{2\ell_k}(x) \right] ds \right)^2 \right] \right\}^{1/2} \right)^2. \quad (5.14)$$

Finally, the proof of (5.12) can almost be resumed in one general statement, which we are going to apply several times. Let us state here our main estimate:

Proposition 5.2. *Let \mathbf{A}, \mathbf{B} be two subsets of Λ_n , and let us denote by $\#\mathbf{B}$ the cardinality of \mathbf{B} . We assume that: for all $x \in \mathbf{A}$, the translated set*

$$\tau_x \mathbf{B} = \{x + y ; y \in \mathbf{B}\}$$

satisfies $\tau_x \mathbf{B} \subset \Lambda_n$. Consider $g : \Omega_n \rightarrow \mathbb{R}$ a local function whose support does not intersect \mathbf{B} , namely: $\text{Supp}(g) \cap \mathbf{B} = \emptyset$, and which has mean zero with respect to ν_{ρ}^n . Then, there exists $C > 0$ such that, for any $n \in \mathbb{N}$ and $t > 0$,

$$\mathbb{E}_{\rho} \left[\left(\int_0^t \sum_{x \in \mathbf{A}} \left\{ v(x) \tau_x g(\eta_{sn^2}) \sum_{z \in \tau_x \mathbf{B}} (\overline{\eta}_{sn^2}(z) - \overline{\eta}_{sn^2}(z + 1)) \right\} ds \right)^2 \right] \leq \frac{C t (\#\mathbf{B})^2}{n} \|g\|_2^2 \|\nu\|_{2,n}^2. \quad (5.15)$$

Proof of Proposition 5.2. We prove it in Section 5.3. \square

Before that, let us apply it to our purposes.

5.2. End of the proof of Theorem 4.2. First, let us prove that we can apply Proposition 5.2 in order to estimate $\mathfrak{I}_{t,n}^{\text{left}}(g_w^{\rightarrow})$, for $w \in \{\text{I}, \dots, \text{VI}\}$. The only estimate that has to be considered separately is the one involving $g_{\text{VII}}^{\rightarrow}$.

We prove that the assumptions of Proposition 5.2 are satisfied for $g_{\text{III}}^{\rightarrow}$ and $g_{\text{IV}}^{\rightarrow}$ (see also (5.13) and (5.14)) and we let the reader to check the other ones. First, recall (5.13) and notice that

$$\overrightarrow{\eta}^{\ell_k}(x) \left[\overrightarrow{\eta}^{\ell_k}(x + \ell_k) - \overrightarrow{\eta}^{2\ell_k}(x + 2\ell_k) \right] = \frac{\overrightarrow{\eta}^{\ell_k}(x)}{2\ell_k} \sum_{y=x+1}^{x+\ell_k} \sum_{z=y+\ell_k}^{y+2\ell_k-1} (\overline{\eta}(z) - \overline{\eta}(z + 1)) \quad (5.16)$$

$$+ \frac{\overrightarrow{\eta}^{\ell_k}(x)}{2\ell_k} \sum_{y=x+1}^{x+\ell_k} \sum_{z=y+\ell_k}^{y+3\ell_k-1} (\overline{\eta}(z) - \overline{\eta}(z + 1)). \quad (5.17)$$

Note that the above identity can be easily obtained by splitting each average $\overrightarrow{\eta}^{\ell_k}(x + \ell_k)$ and $\overrightarrow{\eta}^{2\ell_k}(x + 2\ell_k)$ in two parts, as follows:

$$\begin{aligned}\overrightarrow{\eta}^{\ell_k}(x + \ell_k) &= \frac{1}{2\ell_k} \sum_{y=1}^{\ell_k} \overrightarrow{\eta}(y + x + \ell_k) + \frac{1}{2\ell_k} \sum_{y=1}^{\ell_k} \overrightarrow{\eta}(y + x + \ell_k) \\ \overrightarrow{\eta}^{2\ell_k}(x + 2\ell_k) &= \frac{1}{2\ell_k} \sum_{y=1}^{\ell_k} \overrightarrow{\eta}(y + x + 2\ell_k) + \frac{1}{2\ell_k} \sum_{y=1}^{\ell_k} \overrightarrow{\eta}(y + x + 3\ell_k),\end{aligned}$$

and subtracting the sums one by one. Let us first deal with (5.16): we can use Proposition 5.2 with

$$g(\eta) = \frac{\overrightarrow{\eta}^{\ell_k}(0)}{2\ell_k}, \quad \text{and} \quad \mathbf{B} = \bigcup_{y=1}^{\ell_k} \{y + \ell_k, \dots, y + 2\ell_k - 1\}.$$

Note that $\|g\|_2^2 = C/\ell_k^3$ and $\#\mathbf{B} = \ell_k^2$, and remember that $\ell_k = 2^k \ell_0$. Next, we deal with (5.17): we only need to change \mathbf{B} which now reads

$$\mathbf{B} = \bigcup_{y=1}^{\ell_k} \{y + \ell_k, \dots, y + 3\ell_k - 1\}, \quad \text{hence} \quad \#\mathbf{B} = 2\ell_k^2.$$

Therefore, one can see from (5.13) and Proposition 5.2 that

$$\mathfrak{J}_{t,n}^{\text{left}}(g_{\text{III}}^{\rightarrow}) \leq C \left(\sum_{k=0}^{M-1} \left\{ \frac{t\ell_k^4}{n} \frac{1}{\ell_k^2} \|v\|_{2,n}^2 \right\}^{1/2} \right)^2 \leq C t \frac{\ell_M}{n} \|v\|_{2,n}^2.$$

Let us treat similarly $g_{\text{IV}}^{\rightarrow}$: recall (5.14) and write

$$\overrightarrow{\eta}^{2\ell_k}(x + 2\ell_k) [\overrightarrow{\eta}^{\ell_k}(x) - \overrightarrow{\eta}^{2\ell_k}(x)] = \frac{\overrightarrow{\eta}^{2\ell_k}(x + 2\ell_k)}{2\ell_k} \sum_{y=x+1}^{x+\ell_k} \sum_{z=y}^{y+\ell_k-1} (\overrightarrow{\eta}(z) - \overrightarrow{\eta}(z+1)). \quad (5.18)$$

Here we choose

$$g(\eta) = \frac{\overrightarrow{\eta}^{2\ell_k}(2\ell_k)}{2\ell_k} \quad \text{and} \quad \mathbf{B} = \bigcup_{y=1}^{\ell_k} \{y, \dots, y + \ell_k - 1\}$$

and then apply Proposition 5.2, which gives the same bound as before:

$$\mathfrak{J}_{t,n}^{\text{left}}(g_{\text{IV}}^{\rightarrow}) \leq C t \frac{\ell_M}{n} \|v\|_{2,n}^2.$$

Performing similar arguments and using Proposition 5.2 together with Minkowski's inequality, we get that, for any $\ell, n \in \mathbb{N}$ such that $\ell < \frac{n}{4}$, and any $t > 0$,

$$\begin{aligned}\mathfrak{J}_{t,n}^{\text{left}}(g_{\text{I}}^{\rightarrow}) &\leq C t \frac{\ell_0^2}{n} \|v\|_{2,n}^2, & \mathfrak{J}_{t,n}^{\text{left}}(g_{\text{II}}^{\rightarrow}) &\leq C t \frac{\ell_0}{n} \|v\|_{2,n}^2, \\ \mathfrak{J}_{t,n}^{\text{left}}(g_{\text{w}}^{\rightarrow}) &\leq C t \frac{\ell}{n} \|v\|_{2,n}^2, & \text{for any } w \in \{\text{III}, \text{IV}, \text{V}, \text{VI}\}.\end{aligned}$$

Finally, we have to estimate the last remaining term involving $g_{\text{VII}}^{\rightarrow}$, which is treated separately. More precisely, in Section 5.4 we will prove the following:

Proposition 5.3. *For any $\ell, n \in \mathbb{N}$ such that $\ell < \frac{n}{4}$, and any $t > 0$,*

$$\mathfrak{J}_{t,n}^{\text{left}}(g_{\text{VII}}^{\rightarrow}) = \mathbb{E}_\rho \left[\left(\int_0^t \sum_{x=1}^{n-2\ell-1} v(x) \left\{ \overrightarrow{\eta}_{sn^2}^\ell(x) [\overrightarrow{\eta}_{sn^2}(x) - \overrightarrow{\eta}_{sn^2}^\ell(x)] + \frac{\chi(\rho)}{\ell} \right\} ds \right)^2 \right]$$

$$\leq Ct \left(\frac{\ell}{n} + \frac{tn}{\ell^2} \right) \|v\|_{2,n}^2. \quad (5.19)$$

Putting together the previous estimates into our decomposition (5.5)–(5.11) of g_ℓ^\rightarrow , we obtain straightforwardly the final bound (5.1). We let the reader repeat all the arguments above to obtain the second part, namely (5.2). Theorem 4.2 then easily follows.

The next two sections are devoted to the proofs of Proposition 5.2 and Proposition 5.3.

5.3. Proof of Proposition 5.2. Take \mathbf{A}, \mathbf{B} two subsets of Λ_n such that, for all $x \in \mathbf{A}$, $\tau_x \mathbf{B} \subset \Lambda_n$, and take $g : \Omega_n \rightarrow \mathbb{R}$ a mean zero function with respect to ν_ρ^n such that $\text{Supp}(g) \cap \mathbf{B} = \emptyset$. From [32, Lemma 2.4], we bound the left hand side of (5.15) by

$$Ct \sup_{f \in L^2(\nu_\rho^n)} \left\{ 2 \int f(\eta) \sum_{x \in \mathbf{A}} v(x) \tau_x g(\eta) \sum_{z \in \tau_x \mathbf{B}} (\bar{\eta}(z) - \bar{\eta}(z+1)) \nu_\rho^n(d\eta) - n^2 \mathfrak{D}_n(f) \right\},$$

where \mathfrak{D}_n is the Dirichlet form introduced in (2.2). We write the previous expectation as twice its half, and in one of the integrals we make the change of variables $\eta \mapsto \sigma^{z,z+1} \eta$ to rewrite it as

$$\sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} v(x) \int \left\{ f(\eta) \tau_x g(\eta) - f(\sigma^{z,z+1} \eta) \tau_x g(\sigma^{z,z+1} \eta) \right\} (\bar{\eta}(z) - \bar{\eta}(z+1)) \nu_\rho^n(d\eta)$$

With our assumption, we have $\text{Supp}(\tau_x g) \cap \tau_x \mathbf{B} = \emptyset$ for every $x \in \mathbf{A}$. Therefore: $\tau_x g(\sigma^{z,z+1} \eta) = \tau_x g(\eta)$, for all $z \in \tau_x \mathbf{B}$, and as a consequence the last expression equals

$$\sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} v(x) \int \left\{ \tau_x g(\eta) (\bar{\eta}(z) - \bar{\eta}(z+1)) (f(\eta) - f(\sigma^{z,z+1} \eta)) \right\} \nu_\rho^n(d\eta).$$

For any $x \in \mathbf{A}$ and $z \in \tau_x \mathbf{B}$, we use Young's inequality with $\varepsilon_x > 0$, and we bound the previous expression from above by

$$\sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} \frac{\varepsilon_x}{2} (v(x))^2 \int \left\{ \tau_x g(\eta) (\bar{\eta}(z) - \bar{\eta}(z+1)) \right\}^2 \nu_\rho^n(d\eta) \quad (5.20)$$

$$+ \sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} \frac{1}{2\varepsilon_x} \int \left\{ f(\eta) - f(\sigma^{z,z+1} \eta) \right\}^2 \nu_\rho^n(d\eta). \quad (5.21)$$

Now, since ν_ρ^n is invariant under translations and since $\mathbf{A} \subset \Lambda_n$, it is easy to see that

$$\frac{1}{\#\mathbf{B}} \sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} \int \left\{ f(\eta) - f(\sigma^{z,z+1} \eta) \right\}^2 \nu_\rho^n(d\eta) \leq \mathfrak{D}_n(f).$$

As a result, if we chose $2\varepsilon_x = \#\mathbf{B}/n^2$, we have that (5.21) is bounded by $n^2 \mathfrak{D}_n(f)$, and (5.20) is bounded by

$$\sum_{x \in \mathbf{A}} \sum_{z \in \tau_x \mathbf{B}} \frac{\#\mathbf{B}}{4n^2} (v(x))^2 \|g\|_2^2 \leq \frac{(\#\mathbf{B})^2}{4n} \|v\|_{2,n}^2 \|g\|_2^2.$$

This ends the proof.

5.4. **Proof of Proposition 5.3.** We have

$$\mathfrak{J}_{t,n}^{\text{left}}(g_{\text{VII}}^{\rightarrow}) \leq 2\mathfrak{J}_{t,n}^{\text{left}}(g_1) + 2\mathfrak{J}_{t,n}^{\text{left}}(g_2),$$

where we define

$$\begin{aligned} g_1(\eta) &= \overrightarrow{\eta}^\ell(0) [\overrightarrow{\eta}(0) - \overrightarrow{\eta}^\ell(0)] + \frac{1}{2\ell} (\eta(0) - \eta(1))^2, \\ g_2(\eta) &= \frac{\chi(\rho)}{\ell} - \frac{1}{2\ell} (\eta(0) - \eta(1))^2. \end{aligned}$$

Let us start with the easiest term to estimate, namely $\mathfrak{J}_{t,n}^{\text{left}}(g_2)$. From the Cauchy-Schwarz inequality, together with the independence of $\eta(x)$ and $\eta(y)$ under the invariant measure ν_ρ^n (as soon as $x \neq y$), one can easily show that

$$\mathfrak{J}_{t,n}^{\text{left}}(g_2) \leq C t^2 \frac{n}{\ell^2} \|\nu\|_{2,n}^2.$$

Now let us look at the term with g_1 . From [32, Lemma 2.4], we bound $\mathfrak{J}_{t,n}^{\text{left}}(g_1)$ by

$$\begin{aligned} \sup_{f \in L^2(\nu_\rho^n)} \left\{ 2 \int \sum_{x=1}^{n-2\ell-1} \nu(x) \left\{ \overrightarrow{\eta}^\ell(x) [\overrightarrow{\eta}(x) - \overrightarrow{\eta}^\ell(x)] + \frac{1}{2\ell} (\overrightarrow{\eta}(x) - \overrightarrow{\eta}(x+1))^2 \right\} f(\eta) \nu_\rho^n(d\eta) \right. \\ \left. - n^2 \mathfrak{D}_n(f) \right\}. \quad (5.22) \end{aligned}$$

We note that

$$2 \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \{ \overrightarrow{\eta}(x) - \overrightarrow{\eta}^\ell(x) \} f(\eta) \nu_\rho^n(d\eta) \quad (5.23)$$

$$= 2 \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \{ \overrightarrow{\eta}(x) - \overrightarrow{\eta}(x+1) \} f(\eta) \nu_\rho^n(d\eta) \quad (5.24)$$

$$+ 2 \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \frac{\ell-1}{\ell} \{ \overrightarrow{\eta}(x+1) - \overrightarrow{\eta}(x+2) \} f(\eta) \nu_\rho^n(d\eta)$$

$$+ \dots + 2 \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \frac{1}{\ell} \{ \overrightarrow{\eta}(x+\ell-1) - \overrightarrow{\eta}(x+\ell) \} f(\eta) \nu_\rho^n(d\eta).$$

For each term of the last sum above, we do the following procedure: we write it as twice its half, and in one of the integrals we make the change of variables η to $\sigma^{z,z+1}\eta$ (for some suitable z), for which the measure ν_ρ^n is invariant. After doing this, one can check that the last expression equals

$$\begin{aligned} & \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \{ \overrightarrow{\eta}(x) - \overrightarrow{\eta}(x+1) \} (f(\eta) - f(\sigma^{x,x+1}\eta)) \nu_\rho^n(d\eta) \quad (5.25) \\ & + \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \frac{\ell-1}{\ell} \{ \overrightarrow{\eta}(x+1) - \overrightarrow{\eta}(x+2) \} (f(\eta) - f(\sigma^{x+1,x+2}\eta)) \nu_\rho^n(d\eta) \\ & + \dots + \\ & + \int \sum_{x=1}^{n-2\ell-1} \nu(x) \overrightarrow{\eta}^\ell(x) \frac{1}{\ell} \{ \overrightarrow{\eta}(x+\ell-1) - \overrightarrow{\eta}(x+\ell) \} (f(\eta) - f(\sigma^{x+\ell-1,x+\ell}\eta)) \nu_\rho^n(d\eta) \end{aligned}$$

$$+ \int \sum_{x=1}^{n-2\ell-1} v(x) \frac{\bar{\eta}(x+1) - \bar{\eta}(x)}{\ell} \{ \bar{\eta}(x) - \bar{\eta}(x+1) \} f(\eta) \nu_\rho^n(d\eta). \quad (5.26)$$

Note that the last term (5.26) comes from the change of variables η to $\sigma^{x,x+1}\eta$ in the first term (5.24) above, as well as (5.25). The whole sum can be rewritten as

$$\int \sum_{x=1}^{n-2\ell-1} v(x) \bar{\eta}^\ell(x) \frac{1}{\ell} \sum_{y=x+1}^{x+\ell} \sum_{z=x}^{y-1} \{ \bar{\eta}(z) - \bar{\eta}(z+1) \} \{ f(\eta) - f(\sigma^{z,z+1}\eta) \} \nu_\rho^n(d\eta) \quad (5.27)$$

$$- \int \sum_{x=1}^{n-2\ell-1} v(x) \frac{1}{\ell} (\bar{\eta}(x) - \bar{\eta}(x+1))^2 f(\eta) \nu_\rho^n(d\eta). \quad (5.28)$$

The integral in (5.22) is exactly equal to the sum of (5.23) and (5.28), therefore it is bounded by the first term in the previous expression, namely (5.27). Now, we use the same arguments as in the proof of Proposition 5.2, namely, Young's inequality with $2\varepsilon_x = \ell/n^2$ and we bound it by

$$\begin{aligned} C(\rho) \frac{\ell}{n^2} \sum_{x=1}^{n-2\ell-1} v^2(x) + \frac{n^2}{\ell^2} \sum_{x=1}^{n-2\ell-1} \sum_{y=x+1}^{x+\ell} \sum_{z=x}^{y-1} \int \{ f(\eta) - f(\sigma^{z,z+1}\eta) \}^2 \nu_\rho^n(d\eta) \\ \leq C(\rho) \frac{\ell}{n} \|v\|_{2,n}^2 + n^2 \mathfrak{D}_n(f), \end{aligned}$$

which proves that

$$\mathfrak{I}_{t,n}^{\text{left}}(g_1) \leq C t \frac{\ell}{n} \|v\|_{2,n}^2,$$

so that the proof ends. \square

6. UNIQUENESS OF ENERGY SOLUTIONS

In this section we give all the details for the proof of Theorem 3.3 and we show how the same arguments also apply to the proof of Theorem 3.7.

Recall that we are interested in energy solutions to

$$d\mathcal{Y}_t = A \Delta_{\text{Dir}} \mathcal{Y}_t dt + \bar{E} \nabla_{\text{Dir}} (\mathcal{Y}_t^2) dt + \sqrt{D} \nabla_{\text{Dir}} d\mathcal{W}_t, \quad (6.1)$$

where $A, D > 0$ and $\bar{E} \in \mathbb{R}$. But as in [25, Remark 2.6] we can show that $\{\mathcal{Y}_t; t \in [0, T]\}$ is an energy solution to (6.1) if and only if $\{\sqrt{2A/D} \mathcal{Y}_{t/A}; t \in [0, AT]\}$ solves (6.1) with $A \mapsto 1$, $D \mapsto 2$, and $\bar{E} \mapsto \bar{E} \sqrt{D/(2A^3)}$.

So from now on we assume without loss of generality that $A = 1$ and $D = 2$, and to simplify notation we write E instead of \bar{E} , and we show in this section that the equation

$$d\mathcal{Y}_t = \Delta_{\text{Dir}} \mathcal{Y}_t dt + E \nabla_{\text{Dir}} (\mathcal{Y}_t^2) dt + \sqrt{2} \nabla_{\text{Dir}} d\mathcal{W}_t, \quad (6.2)$$

has a unique solution for any $E \in \mathbb{R}$, where the notion of Dirichlet boundary conditions has been properly defined in Theorem 3.3.

To prove the uniqueness of energy solutions we will use the exact same strategy as in [25] and we will sometimes refer to that paper for additional details. The main idea consists in: first, mollifying an energy solution, then mapping the mollified process through the Cole-Hopf transform to a new process, and then taking the mollification away in order to show that the transformed process solves in the limit the linear multiplicative stochastic heat equation with certain boundary conditions. However, even

using the strategy of [25], we have to redo all computations because our setting is somewhat different, and the boundary condition of the stochastic heat equation actually changes as we pass to the limit: in particular, it is not equal to the one we would naively guess.

Remark 6.1. *Let us notice that the definition given in [25] is not exactly the same as the one we adopted in Theorem 3.3, but it is not difficult to check that they are indeed equivalent (see [20, Proposition 4]), and therefore the same strategy can be implemented.*

This section is split as follows: in Section 6.1 we give some tools that will be used in Section 6.2 in order to show that in the definition of the Burgers non-linearity (namely the process \mathcal{A} of Theorem 3.3) we can replace ι_ε by different approximations of the identity. We conclude in Section 6.3 with the proof of uniqueness for the energy solution \mathcal{Y} , by developing the strategy explained above. In the following we always denote by μ the law of the standard white noise on $\mathcal{S}'_{\text{Dir}}$. If $f : \mathbb{R}^n \rightarrow \mathbb{R}$ is multidimensional, then we denote by $\partial_\alpha f$ its derivative of order $\alpha \in \mathbb{N}_0^n$.

6.1. Preliminaries. In this section we give two ways to handle functionals written in the form $\int_0^\cdot F(\mathcal{Y}_s) ds$, where \mathcal{Y} shall be the energy solution to (6.2) as defined in Theorem 3.3, and F belongs to some general class of functions.

6.1.1. Itô trick and Kipnis-Varadhan Lemma. We write \mathcal{C} for the space of cylinder functions $F : \mathcal{S}'_{\text{Dir}} \rightarrow \mathbb{R}$, which are such that there exist $d \in \mathbb{N}$ and $\varphi_i \in \mathcal{S}_{\text{Dir}}$ ($i = 1, \dots, d$) with

$$F(\mathcal{Y}) = f(\mathcal{Y}(\varphi_1), \dots, \mathcal{Y}(\varphi_d)),$$

where $f \in \mathcal{C}^2(\mathbb{R}^d)$ has polynomial growth of all partial derivatives up to order 2. For $F \in \mathcal{C}$ we define the operator L_0 as

$$\begin{aligned} L_0 F(\mathcal{Y}) &= \sum_{i=1}^d \partial_i f(\mathcal{Y}(\varphi_1), \dots, \mathcal{Y}(\varphi_n)) \mathcal{Y}(\Delta \varphi_i) \\ &\quad + \sum_{i,j=1}^d \partial_{ij} f(\mathcal{Y}(\varphi_1), \dots, \mathcal{Y}(\varphi_n)) \langle \nabla \varphi_i, \nabla \varphi_j \rangle_{L^2([0,1])}, \end{aligned}$$

and its domain $\text{Dom}(L_0)$ is defined as the closure in $L^2(\mu)$ of \mathcal{C} with respect to the norm

$$\|F\|_{L^2(\mu)} + \|L_0 F\|_{L^2(\mu)}.$$

First, let us take $\tilde{\mathcal{Y}}$ as the Ornstein-Uhlenbeck process with Dirichlet boundary conditions as defined in Proposition 3.1 (with $A = 1$, $D = 2$) – or, equivalently, an energy solution of (6.2) with $E = 0$ (as defined in Theorem 3.3). Then, we have for every $F \in \mathcal{C}$,

$$F(\tilde{\mathcal{Y}}_t) = F(\tilde{\mathcal{Y}}_0) + \int_0^t L_0 F(\tilde{\mathcal{Y}}_s) ds + M_t^F,$$

where M^F is a continuous martingale with quadratic variation

$$\langle M^F \rangle_t = \int_0^t 2 \left\| \nabla_u D_u F(\tilde{\mathcal{Y}}_s) \right\|_{L_u^2([0,1])}^2 ds,$$

and where ∇_u is the usual derivative w.r.t. u and $D_u F$ denotes the Malliavin derivative defined in terms of the law of the white noise, i.e.

$$D_u F(\mathcal{Y}) = \sum_{i=1}^d \partial_i f(\mathcal{Y}(\varphi_1), \dots, \mathcal{Y}(\varphi_n)) \varphi_i(u), \quad \text{for any } u \in [0, 1].$$

In the following, for any $\mathcal{Y} \in \mathcal{S}'_{\text{Dir}}$ we denote

$$\mathcal{E}(F)(\mathcal{Y}) = 2 \left\| \nabla_u D_u F(\mathcal{Y}) \right\|_{L^2_u([0,1])}^2. \quad (6.3)$$

Now, let \mathcal{Y} be an energy solution to (6.2) the stochastic Burgers equation with Dirichlet boundary conditions, as defined in Theorem 3.3. Recall that \mathcal{Y}_0 is a $\mathcal{S}'_{\text{Dir}}$ -valued white noise, hence has law μ .

Since (3.5) implies that \mathcal{A} has zero quadratic variation – see [20, Proposition 4] for a proof – the Itô trick for additive functionals of the form $\int_0^\cdot L_0 F(\mathcal{Y}_s) ds$ follows by the same arguments as in [25, Proposition 3.2], i.e. by applying Itô's formula to the forward and backward process and adding up the resulting expressions to obtain a sort of “Lyons-Zheng decomposition” for additive functionals $\int_0^\cdot L_0 F(\mathcal{Y}_s) ds$ as sums of a forward and a backward martingales. Thus, we can prove that for all $F \in \mathcal{C}$ and $p \geq 1$

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t L_0 F(\mathcal{Y}_s) ds \right|^p \right] \lesssim T^{p/2} \mathbb{E} [\mathcal{E}(F)(\mathcal{Y}_0)^{p/2}], \quad (6.4)$$

and for $p = 2$ we get in particular

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t L_0 F(\mathcal{Y}_s) ds \right|^2 \right] \lesssim T \mathbb{E} [\mathcal{E}(F)(\mathcal{Y}_0)]. \quad (6.5)$$

For the sake of clarity, let us define from now on:

$$\|F\|_{1,0}^2 = \mathbb{E} [\mathcal{E}(F)(\mathcal{Y}_0)] = 2 \mathbb{E} [F(\mathcal{Y}_0)(-L_0 F(\mathcal{Y}_0))],$$

where the second equality follows from the Gaussian integration by parts rule, see [35, Lemma 1.2.1].

From this, let us now define two Hilbert spaces which will be useful in controlling additive functionals of \mathcal{Y} .

Definition 6.2. *Let us introduce an equivalence relation on \mathcal{C} by identifying F and G if $\|F - G\|_{1,0} = 0$, so that $\|\cdot\|_{1,0}$ becomes a norm on the equivalence classes. We write \mathfrak{H}_0^1 for the completion of the equivalence classes with respect to $\|\cdot\|_{1,0}$.*

For $F \in \mathcal{C}$ we define

$$\|F\|_{-1,0}^2 = \sup_{G \in \mathcal{C}} \{2\mathbb{E}[F(\mathcal{Y}_0)G(\mathcal{Y}_0)] - \|G\|_{1,0}^2\} = \sup_{\substack{G \in \mathcal{C} \\ \|G\|_{1,0}=1}} \mathbb{E}[F(\mathcal{Y}_0)G(\mathcal{Y}_0)]$$

and we identify F and G if $\|F - G\|_{-1,0} = 0$ and $\|F\|_{-1,0} < \infty$. We write \mathfrak{H}_0^{-1} for the completion of the equivalence classes with respect to $\|\cdot\|_{-1,0}$.

Remark 6.3. *It is possible to show that $\|F\|_{1,0} = 0$ if and only if F is constant, but we do not need this.*

Let us now extend the Itô trick to the entire domain of L_0 :

Lemma 6.4 (Itô trick). *Let $F \in \text{Dom}(L_0)$. Then $F \in \mathfrak{H}_0^1$ and*

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t L_0 F(\mathcal{Y}_s) ds \right|^2 \right] \lesssim T \|F\|_{1,0}^2. \quad (6.6)$$

If also $\mathcal{E}(F) \in L^{p/2}(\mu)$ for $p \geq 1$, then

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t L_0 F(\mathcal{Y}_s) ds \right|^p \right] \lesssim T^{p/2} \mathbb{E}[\mathcal{E}(F)(\mathcal{Y}_0)^{p/2}]. \quad (6.7)$$

Proof of Lemma 6.4. Since $F \in \text{Dom}(L_0)$, there exists a sequence of cylinder functions $\{F_n\}_{n \in \mathbb{N}}$ in \mathcal{C} such that $\|F_n - F\|_{L^2} + \|L_0(F_n - F)\|_{L^2}$ converges to zero. But then also

$$\lim_{n \rightarrow \infty} \|F_n - F\|_{1,0}^2 = \lim_{n \rightarrow \infty} \left\{ -2\mathbb{E}[(F_n - F)(\mathcal{Y}_0)L_0(F_n - F)(\mathcal{Y}_0)] \right\} = 0,$$

so $F \in \mathfrak{H}_0^1$ and (6.6) follows from (6.5). The second estimate follows by approximation from (6.4). \square

Remark 6.5. *If moreover $F(\mathcal{Y}_0)$ has a finite chaos expansion of length d (see Subsection 6.1.2 below for the definition of the chaos expansion), then $\mathcal{E}(F)(\mathcal{Y}_0)$ has a chaos expansion of length $2 \times (d - 1)$, and therefore all its moments are comparable and we can estimate $\mathbb{E}[\mathcal{E}(F)(\mathcal{Y}_0)^{p/2}] \simeq \mathbb{E}[\mathcal{E}(F)(\mathcal{Y}_0)]^{p/2} = \|F\|_{1,0}^p$.*

Corollary 6.6 (Kipnis-Varadhan inequality). *Let $F \in L^2(\mu) \cap \mathfrak{H}_0^{-1}$. Then*

$$\mathbb{E} \left[\left\| \int_0^\cdot F(\mathcal{Y}_s) ds \right\|_{p\text{-var}; [0, T]}^2 \right] + \mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t F(\mathcal{Y}_s) ds \right|^2 \right] \lesssim T \|F\|_{-1,0}^2,$$

where the p -variation norm $\|f\|_{p\text{-var}}$ of a function $f : [a, b] \rightarrow \mathbb{R}$ is defined as

$$\|f\|_{p\text{-var}; [a, b]}^p = \sup \left\{ \sum_{k=0}^{d-1} |f(t_{k+1}) - f(t_k)|^p ; d \in \mathbb{N}, a = t_0 < t_1 < \dots < t_d = b \right\}. \quad (6.8)$$

Proof of Corollary 6.6. The usual proof by duality works. The statement about the p -variation is shown as in [25, Corollary 3]. \square

As a result, with Lemma 6.4 and Corollary 6.6 we are provided with two important tools, which allow us to control in some sense $\int_0^\cdot F(\mathcal{Y}_s) ds$. Note that Lemma 6.4 is convenient only if one is able to write F as $L_0 G$, which may not be easy. If one cannot solve the Poisson equation $F = L_0 G$, then one relies on the variational norm of F given by Corollary 6.6.

The next paragraph is devoted to constructing solutions to the Poisson equation $L_0 G = F$ using the Gaussian structure of $L^2(\mu)$, which is now standard and fully detailed in [25, Section 3.2].

6.1.2. Gaussian analysis. In the following we develop some Gaussian analysis that is helpful for estimating the $\|\cdot\|_{1,0}$ and $\|\cdot\|_{-1,0}$ norms from above. We refer the reader to [35, 29] for details on Malliavin calculus and chaos decompositions.

Let \mathcal{Y} be a white noise on $\mathcal{S}'_{\text{Dir}}$ and write $\sigma(\mathcal{Y})$ for the sigma algebra generated by \mathcal{Y} . Then we can define a *chaos expansion* in $L^2(\sigma(\mathcal{Y}))$ as follows: for $d \in \mathbb{N}_0$ and $f \in L^2([0, 1]^d)$ we write $W_d(f)(\mathcal{Y})$ for the d -th order Wiener Itô-integral of f against \mathcal{Y} ,

$$W_d(f)(\mathcal{Y}) = \int_{[0, 1]^d} f(y_1, \dots, y_d) \mathcal{Y}(y_1) \cdots \mathcal{Y}(y_d) dy_1 \cdots dy_d,$$

see [35, Section 1.1.2] for the construction; occasionally we will simply write $W_d(f)$ instead of $W_d(f)(\mathcal{Y})$. Recall that $W_d(f) = W_d(\tilde{f})$ for $\tilde{f}(u) = \frac{1}{d!} \sum_{\sigma \in \mathfrak{S}_d} f(u_{\sigma(1)}, \dots, u_{\sigma(d)})$, where \mathfrak{S}_d is the symmetric group on $\{1, \dots, d\}$.

The chaos expansion of $F \in L^2(\sigma(\mathcal{Y}))$ is then given by

$$F = \sum_{d=0}^{\infty} W_d(f_d),$$

where $\{f_d \in L^2([0, 1]^d)\}_{d \in \mathbb{N}_0}$ are uniquely determined deterministic functions.

Let us denote by $L_{\text{Sym}}^2([0, 1]^d)$ the space of symmetric L^2 -functions. Since the d -th homogeneous chaos $\{W_d(f_d) ; f_d \in L_{\text{Sym}}^2([0, 1]^d)\}$ is also the closure of the span of all random variables of the form $\mathcal{Y} \mapsto H_d(\mathcal{Y}(\varphi))$, where H_d is the d -th Hermite polynomial and $\varphi \in \mathcal{S}_{\text{Dir}}$ with $\|\varphi\|_{L^2([0, 1])} = 1$, it will be convenient to write down the action of L_0 onto these random variables. Indeed, it is well-known that

$$H_d(\mathcal{Y}(\varphi)) = W_d(\varphi^{\otimes d})(\mathcal{Y})$$

for all $\varphi \in L^2([0, 1])$ with $\|\varphi\|_{L^2([0, 1])} = 1$. For all $\varphi \in \mathcal{S}_{\text{Dir}}$ we have

$$\begin{aligned} L_0 H_d(\mathcal{Y}(\varphi)) &= H'_d(\mathcal{Y}(\varphi)) \mathcal{Y}(\Delta \varphi) + H''_d(\mathcal{Y}(\varphi)) \langle \nabla \varphi, \nabla \varphi \rangle_{L^2([0, 1])} \\ &= H'_d(\mathcal{Y}(\varphi)) \mathcal{Y}(\Delta \varphi) - H''_d(\mathcal{Y}(\varphi)) \langle \varphi, \Delta \varphi \rangle_{L^2([0, 1])}. \end{aligned} \quad (6.9)$$

From here let us define

$$\begin{aligned} \mathcal{S}_{\text{Dir}}([0, 1]^d) &= \left\{ g \in \mathcal{C}^\infty([0, 1]^d) ; \partial_\alpha g(u) = 0, \forall u \in \partial([0, 1]^d), \right. \\ &\quad \left. \forall \alpha = (2k_1, \dots, 2k_d) \in \mathbb{N}_0^d \right\}. \end{aligned}$$

Then, from (6.9), the same arguments as in [25, Lemma 3.7] show that for all symmetric functions $f_d \in \mathcal{S}_{\text{Dir}}([0, 1]^d)$ we have

$$L_0 W_d(f_d) = W_d(\Delta f_d), \quad (6.10)$$

with $\Delta := \sum_{k=1}^d \partial_{kk}$. Therefore, the operator L_0 leaves the d -th chaos invariant, and this will be useful to solve the Poisson equation. It only remains to compute its norm $\|W_d(f_d)\|_{1,0}$, which is the goal of the remainder of this section.

To do so, let us introduce other notations: similarly to our definitions in Section 2.3, we denote by $\mathcal{H}_{\text{Dir}}^1([0, 1]^d)$ the completion of $\mathcal{S}_{\text{Dir}}([0, 1]^d)$ with respect to the norm

$$\|g\|_{\mathcal{H}_{\text{Dir}}^1([0, 1]^d)} = \|\nabla g\|_{L^2([0, 1]^d)} = \left(\sum_{i=1}^d \|\partial_i g\|_{L^2([0, 1]^d)}^2 \right)^{1/2}$$

and $\mathcal{H}_{\text{Dir}, \text{Sym}}^1([0, 1]^d)$ denotes the completion of $\mathcal{S}_{\text{Dir}}([0, 1]^d)$ with respect to the norm

$$\|g\|_{\mathcal{H}_{\text{Dir}, \text{Sym}}^1([0, 1]^d)} = \|\tilde{g}\|_{\mathcal{H}_{\text{Dir}}^1([0, 1]^d)} \quad \text{for } \tilde{g}(u) = \frac{1}{d!} \sum_{\sigma \in \mathfrak{S}_d} g(u_{\sigma(1)}, \dots, u_{\sigma(d)}),$$

and therefore we identify non-symmetric functions g in $\mathcal{S}_{\text{Dir}}([0, 1]^d)$ with their symmetrizations \tilde{g} . From the Poincaré inequality we obtain that $\|g\|_{L^2([0, 1]^d)} \lesssim \|g\|_{\mathcal{H}_{\text{Dir}}^1([0, 1]^d)}$ for all $g \in \mathcal{S}_{\text{Dir}}([0, 1]^d)$, and therefore $\mathcal{H}_{\text{Dir}, \text{Sym}}^1([0, 1]^d)$ is contained in $L_{\text{Sym}}^2([0, 1]^d)$. Note also that the symmetric $\mathcal{S}_{\text{Dir}}([0, 1]^d)$ -functions are dense in $\mathcal{H}_{\text{Dir}, \text{Sym}}^1([0, 1]^d)$.

From (6.10), in the same way as in [25, Lemma 3.13], we obtain:

Lemma 6.7. *For symmetric $f_d \in \mathcal{S}_{\text{Dir}}([0, 1]^d)$ we have*

$$\begin{aligned} \|W_d(f_d)\|_{1,0}^2 &= -2\mathbb{E}[W_d(f_d)W_d(\Delta f_d)] = -2(d!)\langle f_d, \Delta f_d \rangle_{L^2([0,1]^d)}^2 \\ &= 2(d!)\|\nabla f_d\|_{L^2([0,1]^d)}^2 = 2(d!)\|f_d\|_{\mathcal{H}_{\text{Dir}}^1([0,1]^d)}^2. \end{aligned}$$

Moreover, for general (not necessarily symmetric) $f_d \in \mathcal{S}_{\text{Dir}}([0, 1]^d)$ we obtain the inequality

$$\|W_d(f_d)\|_{1,0}^2 = 2(d!)\|f_d\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^d)}^2 \leq 2(d!)\|f_d\|_{\mathcal{H}_{\text{Dir}}^1([0,1]^d)}^2.$$

Proof of Lemma 6.7. Integration by parts works without boundary terms because $f_d \in \mathcal{S}_{\text{Dir}}([0, 1]^d)$. The bound $\|f_d\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^d)}^2 \leq \|f_d\|_{\mathcal{H}_{\text{Dir}}^1([0,1]^d)}^2$ follows from the triangle inequality. \square

We also write $\mathcal{H}_{\text{Dir,Sym}}^{-1}([0, 1]^d)$ for the completion of $L_{\text{Sym}}^2([0, 1]^d)$ with respect to the norm

$$\|f\|_{\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^d)} = \sup_{\substack{g \in \mathcal{H}_{\text{Dir,Sym}}^1([0,1]^d) \\ \|g\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^d)} = 1}} \langle f, g \rangle_{L_{\text{Sym}}^2([0,1]^d)}. \quad (6.11)$$

In all what follows, to simplify notation, whenever $d = 1$ we will simply denote, as before

$$\mathcal{H}_{\text{Dir}}^1 = \mathcal{H}_{\text{Dir}}^1([0, 1]) = \mathcal{H}_{\text{Dir,Sym}}^1([0, 1]), \quad \text{and} \quad \mathcal{H}_{\text{Dir}}^{-1} = \mathcal{H}_{\text{Dir,Sym}}^{-1}([0, 1]).$$

If we need to make the variable y precise we will highlight it by writing $\mathcal{H}_{\text{Dir}}^1(y)$ and $\mathcal{H}_{\text{Dir}}^{-1}(y)$.

Lemma 6.8. *For $f \in \mathcal{H}_{\text{Dir,Sym}}^1([0, 1]^d)$ we have the more explicit representation*

$$\|f\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^d)}^2 = \sum_{k \in \mathbb{N}^d} |\pi k|^2 \langle f, e_{k_1} \otimes \cdots \otimes e_{k_d} \rangle_{L^2([0,1]^d)}^2$$

and for any $g \in \mathcal{H}_{\text{Dir,Sym}}^{-1}([0, 1]^d)$

$$\|g\|_{\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^d)}^2 = \sum_{k \in \mathbb{N}^d} |\pi k|^{-2} \langle g, e_{k_1} \otimes \cdots \otimes e_{k_d} \rangle_{L^2([0,1]^d)}^2,$$

where $e_m(u) = \sqrt{2} \sin(m\pi u)$ has been already defined in Section 2.3 and where $|k|^2 = |k_1|^2 + \cdots + |k_d|^2$ is the squared Euclidean norm of $k = (k_1, \dots, k_d) \in \mathbb{N}^d$.

Proof of Lemma 6.8. Recall that $\{e_m = \sqrt{2} \sin(m\pi \cdot) ; m \geq 1\}$, is an orthonormal basis of $L^2([0, 1])$, as is $\{\tilde{e}_0 \equiv 1, \tilde{e}_m = \sqrt{2} \cos(m\pi \cdot) ; m \geq 1\}$. Therefore,

$$\{e_{k_1} \otimes \cdots \otimes e_{k_{i-1}} \otimes \tilde{e}_\ell \otimes e_{k_i} \otimes \cdots \otimes e_{k_{d-1}} ; k \in \mathbb{N}^{d-1}, \ell \in \mathbb{N}_0\}$$

is an orthonormal basis of $L^2([0, 1]^d)$ for all $i = 1, \dots, d$. If we use this orthonormal basis to compute $\|\partial_i f\|_{L_{\text{Sym}}^2([0,1]^d)}^2$ and apply integration by parts, then the first equality follows. For the second equality we use the duality of $\mathcal{H}_{\text{Dir,Sym}}^{-1}([0, 1]^d)$ and $\mathcal{H}_{\text{Dir,Sym}}^1([0, 1]^d)$. \square

We have now all at hand to state and prove the main result of this section:

Corollary 6.9. *The Poisson equation*

$$L_0 F = W_d(g)$$

is solvable for all $g \in L^2_{\text{Sym}}([0, 1]^d)$ and moreover

$$\|F\|_{1,0}^2 = 2(d!) \|g\|_{\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^d)}^2.$$

Proof of Corollary 6.9. Since L_0 leaves the d -th chaos invariant, F must be of the form $F = W_d(f)$ and thus $L_0 F = W_d(g)$ is equivalent to $\Delta f = g$, which has the explicit solution

$$f = \sum_{k \in \mathbb{N}^d} |\pi k|^{-2} \langle g, e_{k_1} \otimes \cdots \otimes e_{k_d} \rangle e_{k_1} \otimes \cdots \otimes e_{k_d},$$

so Lemma 6.7 and Lemma 6.8 give

$$\begin{aligned} \|F\|_{1,0}^2 &= 2(d!) \|f\|_{\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^d)}^2 = 2(d!) \sum_{k \in \mathbb{N}^d} 2^d |\pi k|^2 \langle f, \sin(\pi k_1 \cdot) \otimes \cdots \otimes \sin(\pi k_d \cdot) \rangle^2 \\ &= \sum_{k \in \mathbb{N}^d} 2^d |\pi k|^{-2} \langle g, \sin(\pi k_1 \cdot) \otimes \cdots \otimes \sin(\pi k_d \cdot) \rangle^2 = 2(d!) \|g\|_{\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^d)}^2. \end{aligned}$$

□

6.2. Burgers/KPZ non-linearity and existence of the process \mathcal{A}/\mathcal{B} . All along this section, the integration spaces denoted L^p , if not made precise, are in fact $L^p([0, 1])$.

Lemma 6.10. *There exists a unique process $\int_0^\cdot \mathcal{Y}_s^{\circ 2} ds \in \mathcal{C}(\mathbb{R}_+, \mathcal{S}'_{\text{Neu}})$ such that for all $\psi \in \mathcal{S}_{\text{Neu}}$, all $p \geq 1$ and all $\rho \in L^\infty([0, 1]^2)$ the following bound holds:*

$$\begin{aligned} \mathbb{E} \left[\sup_{t \in [0, T]} \left| \left(\int_0^t \mathcal{Y}_s^{\circ 2} ds \right) (\psi) - \int_0^t \left(\int_0^1 \{ \mathcal{Y}_s^2(\rho(u, \cdot)) - \|\rho(u, \cdot)\|_{L^2}^2 \} \psi(u) du \right) ds \right|^p \right] \\ \lesssim T^{p/2} \left(\sup_{u \in [0, 1]} \{ \|\rho(u, \cdot)\|_{L^1} \times \|\cdot - u\|_{L^1}^{1/2} \} + \|\langle \rho(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \right)^p \\ \times \left(\sup_{u \in [0, 1]} \{ \|\rho(\cdot, u)\|_{L^1} + \|\rho(u, \cdot)\|_{L^1} \} \right)^p \|\psi\|_{L^\infty}^p. \end{aligned} \quad (6.12)$$

Proof of Lemma 6.10. Let us define

$$\tilde{\mathcal{Y}}^{\circ 2}(\rho(u, \cdot)) = \mathcal{Y}^2(\rho(u, \cdot)) - \|\rho(u, \cdot)\|_{L^2}^2.$$

The strategy to define the process $(\int_0^\cdot \mathcal{Y}_s^{\circ 2} ds)(\psi)$ is to obtain it as the limit of

$$\int_0^\cdot \int_0^1 \tilde{\mathcal{Y}}_s^{\circ 2}(\rho_m(u, \cdot)) \psi(u) du ds \quad (6.13)$$

for some suitable sequence $\{\rho_m\}_{m \in \mathbb{N}}$. First, observe that

$$\int_0^1 \tilde{\mathcal{Y}}^{\circ 2}(\rho(u, \cdot)) \psi(u) du = W_2 \left(\int_0^1 \rho(u, \cdot)^{\otimes 2} \psi(u) du \right) (\mathcal{Y}). \quad (6.14)$$

So if $\{\rho_m\}_{m \in \mathbb{N}}$ is a sequence of functions in $L^\infty([0, 1]^2)$, we can use Corollary 6.9 to solve the Poisson equation with (6.14) and $\rho = \rho_m$ and then estimate the norm of the solution with the Itô trick given in Lemma 6.4. This gives, for any $m, n \in \mathbb{N}$

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t \left(\int_0^1 \tilde{\mathcal{Y}}_s^{\circ 2}(\rho_m(u, \cdot)) \psi(u) du \right) ds - \int_0^t \left(\int_0^1 \tilde{\mathcal{Y}}_s^{\circ 2}(\rho_n(u, \cdot)) \psi(u) du \right) ds \right|^p \right]$$

$$\lesssim T^{p/2} \left\| \int_0^1 (\rho_m(u, \cdot)^{\otimes 2} - \rho_n(u, \cdot)^{\otimes 2}) \psi(u) du \right\|_{\mathcal{H}_{\text{Dir}, \text{Sym}}^{-1}([0,1]^2)}^p.$$

To bound the norm on the right hand side we argue by duality and apply (6.11): let f be a symmetric function in $\mathcal{S}_{\text{Dir}}([0,1]^2)$ and consider

$$\begin{aligned} & \left\langle \int_0^1 (\rho_m(u, \cdot)^{\otimes 2} - \rho_n(u, \cdot)^{\otimes 2}) \psi(u) du, f \right\rangle_{L^2([0,1]^2)} \\ &= \int_{[0,1]^3} dudv_1 dv_2 (\rho_m(u, v_1) \rho_m(u, v_2) - \delta_0(u - v_1) \delta_0(u - v_2)) \psi(u) f(v_1, v_2) \\ &+ \int_{[0,1]^3} dudv_1 dv_2 (\delta_0(u - v_1) \delta_0(u - v_2) - \rho_n(u, v_1) \rho_n(u, v_2)) \psi(u) f(v_1, v_2), \end{aligned}$$

where $\delta_a(\cdot)$ denotes the Dirac delta function at point a . Both terms on the right hand side are of the same form, so we argue for the first one only. For this purpose we decompose

$$\begin{aligned} & \rho_m(u, v_1) \rho_m(u, v_2) - \delta_0(u - v_1) \delta_0(u - v_2) \\ &= (\rho_m(u, v_1) - \delta_0(u - v_1)) \rho_m(u, v_2) + \delta_0(u - v_1) (\rho_m(u, v_2) - \delta_0(u - v_2)) \end{aligned}$$

and again only treat the first contribution. In the following list of inequalities, each step will be made clear by using a notation of the form $(*)$ over the inequality \leq in order to explain where does the inequality come from. We are going to: sum and subtract $f(u, v_2)$ and use the triangular inequality (denoted by $\pm f(u, v_2)$); use the Cauchy-Schwartz inequality (C-S); use L^∞ bounds (L^∞); and finally, use the fact that for any $\varphi \in \mathcal{S}_{\text{Dir}}$, $\|\varphi\|_{L^\infty} \leq \|\nabla \varphi\|_{L^2}$, which will be denoted below by $(\star\star)$. Let us bound as follows:

$$\begin{aligned} & \left| \int_{[0,1]^3} dudv_1 dv_2 (\rho_m(u, v_1) - \delta_0(u - v_1)) \rho_m(u, v_2) \psi(u) f(v_1, v_2) \right| \\ & \stackrel{\pm f(u, v_2)}{\leq} \left| \int_{[0,1]^3} dudv_1 dv_2 \rho_m(u, v_1) (f(v_1, v_2) - f(u, v_2)) \rho_m(u, v_2) \psi(u) \right| \\ & + \left| \int_{[0,1]^2} dudv_2 (\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1) f(u, v_2) \rho_m(u, v_2) \psi(u) \right| \\ & \stackrel{(C-S)}{\leq} \left| \int_{[0,1]^3} dudv_1 dv_2 \rho_m(u, v_1) \int_u^{v_1} dw \partial_1 f(w, v_2) \rho_m(u, v_2) \psi(u) \right| \\ & + \|\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \left(\int_0^1 du \left| \int_0^1 dv_2 f(u, v_2) \rho_m(u, v_2) \psi(u) \right|^2 \right)^{1/2} \\ & \stackrel{(C-S)}{\leq} \int_{[0,1]^3} dudv_1 dv_2 |\rho_m(u, v_1)| \times |v_1 - u|^{1/2} \|\partial_1 f(\cdot, v_2)\|_{L^2} |\rho_m(u, v_2) \psi(u)| \\ & + \|\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \left(\int_0^1 du \|\rho_m(u, \cdot)\|_{L^1}^2 \|f(u, \cdot)\|_{L^\infty}^2 \|\psi\|_{L^\infty}^2 \right)^{1/2} \\ & \stackrel{L^\infty + (\star\star)}{\leq} \sup_{u \in [0,1]} \left\{ \|\rho_m(u, \cdot)\|_{L^1} \times |\cdot - u|^{1/2} \right\} \sup_{v \in [0,1]} \left\{ \|\rho_m(\cdot, v)\|_{L^1} \right\} \|\partial_1 f\|_{L^2([0,1]^2)} \|\psi\|_{L^\infty} \end{aligned}$$

$$\begin{aligned}
& + \|\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \sup_{u \in [0,1]} \left\{ \|\rho_m(u, \cdot)\|_{L^1} \right\} \|\partial_2 f\|_{L^2([0,1]^2)} \|\psi\|_{L^\infty} \\
& \leq \left(\sup_{u \in [0,1]} \left\{ \|\rho_m(u, \cdot)\|_{L^1} \times \|\cdot - u\|_{L^1}^{1/2} \right\} + \|\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \right) \\
& \quad \times \sup_{u \in [0,1]} \left\{ \|\rho_m(\cdot, u)\|_{L^1} + \|\rho_m(u, \cdot)\|_{L^1} \right\} \times \|f\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^2)} \|\psi\|_{L^\infty},
\end{aligned}$$

where in the last step we used the definition of the norm $\|\cdot\|_{\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^2)}$. By the density of the symmetric $\mathcal{S}_{\text{Dir}}([0,1]^2)$ -functions in $\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^2)$ and the duality of $\mathcal{H}_{\text{Dir,Sym}}^{-1}([0,1]^2)$ and $\mathcal{H}_{\text{Dir,Sym}}^1([0,1]^2)$ we conclude that

$$\begin{aligned}
& \mathbb{E} \left[\sup_{t \in [0,T]} \left| \int_0^t \left(\int_0^1 \tilde{\mathcal{V}}_s^{\circ 2}(\rho_m(u, \cdot)) \psi(u) du \right) ds - \int_0^t \left(\int_0^1 \tilde{\mathcal{V}}_s^{\circ 2}(\rho_n(u, \cdot)) \psi(u) du \right) ds \right|^p \right] \\
& \lesssim T^{p/2} \|\psi\|_{L^\infty}^p \\
& \quad \times \left[\left(\sup_{u \in [0,1]} \left\{ \|\rho_m(u, \cdot)\|_{L^1} \times \|\cdot - u\|_{L^1}^{1/2} \right\} + \|\langle \rho_m(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \right) \right. \\
& \quad \times \sup_{u \in [0,1]} \left\{ \|\rho_m(\cdot, u)\|_{L^1} + \|\rho_m(u, \cdot)\|_{L^1} \right\} \\
& \quad + \left(\sup_{u \in [0,1]} \left\{ \|\rho_n(u, \cdot)\|_{L^1} \times \|\cdot - u\|_{L^1}^{1/2} \right\} + \|\langle \rho_n(u, \cdot), 1 \rangle_{L^2} - 1\|_{L_u^2} \right) \\
& \quad \times \sup_{u \in [0,1]} \left\{ \|\rho_n(\cdot, u)\|_{L^1} + \|\rho_n(u, \cdot)\|_{L^1} \right\} \left. \right]^p.
\end{aligned}$$

Now let us choose the sequence $\{\rho_m\}_{m \in \mathbb{N}}$. For that purpose, let p^{Dir} be the Dirichlet heat kernel on $[0, 1]$, i.e.

$$p_t^{\text{Dir}}(u, v) = \sum_{k=1}^{\infty} e^{-t\pi^2 k^2} e_k(u) e_k(v), \quad u, v \in [0, 1], t > 0, \quad (6.15)$$

and let us set $\rho_m(u, v) = p_{1/m}^{\text{Dir}}(u, v)$ and $m \leq n$. Note that we have, for any $f \in L^2([0, 1])$,

$$\lim_{\varepsilon \rightarrow 0} \int_0^1 |\langle p_\varepsilon^{\text{Dir}}(u, \cdot), f \rangle_{L^2([0,1])} - f(u)|^2 du = 0.$$

The other properties of p^{Dir} are given in Appendix C: from Lemma C.1 we obtain that

$$\begin{aligned}
& \mathbb{E} \left[\sup_{t \in [0,T]} \left| \int_0^t \left(\int_0^1 \tilde{\mathcal{V}}_s^{\circ 2}(\rho_m(u, \cdot)) \psi(u) du \right) ds - \int_0^t \left(\int_0^1 \tilde{\mathcal{V}}_s^{\circ 2}(\rho_n(u, \cdot)) \psi(u) du \right) ds \right|^p \right] \\
& \lesssim T^{p/2} \|\psi\|_{L^\infty}^p (m^{-1/4} + n^{-1/4})^p \lesssim T^{p/2} \|\psi\|_{L^\infty}^p m^{-p/4}.
\end{aligned}$$

Therefore, the sequence (6.13) is Cauchy and there exists a limit in $\mathcal{C}(\mathbb{R}_+, \mathbb{R})$ which we denote with $(\int_0^\cdot \mathcal{V}_s^{\circ 2} ds)(\psi)$. Making use of the bound in terms of $\|\psi\|$, similar arguments as in Section 4.3 show that even $\psi \mapsto \int_0^\cdot \int_0^1 \tilde{\mathcal{V}}_s^{\circ 2}(\rho_m(u, \cdot)) \psi(u) du ds$ converges in $\mathcal{C}(\mathbb{R}_+, \mathcal{S}'_{\text{Neu}})$. By the computation above the limit satisfies (6.12) and clearly that estimate identifies $\int_0^\cdot \mathcal{V}_s^{\circ 2} ds$ uniquely as the limit of $\int_0^\cdot \tilde{\mathcal{V}}_s^{\circ 2}(\rho_m(u, \cdot)) ds$. \square

Corollary 6.11. *We have $-(\int_0^\cdot \mathcal{Y}_s^{\circ 2} ds)(\nabla \varphi) = \mathcal{A}(\varphi)$ for all $\varphi \in \mathcal{S}_{\text{Dir}}$, where \mathcal{A} denotes the process defined in (3.7) from the statement of Theorem 3.3.*

Similarly $(\int_0^\cdot \mathcal{Y}_s^{\circ 2} ds)(\varphi) = \mathcal{B}(\varphi)$ for all $\varphi \in \mathcal{S}_{\text{New}}$, where \mathcal{B} is the process defined in (3.10).

Proof of Corollary 6.11. According to Lemma 6.10 we only need to verify that the approximation of the identity given in Definition 3.2, namely:

$$\iota_\varepsilon(u, v) = \begin{cases} \varepsilon^{-1} \mathbf{1}_{]u, u+\varepsilon]}(v) & \text{if } u \in [0, 1-2\varepsilon), \\ \varepsilon^{-1} \mathbf{1}_{[u-\varepsilon, u[}(v) & \text{if } u \in [1-2\varepsilon, 1]. \end{cases}$$

satisfies

$$\lim_{\varepsilon \rightarrow 0} \left\{ \left(\sup_{u \in [0, 1]} \left\{ \|\iota_\varepsilon(u, \cdot)\| \times \|\cdot - u\|^{1/2} \right\}_{L^1} + \|\langle \iota_\varepsilon(u, \cdot), 1 \rangle - 1\|_{L^2_u} \right)^p \times \left(\sup_{v \in [0, 1]} \left\{ \|\iota_\varepsilon(\cdot, v)\|_{L^1} + \|\iota_\varepsilon(v, \cdot)\|_{L^1} \right\} \right)^p \right\} = 0.$$

Clearly $\langle \iota_\varepsilon(u, \cdot), 1 \rangle_{L^2} - 1 = 0$ and $\|\iota_\varepsilon(u, \cdot)\|_{L^1} \equiv 1$ for all $u \in [0, 1]$, while

$$\|\iota_\varepsilon(u, \cdot)\| \times \|\cdot - u\|^{1/2}_{L^1} \leq \varepsilon^{1/2} \|\iota_\varepsilon(u, \cdot)\|_{L^1} = \varepsilon^{1/2},$$

and therefore it remains to show that $\|\iota_\varepsilon(\cdot, v)\|_{L^1}$ is uniformly bounded in $v \in [0, 1]$ and $\varepsilon \in (0, 1]$. But $\iota_\varepsilon(u, v) = 0$ unless $|v - u| \leq \varepsilon$, and therefore $\|\iota_\varepsilon(\cdot, v)\|_{L^1} \leq 2$. \square

By a similar interpolation argument as in the proof of [25, Corollary 3.17] we get the following result:

Lemma 6.12. *For all $\alpha < \frac{3}{4}$ and $\varphi \in \mathcal{S}_{\text{New}}$, the process*

$$\left\{ \left(\int_0^t \mathcal{Y}_s^{\circ 2} ds \right)(\varphi); t \in [0, T] \right\}$$

is almost surely in $\mathcal{C}^\alpha([0, T], \mathbb{R})$, the space of α -Hölder continuous functions from $[0, T]$ to \mathbb{R} . Moreover, writing $p_\varepsilon^{\text{Dir}}$ for the Dirichlet heat kernel as defined in (6.15) and

$$\|f\|_{\mathcal{C}^\alpha([0, T], \mathbb{R})} := \sup_{0 \leq s < t \leq T} \frac{|f(t) - f(s)|}{|t - s|^\alpha},$$

we have for all $p \in [1, \infty)$

$$\mathbb{E} \left[\left\| \left(\int_0^\cdot \mathcal{Y}_s^{\circ 2} ds \right)(\varphi) - \int_0^\cdot \int_0^1 \left\{ \mathcal{Y}_s^2(p_\varepsilon^{\text{Dir}}(u, \cdot)) - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 \right\} \varphi(u) du ds \right\|_{\mathcal{C}^\alpha([0, T], \mathbb{R})}^p \right] \xrightarrow{\varepsilon \rightarrow 0} 0.$$

In particular, it follows together with Corollary 6.11 that $\mathcal{B}(\varphi)$ has zero quadratic variation.

6.3. Mapping to the stochastic heat equation and conclusion. To prove the uniqueness of our energy solution \mathcal{Y} to (6.2), we would like to apply the Cole-Hopf transform to map \mathcal{Y} to a solution of the well posed stochastic heat equation. To do so, we should integrate \mathcal{Y} in the space variable and then exponentiate the resulting process. But since we only have an explicit description of the dynamics of \mathcal{Y} after testing against a test function $\varphi \in \mathcal{S}_{\text{Dir}}$, we should first mollify \mathcal{Y} with a kernel in \mathcal{S}_{Dir} before carrying out this program.

We find it convenient to mollify \mathcal{Y} with the Dirichlet heat kernel p^{Dir} which was defined in (6.15), and we set

$$\mathcal{Y}_t^\varepsilon(u) = \mathcal{Y}_t(p_\varepsilon^{\text{Dir}}(u, \cdot)).$$

The unique antiderivative $\Theta_u^\varepsilon(v)$ which satisfies $\nabla_u \Theta_u^\varepsilon(v) = p_\varepsilon^{\text{Dir}}(u, v)$ and $\int_0^1 \Theta_u^\varepsilon(v) du = 0$ for all v is explicitly given as

$$\Theta_u^\varepsilon(v) = \sum_{\ell=1}^{\infty} e^{-\varepsilon \pi^2 \ell^2} (-\ell \pi)^{-1} 2 \cos(\ell \pi u) \sin(\ell \pi v), \quad u, v \in [0, 1], \varepsilon > 0.$$

Note that $\Theta_u^\varepsilon \in \mathcal{S}_{\text{Dir}}$ for all $u \in [0, 1]$ and $\varepsilon > 0$, so to integrate \mathcal{Y}^ε in the space variable we set

$$\mathcal{Z}_t^\varepsilon(u) = \mathcal{Y}_t(\Theta_u^\varepsilon),$$

which is a smooth function with $\nabla \mathcal{Z}^\varepsilon = \mathcal{Y}^\varepsilon$. Then we obtain

$$d\mathcal{Z}_t^\varepsilon(u) = \mathcal{Y}_t(\Delta \Theta_u^\varepsilon) dt + E d\mathcal{A}_t(\Theta_u^\varepsilon) + \sqrt{2} d\nabla_{\text{Dir}} \mathcal{W}_t(\Theta_u^\varepsilon),$$

where \mathcal{A} is the non-linearity that has been previously defined and

$$\nabla_{\text{Dir}} \mathcal{W}_t(\Theta_u^\varepsilon) = -\mathcal{W}_t(\nabla \Theta_u^\varepsilon) = \mathcal{W}_t(p_\varepsilon^{\text{Neu}}(u, \cdot) - 1),$$

where $p_\varepsilon^{\text{Neu}}$ is the heat kernel with Neumann boundary conditions, namely

$$p_\varepsilon^{\text{Neu}}(u, v) = -\nabla_v \Theta_u^\varepsilon(v) + 1 = 1 + \sum_{\ell=1}^{\infty} e^{-\varepsilon \pi^2 \ell^2} 2 \cos(\ell \pi u) \cos(\ell \pi v), \quad (6.16)$$

for $u, v \in [0, 1]$ and $\varepsilon > 0$. To shorten the notation we write

$$J_u^\varepsilon(\cdot) = p_\varepsilon^{\text{Neu}}(u, \cdot) - 1 = -\nabla \Theta_u^\varepsilon(\cdot) \in \mathcal{S}_{\text{Neu}},$$

so that Corollary 6.11 gives $\mathcal{A}_t(\Theta_u^\varepsilon) = \mathcal{B}_t(J_u^\varepsilon) = (\int_0^t \mathcal{Y}_s^{\circ 2} ds)(J_u^\varepsilon)$. Furthermore, note that $\Delta \Theta_u^\varepsilon(v) = \Delta_u \Theta_u^\varepsilon(v)$, from where we get

$$d\mathcal{Z}_t^\varepsilon(u) = \Delta \mathcal{Z}_t^\varepsilon(u) dt + E d\mathcal{B}_t(J_u^\varepsilon) + \sqrt{2} d\mathcal{W}_t(J_u^\varepsilon),$$

and $d\langle \mathcal{Z}^\varepsilon(u) \rangle_t = 2\|J_u^\varepsilon\|_{L^2}^2 dt$ because $\mathcal{B}(J_u^\varepsilon)$ has zero quadratic variation by Lemma 6.12.

Now, let us consider the process $\Psi_t^\varepsilon(u) = e^{E\mathcal{Z}_t^\varepsilon(u)}$. By the Cole-Hopf transform for the KPZ equation we expect that Ψ^ε solves an approximate version of the stochastic heat equation as $\varepsilon \rightarrow 0$, and our goal for the remainder of this section is to prove this. Itô's formula applied to $\Psi_t^\varepsilon(u)$ gives

$$\begin{aligned} d\Psi_t^\varepsilon(u) &= \Psi_t^\varepsilon(u) (E\Delta \mathcal{Z}_t^\varepsilon(u) dt + E^2 d\mathcal{B}_t(J_u^\varepsilon) + E\sqrt{2} d\mathcal{W}_t(J_u^\varepsilon) + E^2 \|J_u^\varepsilon\|_{L^2}^2 dt) \\ &= \Delta \Psi_t^\varepsilon(u) dt + E^2 \Psi_t^\varepsilon(u) (\|J_u^\varepsilon\|_{L^2}^2 - (\mathcal{Y}_t^\varepsilon(u))^2) dt \\ &\quad + E^2 \Psi_t^\varepsilon(u) (d\mathcal{B}_t(p_\varepsilon^{\text{Neu}}(u, \cdot)) - d\langle \mathcal{B}_t, 1 \rangle + E^{-1} \sqrt{2} d\mathcal{W}_t(J_u^\varepsilon)), \end{aligned}$$

where we wrote $\langle \mathcal{B}_t, 1 \rangle$ instead of $\mathcal{B}_t(1)$ to avoid confusion between testing against the constant function 1 and evaluating in the point 1, and where in the second step we applied the chain rule for the Laplacian, namely $\Delta e^{E\mathcal{Z}^\varepsilon} = E e^{E\mathcal{Z}^\varepsilon} \Delta \mathcal{Z}^\varepsilon + E^2 e^{E\mathcal{Z}^\varepsilon} (\nabla \mathcal{Z}^\varepsilon)^2$. Next, recall that $\nabla \mathcal{Z}_t^\varepsilon(0) = \mathcal{Y}_t(p_\varepsilon^{\text{Dir}}(0, \cdot)) = \mathcal{Y}_t(0) = 0$ and similarly $\nabla \mathcal{Z}_t^\varepsilon(1) = 0$, which means that for any $\varphi \in \mathcal{C}^2([0, 1])$ we can apply integration by parts to obtain for all $t \geq 0$

$$\int_0^1 \Delta \Psi_t^\varepsilon(u) \varphi(u) du = \int_0^1 \Psi_t^\varepsilon(u) \Delta \varphi(u) du - (\nabla \varphi(1) \Psi_t^\varepsilon(1) - \nabla \varphi(0) \Psi_t^\varepsilon(0)).$$

We apply this to derive the following weak formulation, for $\varphi \in \mathcal{C}^2([0, 1])$,

$$\begin{aligned} d\Psi_t^\varepsilon(\varphi) &= \Psi_t^\varepsilon(\Delta\varphi)dt + E\sqrt{2} \int_0^1 \Psi_t^\varepsilon(u)\varphi(u)d\mathcal{W}_t(J_u^\varepsilon)du - E^2\Psi_t^\varepsilon(\varphi)d\langle \mathcal{B}_t, 1 \rangle \\ &\quad + E^2 \int_0^1 (K_u^\varepsilon - 1)\Psi_t^\varepsilon(u)\varphi(u)dudt + \frac{E^2}{2}(\Psi_t^\varepsilon(0)\varphi(0) + \Psi_t^\varepsilon(1)\varphi(1))dt \\ &\quad + E^2dR_t^\varepsilon(\varphi) + E^2dQ_t^\varepsilon(\varphi), \end{aligned}$$

where

$$R_t^\varepsilon(\varphi) = \int_0^t \int_0^1 \Psi_s^\varepsilon(u) \left[d\mathcal{B}_s(p_\varepsilon^{\text{Neu}}(u, \cdot)) - ((\mathcal{Y}_s^\varepsilon(u))^2 - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2)ds - K_u^\varepsilon ds \right] \varphi(u)du \quad (6.17)$$

for the deterministic function

$$K_u^\varepsilon = E^2 \left(\int_0^1 p_\varepsilon^{\text{Neu}}(u, w)(\Theta_u^\varepsilon(w))^2 dw - \left(\int_0^1 p_\varepsilon^{\text{Dir}}(u, v)\Theta_u^\varepsilon(v)dv \right)^2 \right), \quad (6.18)$$

and

$$\begin{aligned} Q_t^\varepsilon(\varphi) &= \int_0^t \left[\int_0^1 \Psi_s^\varepsilon(u) \left(\|J_u^\varepsilon\|_{L^2}^2 - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 + 1 - \frac{1}{2}(\delta_0(u) + \delta_1(u)) \right) \varphi(u)du \right. \\ &\quad \left. - E^{-2}(\nabla\varphi(1)\Psi_s^\varepsilon(1) - \nabla\varphi(0)\Psi_s^\varepsilon(0)) \right] ds. \quad (6.19) \end{aligned}$$

In the next section we will show the following results concerning the three additional terms $R_t^\varepsilon(\varphi)$, $Q_t^\varepsilon(\varphi)$ and K_u^ε :

Lemma 6.13. *We have for all $p > 2$ and $\varphi \in \mathcal{C}([0, 1])$*

$$\lim_{\varepsilon \rightarrow 0} \left\{ \mathbb{E}[\|R^\varepsilon(\varphi)\|_{p\text{-var}[0, T]}^2] + \mathbb{E}\left[\sup_{t \in [0, T]} |R_t^\varepsilon(\varphi)|^2\right] \right\} = 0. \quad (6.20)$$

To control the term $Q^\varepsilon(\varphi)$, we need to assume that φ satisfies suitable boundary conditions. But this will not be a problem because the boundary condition is compatible with our formulation of the stochastic heat equation with Robin boundary condition, as defined in Appendix 3.4.

Lemma 6.14. *For all $p > 2$ and $\varphi \in \mathcal{C}^2([0, 1])$ with $\nabla\varphi(0) = \nabla\varphi(1) = 0$ we have almost surely*

$$\lim_{\varepsilon \rightarrow 0} \|Q^\varepsilon(\varphi)\|_{L^\infty([0, T])} = 0,$$

and the 1-variation norm (as defined in (6.8)) of $\{Q^\varepsilon(\varphi)\}_{\varepsilon \in (0, 1]}$ is uniformly bounded.

Remark 6.15. *In the definition of $Q_t^\varepsilon(\varphi)$ we have to subtract the Dirac deltas at 0 and 1 in order to see the convergence to zero. Since these terms are not present in the formula for $d\Psi_t^\varepsilon(\varphi)$ we have to add them back in (times the prefactor $E^2/2$). This is exactly the correction to the boundary condition that we discussed in Remark 3.16.*

Lemma 6.16. *The function K^ε converges in $L^2([0, 1])$ as $\varepsilon \rightarrow 0$ to $E^2/12$.*

Lemma 6.17. *The family of processes $\{\Psi^\varepsilon\}$ converges in $L^2(\Omega \times [0, T] \times [0, 1])$ as $\varepsilon \rightarrow 0$ to $\Psi_t(u) = e^{E\mathcal{A}_t(\Theta_u^0)}$. Moreover, $\Psi \in \mathcal{L}_C([0, T])$ and for all $u \in [0, 1]$ we have $\Psi_t(u) = \lim_{\varepsilon \rightarrow 0} \Psi_t^\varepsilon(u)$ almost surely, where the convergence is in $L^2(\Omega \times [0, T])$.*

Proof. The convergence in $L^2(\Omega \times [0, T] \times [0, 1])$ easily follows from the inequality $|e^x - e^y| \leq (e^x + e^y)|x - y|$ and the $L^2([0, 1])$ -convergence of Θ_u^ε to Θ_u^0 , where the latter holds uniformly in $u \in [0, 1]$. To see that $\Psi \in \mathcal{L}_C([0, T])$ note that we have uniformly in $\varepsilon \in (0, 1]$ and $t \in [0, T]$

$$\begin{aligned} \mathbb{E}[|\Psi_t^\varepsilon(u) - \Psi_t^\varepsilon(u')|^2] &\lesssim \mathbb{E}[|\mathcal{Y}_t(\Theta_u^\varepsilon) - \mathcal{Y}_t(\Theta_{u'}^\varepsilon)|^4]^{1/2} \simeq \int |\Theta_u^\varepsilon(v) - \Theta_{u'}^\varepsilon(v)|^2 dv \\ &\simeq \sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} \ell^{-2} |\cos(\ell\pi u) - \cos(\ell\pi u')|^2 \lesssim \sum_{\ell=1}^{\infty} \ell^{-3/2} |u - u'|^{1/2}. \end{aligned}$$

This uniform continuity in u also gives $\Psi_t(u) = \lim_{\varepsilon \rightarrow 0} \Psi_t^\varepsilon(u)$. \square

Given these lemmas, the same arguments as in the proof of [25, Theorem 2.4] (see more precisely [25, Section 4.1]) show that the process Ψ solves for all $\varphi \in \mathcal{S}_{\text{Neu}}$

$$\begin{aligned} \Psi_t(\varphi) &= \langle e^{E\mathcal{B}_0(\Theta_u^0)}, \varphi \rangle + \int_0^t \Psi_s(\Delta\varphi) ds \\ &\quad + \left(\frac{E^4}{12} - E^2 \right) \int_0^t \Psi_s(\varphi) ds - E^2 \int_0^t \Psi_s(\varphi) d\langle \mathcal{B}_s, 1 \rangle \\ &\quad + \frac{E^2}{2} (\Psi_t(0)\varphi(0) + \Psi_t(1)\varphi(1)) dt \\ &\quad + E\sqrt{2} \left(\int_0^1 \int_0^t \Psi_s(u)\varphi(u) d\mathcal{W}_s(u) du - \int_0^t \Psi_s(\varphi) d\langle \mathcal{W}_s, 1 \rangle \right). \end{aligned}$$

Therefore, setting

$$\mathcal{X}_t = E\sqrt{2}\langle \mathcal{W}_t, 1 \rangle - \frac{E^4}{12}t + E^2\langle \mathcal{B}_t, 1 \rangle \quad \text{and} \quad \Phi_t(u) = \Psi_t(u)e^{\mathcal{X}_t}, \quad (6.21)$$

we have $\Phi_0 = \Psi_0$ and $d\langle \mathcal{X} \rangle_t = E^2 2dt$ (recall that $\langle B, 1 \rangle$ has zero quadratic variation). Moreover, for $\varphi \in \mathcal{S}_{\text{Neu}}$, Itô's formula gives

$$\begin{aligned} d\Phi_t(\varphi) &= e^{\mathcal{X}_t} d\Psi_t(\varphi) + \Phi_t(\varphi) d\mathcal{X}_t + \frac{1}{2} \Phi_t(\varphi) d\langle \mathcal{X} \rangle_t + d\langle \Psi(\varphi), e^{\mathcal{X}} \rangle_t \\ &= \Phi_t(\Delta\varphi) dt + \frac{E^2}{2} (\Phi_t(0)\varphi(0) + \Phi_t(1)\varphi(1)) dt \\ &\quad + E\sqrt{2} \int_0^1 \Phi_t(u)\varphi(u) d\mathcal{W}_t(u) du, \end{aligned}$$

where we used that

$$\begin{aligned} \langle \Psi(\varphi), e^{\mathcal{X}} \rangle_t &= E^2 2 \left\langle \int_0^1 \int_0^t \Psi_s(u)\varphi(u) d\mathcal{W}_s(u) du - \int_0^t \Psi_s(\varphi) d\langle \mathcal{W}_s, 1 \rangle, \int_0^t e^{\mathcal{X}_s} d\langle \mathcal{W}_s, 1 \rangle \right\rangle_t \\ &= E^2 2 \left(\int_0^1 \int_0^t \Phi_s(u)\varphi(u) ds du - \int_0^t \Phi_s(\varphi) ds \right) = 0. \end{aligned}$$

Moreover, Φ is locally uniformly bounded in $L^2(\mathbb{P})$, and more precisely $\Phi \in \mathcal{L}_C^2([0, T])$, as given by the following:

Lemma 6.18. *Consider the process Φ defined in (6.21). There exists $T > 0$ such that*

$$\sup_{t \in [0, T], u \in [0, 1]} \mathbb{E}[\Phi_t(u)^2] < \infty.$$

Moreover, if $R \geq 0$ and $\sup_{t \in [0, R], u \in [0, 1]} \mathbb{E}[\Phi_t(u)^8] < \infty$, then for some $T > 0$, independent of R ,

$$\sup_{t \in [0, R+T], u \in [0, 1]} \mathbb{E}[\Phi_t(u)^2] < \infty.$$

Proof of Lemma 6.18. The proof is basically the same as for [25, Lemma B.1], and therefore we omit it. \square

In other words, $\{\Phi_t; t \in [0, T]\}$ is a weak solution to the stochastic heat equation

$$d\Phi_t = \Delta_{\text{Rob}} \Phi_t dt + E\sqrt{2}\Phi_t d\mathcal{W}_t, \quad \Phi_0(u) = e^{E\mathcal{W}_0(\Theta_u^0)},$$

with Robin boundary conditions

$$\nabla \Phi_t(0) = -\frac{E^2}{2}\Phi_t(0), \quad \nabla \Phi_t(1) = \frac{E^2}{2}\Phi_t(1),$$

as defined in Proposition 3.10. By the uniqueness property given in that proposition, on $[0, T]$ the process Φ is equal to the unique weak solution of (3.11). Then Lemma D.1 stated ahead gives the better moment bound

$$\sup_{t \in [0, T], u \in [0, 1]} \mathbb{E}[\Phi_t(u)^8] < \infty,$$

and now Lemma 6.18 shows that $\sup_{t \in [0, 2T], u \in [0, 1]} \mathbb{E}[\Phi_t(u)^2] < \infty$, which means that $\{\Phi_t; t \in [0, 2T]\}$ is the unique weak solution of (3.11) on $[0, 2T]$. Now we can keep iterating this argument to see that Φ is uniquely determined on all of $[0, \infty)$.

Since for all $\varphi \in \mathcal{S}_{\text{Dir}}$

$$\mathcal{Y}(\varphi) = -E^{-1} \log \Psi(\nabla \varphi) = -E^{-1} (\log \Phi - \mathcal{X})(\nabla \varphi) = -E^{-1} \log(\Phi)(\nabla \varphi), \quad (6.22)$$

which can be easily verified for \mathcal{Y}^ε and Ψ^ε and then carries over to the limit $\varepsilon \rightarrow 0$, the uniqueness of \mathcal{Y} follows from that of Φ .

As in the proof of [25, Theorem 2.10] we also obtain the uniqueness of the almost stationary energy solution to the KPZ equation, i.e. Theorem 3.7.

6.4. Convergence of the remainders. In this section we prove successively Lemma 6.13, Lemma 6.14 and Lemma 6.16.

6.4.1. Proof of Lemma 6.13. Recall that $R^\varepsilon(\varphi)$ was defined in (6.17) and that $\mathcal{B}_t(\psi) = (\int_0^t \mathcal{Y}_s^{\circ 2} ds)(\psi)$. Given $\delta > 0$ we approximate R^ε by

$$\begin{aligned} R_t^{\varepsilon, \delta}(\varphi) &= \int_0^t \int_0^1 \Psi_s^\varepsilon(u) \left\langle (\mathcal{Y}_s^\delta(\cdot))^2 - \|p_\delta^{\text{Dir}}(\cdot, \nu)\|_{L_v^2}^2, p_\varepsilon^{\text{Neu}}(u, \cdot) \right\rangle_{L^2} \varphi(u) du ds \\ &\quad - \int_0^t \int_0^1 \Psi_s^\varepsilon(u) \left[(\mathcal{Y}_s^\varepsilon(u))^2 - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 + K_u^{\varepsilon, \delta} \right] \varphi(u) du ds \\ &=: \int_0^t r^{\varepsilon, \delta}(\mathcal{Y}_s) ds, \end{aligned} \quad (6.23)$$

where $K^{\varepsilon, \delta}$ will be defined in equation (6.25) below. Provided that $K^{\varepsilon, \delta}$ converges to K^ε in $L^2([0, 1])$ as $\delta \rightarrow 0$ we obtain from Lemma 6.12 that

$$\lim_{\delta \rightarrow 0} \left\{ \mathbb{E} \left[\|R^{\varepsilon, \delta}(\varphi) - R^\varepsilon(\varphi)\|_{p\text{-var}; [0, T]}^2 \right] + \mathbb{E} \left[\sup_{t \in [0, T]} |R_t^{\varepsilon, \delta}(\varphi) - R_t^\varepsilon(\varphi)|^2 \right] \right\} = 0,$$

which holds because we can control

$$\int_0^\cdot \Psi_s^\varepsilon(u) \left\langle (\mathcal{Y}_s^\delta)^2(\cdot) - \|p_\delta^{\text{Dir}}(\cdot, v)\|_{L_v^2}^2, p_\varepsilon^{\text{Neu}}(u, \cdot) \right\rangle_{L^2} ds$$

as a Young integral: indeed, by Lemma 6.12 the integrator converges in α -Hölder norm for any $\alpha < \frac{3}{4}$ and the integrand is almost surely β -Hölder continuous for any $\beta < \frac{1}{2}$ and $\alpha + \beta > 1$; see [42] or [34, Section 1.3] for details on the Young integral. So, to prove the convergence claimed in Lemma 6.13, it suffices to show that $R^{\varepsilon, \delta}(\varphi)$ vanishes in the same sense as in (6.20), as first $\delta \rightarrow 0$ and then $\varepsilon \rightarrow 0$.

By Corollary 6.6 it suffices to show $\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \|r^{\varepsilon, \delta}(\cdot)\|_{-1,0} = 0$, where the random variable $r^{\varepsilon, \delta}$ was defined in (6.23). Note that $r^{\varepsilon, \delta}$ satisfies the assumption of Corollary 6.6: it is clearly in $L^2(\mu)$ because all the kernels appearing in its definition are bounded continuous functions. The fact that $r^{\varepsilon, \delta}$ is in $\mathcal{H}_{\text{Dir}}^{-1}$ is not obvious but will be a consequence of our estimates below.

Recall that $\mathcal{Y}(\psi_1)\mathcal{Y}(\psi_2) = W_2(\psi_1 \otimes \psi_2)(\mathcal{Y}) + \langle \psi_1, \psi_2 \rangle_{L^2}$ for all $\psi_1, \psi_2 \in L^2$, see [35, Proposition 1.1.2], and therefore

$$\left\langle (\mathcal{Y}^\delta(\cdot))^2 - \|p_\delta^{\text{Dir}}(\cdot, v)\|_{L_v^2}^2, p_\varepsilon^{\text{Neu}}(u, \cdot) \right\rangle_{L^2} - (\mathcal{Y}^\varepsilon)^2(u) + \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 = W_2(g_u^{\varepsilon, \delta})(\mathcal{Y})$$

with

$$g_u^{\varepsilon, \delta}(v_1, v_2) = \int_0^1 p_\delta^{\text{Dir}}(w, v_1) p_\delta^{\text{Dir}}(w, v_2) p_\varepsilon^{\text{Neu}}(u, w) dw - p_\varepsilon^{\text{Dir}}(u, v_1) p_\varepsilon^{\text{Dir}}(u, v_2). \quad (6.24)$$

To control the \mathfrak{H}_0^{-1} -norm of $r^{\varepsilon, \delta}(\cdot)$ we need to bound $\mathbb{E}[r^{\varepsilon, \delta}(\mathcal{Y}_0)F(\mathcal{Y}_0)]$ for an arbitrary cylinder function $F \in \mathcal{C}$. As in [25, Lemma 4.4] we use Gaussian (partial) integration by parts to show that, with the choice

$$K_u^{\varepsilon, \delta} = E^2 \int_{[0,1]^2} g_u^{\varepsilon, \delta}(v_1, v_2) \Theta_u^\varepsilon(v_1) \Theta_u^\varepsilon(v_2) dv_1 dv_2, \quad (6.25)$$

we have the following result:

Lemma 6.19. *We can bound*

$$\begin{aligned} \|r^{\varepsilon, \delta}(\cdot)\|_{-1,0}^2 &\leq \mathbb{E} \left[\left\| \int_0^1 du \varphi(u) (W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^\varepsilon(u)) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)}^2 \right] \\ &\quad + \mathbb{E} \left[\left\| \int_0^1 du \varphi(u) \int_0^1 dv_1 g_u^{\varepsilon, \delta}(v_1, v_2) D_{v_1} \Psi_0^\varepsilon(u) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}(v_2)}^2 \right] =: A^{\varepsilon, \delta} + B^{\varepsilon, \delta} \end{aligned}$$

where the \diamond -product is defined below:

$$W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^\varepsilon(u) := W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \Psi_0^\varepsilon(u) - \int_0^1 g_u^{\varepsilon, \delta}(v_1, v_2) D_{v_2} \Psi_0^\varepsilon(u) dv_2.$$

To prove Lemma 6.19, we control separately $A^{\varepsilon, \delta}$ and $B^{\varepsilon, \delta}$ and prove that they vanish respectively in Lemma 6.23 and Lemma 6.22 below. Before that, we need some auxiliary results that we give now:

Lemma 6.20. *The sequence $\|g_u^{\varepsilon, \delta}(\cdot, v_2)\|_{\mathcal{H}_{\text{Dir}}^{-1}}$ is uniformly bounded in $(\delta, u, v_2) \in (0, 1] \times [0, 1]^2$, and we have for almost all (w.r.t. Lebesgue) $v_2 \in [0, 1]$, and all $\varepsilon \in (0, 1]$*

$$\lim_{\delta \rightarrow 0} \|g_u^{\varepsilon, \delta}(\cdot, v_2)\|_{\mathcal{H}_{\text{Dir}}^{-1}} \lesssim |p_\varepsilon^{\text{Neu}}(u, v_2) - p_\varepsilon^{\text{Dir}}(u, v_2)| + p_\varepsilon^{\text{Dir}}(u, v_2) \times (\varepsilon^{-1/4}|v_2 - u| + \varepsilon^{1/4}).$$

Proof of Lemma 6.20. Recall that $g_u^{\varepsilon, \delta}(v_1, v_2)$ is given in (6.24). We use the explicit characterization of the $\mathcal{H}_{\text{Dir}}^{-1}$ -norm in Lemma 6.8 to write

$$\|g_u^{\varepsilon, \delta}(\cdot, v_2)\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 = \pi^{-2} \sum_{k=1}^{\infty} k^{-2} |\langle g_u^{\varepsilon, \delta}(\cdot, v_2), \sin(k\pi \cdot) \rangle_{L^2}|^2,$$

and since $\int_0^1 p_t^{\text{Dir}}(u, v_1) \sin(k\pi v_1) dv_1 = e^{-t\pi^2 k^2} \sin(k\pi u)$ for $k \in \mathbb{N}$, $t > 0$, $u \in [0, 1]$, we have

$$\begin{aligned} & |\langle g_u^{\varepsilon, \delta}(\cdot, v_2), \sin(k\pi \cdot) \rangle_{L^2([0,1])}| \\ &= \left| \int_0^1 e^{-\delta\pi^2 k^2} p_{\delta}^{\text{Dir}}(w, v_2) \sin(k\pi w) p_{\varepsilon}^{\text{Neu}}(u, w) dw - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u) p_{\varepsilon}^{\text{Dir}}(u, v_2) \right|, \end{aligned}$$

from where the uniform bound follows (note that we made no claim about uniform boundedness in ε). Moreover, since p^{Dir} are the transition densities of a killed Brownian motion we have for all $v_2 \in (0, 1)$

$$\begin{aligned} & \lim_{\delta \rightarrow 0} \left| \int_0^1 e^{-\delta\pi^2 k^2} p_{\delta}^{\text{Dir}}(w, v_2) \sin(k\pi w) p_{\varepsilon}^{\text{Neu}}(u, w) dw - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u) p_{\varepsilon}^{\text{Dir}}(u, v_2) \right| \\ &= \left| \sin(k\pi v_2) p_{\varepsilon}^{\text{Neu}}(u, v_2) - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u) p_{\varepsilon}^{\text{Dir}}(u, v_2) \right| \\ &\leq |p_{\varepsilon}^{\text{Neu}}(u, v_2) - p_{\varepsilon}^{\text{Dir}}(u, v_2)| + p_{\varepsilon}^{\text{Dir}}(u, v_2) |\sin(k\pi v_2) - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u)|. \end{aligned}$$

The last factor on the right hand side is bounded by

$$\begin{aligned} |\sin(k\pi v_2) - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u)| &\leq |1 - e^{-\varepsilon\pi^2 k^2}| + |\sin(k\pi v_2) - \sin(k\pi u)| \\ &\lesssim (\varepsilon\pi^2 k^2)^{1/2} + k|v_2 - u| \lesssim k(|v_2 - u| + \varepsilon^{1/2}), \end{aligned}$$

where we applied the bound $|1 - e^{-c}| \leq c^{1/2}$, true for any $c > 0$. We use this bound for $k \leq \varepsilon^{-1/2}$, while for $k > \varepsilon^{-1/2}$ we simply bound $|\sin(k\pi v_2) - e^{-\varepsilon\pi^2 k^2} \sin(k\pi u)| \leq 2$. Thus, we obtain

$$\begin{aligned} \lim_{\delta \rightarrow 0} \|g_u^{\varepsilon, \delta}(\cdot, v_2)\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 &\lesssim \sum_{k=1}^{\infty} k^{-2} \left(|p_{\varepsilon}^{\text{Neu}}(u, v_2) - p_{\varepsilon}^{\text{Dir}}(u, v_2)|^2 \right. \\ &\quad \left. + |p_{\varepsilon}^{\text{Dir}}(u, v_2)|^2 \times [\mathbf{1}_{(0, \varepsilon^{-1/2}]}(k) (k^2(|v_2 - u|^2 + \varepsilon)) + \mathbf{1}_{(\varepsilon^{-1/2}, +\infty)}(k)] \right) \\ &\lesssim \left(|p_{\varepsilon}^{\text{Neu}}(u, v_2) - p_{\varepsilon}^{\text{Dir}}(u, v_2)|^2 \right. \\ &\quad \left. + |p_{\varepsilon}^{\text{Dir}}(u, v_2)|^2 \times [\varepsilon^{-1/2}|v_2 - u|^2 + \varepsilon^{1/2}] \right), \end{aligned}$$

and now it suffices to take the square root on both sides. \square

The following lemma, which is a simple consequence of the estimates we derived so far, will be useful for bounding both constants $A^{\varepsilon, \delta}$ and $B^{\varepsilon, \delta}$ of Lemma 6.19:

Lemma 6.21. *For all $u' \in [0, 1]$ the function $x \mapsto \left\| \int_0^1 g_u^{\varepsilon, \delta}(v_1, \cdot) \Theta_{u'}^{\varepsilon}(v_1) dv_1 \right\|_{\mathcal{H}_{\text{Dir}}^{-1}}$ is uniformly bounded in $(\delta, u) \in (0, 1] \times [0, 1]$ and satisfies, for all $\varepsilon \in (0, 1]$*

$$\sup_{u' \in [0, 1]} \lim_{\delta \rightarrow 0} \left\| \int_0^1 g_u^{\varepsilon, \delta}(v_1, \cdot) \Theta_{u'}^{\varepsilon}(v_1) dv_1 \right\|_{\mathcal{H}_{\text{Dir}}^{-1}} \lesssim \frac{\varepsilon^{1/2}}{u(1-u)} \wedge 1 + \varepsilon^{1/4}.$$

Proof of Lemma 6.21. From Lemma 6.20 (note that $g_u^{\varepsilon, \delta}$ is symmetric in its two arguments) we obtain

$$\begin{aligned}
\lim_{\delta \rightarrow 0} \left\| \int_0^1 g_u^{\varepsilon, \delta}(v_1, \cdot) \Theta_{u'}^{\varepsilon}(v_1) dv_1 \right\|_{\mathcal{H}_{\text{Dir}}^{-1}} &\leq \lim_{\delta \rightarrow 0} \int_0^1 \|g_u^{\varepsilon, \delta}(v_1, \cdot)\|_{\mathcal{H}_{\text{Dir}}^{-1}} |\Theta_{u'}^{\varepsilon}(v_1)| dv_1 \\
&\lesssim \int_0^1 \left(|p_{\varepsilon}^{\text{Neu}}(u, v_1) - p_{\varepsilon}^{\text{Dir}}(u, v_1)| + p_{\varepsilon}^{\text{Dir}}(u, v_1) \times (\varepsilon^{-1/4} |v_1 - u| + \varepsilon^{1/4}) \right) |\Theta_{u'}^{\varepsilon}(v_1)| dv_1 \\
&\leq \|\Theta_{u'}^{\varepsilon}\|_{L^{\infty}} \left(\|p_{\varepsilon}^{\text{Neu}}(u, \cdot) - p_{\varepsilon}^{\text{Dir}}(u, \cdot)\|_{L^1} + \varepsilon^{-1/4} \|p_{\varepsilon}^{\text{Dir}}(u, \cdot) \times |\cdot - u|\|_{L^1} + \varepsilon^{1/4} \right) \\
&\lesssim \frac{\varepsilon^{1/2}}{u(1-u)} \wedge 1 + \varepsilon^{1/4}, \quad \text{for any } u' \in [0, 1],
\end{aligned}$$

where the last estimate comes from Lemma C.1 and Lemma C.2 of the Appendix. \square

Finally we are now able to state and prove that $A^{\varepsilon, \delta}$ and $B^{\varepsilon, \delta}$ vanish:

Lemma 6.22. *We have*

$$\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} B^{\varepsilon, \delta} = 0.$$

Proof of Lemma 6.22. Recall that $\Psi_0^{\varepsilon}(u) = e^{E\mathcal{D}_0^{\varepsilon}(u)} = e^{E\mathcal{D}_0(\Theta_u^{\varepsilon})}$ and therefore its Malliavin derivative is $D_{v_1} \Psi_0^{\varepsilon}(u) = E\Psi_0^{\varepsilon}(u) \Theta_u^{\varepsilon}(v_1)$. As a result,

$$\begin{aligned}
B^{\varepsilon, \delta} &= \mathbb{E} \left[\left\| \int_0^1 du \varphi(u) \int_0^1 dv_1 g_u^{\varepsilon, \delta}(v_1, \cdot) E\Theta_u^{\varepsilon}(v_1) \Psi_0^{\varepsilon}(u) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 \right] \\
&\stackrel{(C-S)}{\leq} \int_0^1 du |\varphi(u)|^2 E^2 \left\| \int_0^1 dv_1 g_u^{\varepsilon, \delta}(v_1, \cdot) \Theta_u^{\varepsilon}(v_1) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 \mathbb{E}[|\Psi_0^{\varepsilon}(u)|^2],
\end{aligned}$$

so from Lemma 6.21 together with the dominated convergence theorem we get

$$\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} B^{\varepsilon, \delta} \lesssim \lim_{\varepsilon \rightarrow 0} \int_0^1 du |\varphi(u)|^2 \left(\left(\frac{\varepsilon^{1/2}}{u(1-u)} \wedge 1 \right) + \varepsilon^{1/4} \right)^2 \mathbb{E}[|\Psi_0^{\varepsilon}(u)|^2] = 0.$$

\square

To complete the proof of Lemma 6.13 we need to control $A^{\varepsilon, \delta}$, which is achieved in the next lemma.

Lemma 6.23. *We have*

$$\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} A^{\varepsilon, \delta} = 0.$$

Proof of Lemma 6.23. We expand

$$\begin{aligned}
A^{\varepsilon, \delta} &= \mathbb{E} \left[\left\| \int_0^1 du \varphi(u) (W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^{\varepsilon}(u)) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)}^2 \right] \\
&= \int_0^1 du \int_0^1 du' \varphi(u) \varphi(u') \\
&\quad \times \mathbb{E} \left[\left\langle W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^{\varepsilon}(u), W_1(g_{u'}^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^{\varepsilon}(u') \right\rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} \right], \quad (6.26)
\end{aligned}$$

and as in the proof of [25, Lemma 4.7] we get from an integration by parts

$$\mathbb{E} \left[\left\langle W_1(g_u^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^{\varepsilon}(u), W_1(g_{u'}^{\varepsilon, \delta}(v_1, \cdot)) \diamond \Psi_0^{\varepsilon}(u') \right\rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} \right]$$

$$\begin{aligned}
&= \mathbb{E}[\Psi_0^\varepsilon(u)\Psi_0^\varepsilon(u')] \int_0^1 \langle g_u^{\varepsilon,\delta}(v_1, v_2), g_{u'}^{\varepsilon,\delta}(v_1, v_2) \rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} dv_2 \\
&\quad + E^2 \mathbb{E}[\Psi_0^\varepsilon(u)\Psi_0^\varepsilon(u')] \\
&\quad \times \int_0^1 dv_2 \int_0^1 dv_3 \Theta_u^\varepsilon(v_3) \Theta_{u'}^\varepsilon(v_2) \langle g_u^{\varepsilon,\delta}(v_1, v_2), g_{u'}^{\varepsilon,\delta}(v_1, v_3) \rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)}. \quad (6.27)
\end{aligned}$$

For the second term on the right hand side we simply estimate from the Cauchy-Schwartz inequality

$$\begin{aligned}
&\lim_{\delta \rightarrow 0} \left| \int_0^1 dv_2 \int_0^1 dv_3 \Theta_u^\varepsilon(v_3) \Theta_{u'}^\varepsilon(v_2) \langle g_u^{\varepsilon,\delta}(v_1, v_2), g_{u'}^{\varepsilon,\delta}(v_1, v_3) \rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} \right| \\
&= \lim_{\delta \rightarrow 0} \left| \left\langle \int_0^1 dv_2 \Theta_{u'}^\varepsilon(v_2) g_u^{\varepsilon,\delta}(v_1, v_2), \int_0^1 dv_3 \Theta_u^\varepsilon(v_3) g_{u'}^{\varepsilon,\delta}(v_1, v_3) \right\rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} \right| \\
&\leq \lim_{\delta \rightarrow 0} \left\| \int_0^1 dv_2 \Theta_{u'}^\varepsilon(v_2) g_u^{\varepsilon,\delta}(\cdot, v_2) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}} \left\| \int_0^1 dv_3 \Theta_u^\varepsilon(v_3) g_{u'}^{\varepsilon,\delta}(\cdot, v_3) \right\|_{\mathcal{H}_{\text{Dir}}^{-1}} \\
&\lesssim \left(\frac{\varepsilon^{1/2}}{u(1-u)} \wedge 1 + \varepsilon^{1/4} \right) \left(\frac{\varepsilon^{1/2}}{u'(1-u')} \wedge 1 + \varepsilon^{1/4} \right),
\end{aligned}$$

where in the last step we used Lemma 6.21 and the symmetry of $g_u^{\varepsilon,\delta}$. So when we plug this contribution into (6.26) inside the integration w.r.t. u and u' , one easily shows that it vanishes for $\varepsilon \rightarrow 0$. We are left with bounding the first contribution coming from (6.27). We set $V^\varepsilon(u) = |\varphi(u)| \mathbb{E}[\Psi_0^\varepsilon(u)^2]^{1/2}$ and obtain from Lemma 6.20 and Lemma C.2 as well as Lemma C.1:

$$\begin{aligned}
&\lim_{\delta \rightarrow 0} \int_0^1 du \int_0^1 du' \varphi(u) \varphi(u') \mathbb{E}[\Psi_0^\varepsilon(u) \Psi_0^\varepsilon(u')] \int_0^1 \langle g_u^{\varepsilon,\delta}(v_1, v_2), g_{u'}^{\varepsilon,\delta}(v_1, v_2) \rangle_{\mathcal{H}_{\text{Dir}}^{-1}(v_1)} dv_2 \\
&\stackrel{(C-S)}{\leq} \lim_{\delta \rightarrow 0} \int_0^1 dv_2 \left(\int_0^1 du V^\varepsilon(u) \|g_u^{\varepsilon,\delta}(\cdot, v_2)\|_{\mathcal{H}_{\text{Dir}}^{-1}} \right)^2 \\
&\stackrel{(\text{Lem. 6.20})}{\lesssim} \int_0^1 dv_2 \left(\int_0^1 du V^\varepsilon(u) [p_\varepsilon^{\text{Neu}}(u, v_2) - p_\varepsilon^{\text{Dir}}(u, v_2)] \right. \\
&\quad \left. + p_\varepsilon^{\text{Dir}}(u, v_2) \times (\varepsilon^{-1/4} |v_2 - u| + \varepsilon^{1/4}) \right]^2 \\
&\lesssim \int_0^1 dv_2 \left(\|V^\varepsilon\|_{L^\infty} \left[\left(\frac{\varepsilon^{1/2}}{v_2(1-v_2)} \wedge 1 \right) + \varepsilon^{1/4} \right] \right)^2 \lesssim \varepsilon^{1/2},
\end{aligned}$$

which also vanishes as $\varepsilon \rightarrow 0$. \square

This concludes the proof of Lemma 6.13.

6.4.2. *Proof of Lemma 6.14.* Recall that $Q^\varepsilon(\varphi)$ was defined in (6.19). One term that appears in its definition is

$$\|J_u^\varepsilon\|_{L^2}^2 - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 = \sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} 2(\cos(\ell\pi u)^2 - \sin(\ell\pi u)^2) = \sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} 2\cos(2\ell\pi u),$$

and for $u \in [0, \frac{1}{2}]$ we have

$$\sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} 2\cos(2\ell\pi u) = p_{2\varepsilon}^{\text{Neu}}(0, 2u) - 1,$$

while for $u \in (\frac{1}{2}, 1]$

$$\begin{aligned} \sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} 2\cos\left(2\ell\pi\left(\frac{1}{2} + (u - \frac{1}{2})\right)\right) &= \sum_{\ell=1}^{\infty} e^{-2\varepsilon\pi^2\ell^2} 2\cos\left(\ell\pi\left[2 - (1 - 2(u - \frac{1}{2}))\right]\right) \\ &= p_{2\varepsilon}^{\text{Neu}}(0, 2 - 2u) - 1. \end{aligned}$$

If $\{f^\varepsilon\}_{\varepsilon>0}$ is a uniformly bounded family of continuous functions that converges to a continuous function f , then

$$\lim_{\varepsilon \rightarrow 0} \left[\int_0^{1/2} p_{2\varepsilon}^{\text{Neu}}(0, 2u) f^\varepsilon(u) du + \int_{1/2}^1 p_{2\varepsilon}^{\text{Neu}}(0, 2 - 2u) f^\varepsilon(u) du \right] = \frac{1}{2}(f(0) + f(1)),$$

which means that for $\varphi \in \mathcal{C}^1([0, 1])$

$$\begin{aligned} Q_t^\varepsilon(\varphi) &= \int_0^t \left[\int_0^1 \Psi_s^\varepsilon(u) \left(\|J_u^\varepsilon\|_{L^2}^2 - \|p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^2}^2 + 1 \right) \varphi(u) du \right. \\ &\quad \left. - \frac{1}{2} (\Psi_s^\varepsilon(0)\varphi(0) + \Psi_s^\varepsilon(1)\varphi(1)) \right. \\ &\quad \left. - E^{-2} (\nabla\varphi(1)\Psi_s^\varepsilon(1) - \nabla\varphi(0)\Psi_s^\varepsilon(0)) \right] ds \\ &\xrightarrow{\varepsilon \rightarrow 0} \int_0^t \left[-E^{-2} (\nabla\varphi(1)\Psi_s(1) - \nabla\varphi(0)\Psi_s(0)) \right] ds, \end{aligned}$$

and if $\nabla\varphi(0) = \nabla\varphi(1) = 0$, the right hand side vanishes. Moreover, since the integrand in the time integral in the definition of $Q^\varepsilon(\varphi)$ converges absolutely, then $\{Q^\varepsilon(\varphi)\}_{\varepsilon \in (0, 1]}$ is uniformly bounded in 1-variation norm.

6.4.3. *Proof of Lemma 6.16.* This is a consequence of the following result, which uses the approximation $K_u^{\varepsilon, \delta}$ defined in (6.25).

Lemma 6.24. *We have*

$$\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} K_u^{\varepsilon, \delta} = \frac{E^2}{12},$$

where the convergence is in $L_u^2([0, 1])$.

Proof of Lemma 6.24. Recall that, by definition

$$\begin{aligned} \lim_{\delta \rightarrow 0} E^{-2} K_u^{\varepsilon, \delta} &= \lim_{\delta \rightarrow 0} \int_{[0, 1]^2} g_u^{\varepsilon, \delta}(v_1, v_2) \Theta_u^\varepsilon(v_1) \Theta_u^\varepsilon(v_2) dv_1 dv_2 \\ &= \lim_{\delta \rightarrow 0} \int_{[0, 1]^3} p_\delta^{\text{Dir}}(w, v_1) p_\delta^{\text{Dir}}(w, v_2) p_\varepsilon^{\text{Neu}}(u, w) \Theta_u^\varepsilon(v_1) \Theta_u^\varepsilon(v_2) dw dv_1 dv_2 \\ &\quad - \int_{[0, 1]^2} p_\varepsilon^{\text{Dir}}(u, v_1) p_\varepsilon^{\text{Dir}}(u, v_2) \Theta_u^\varepsilon(v_1) \Theta_u^\varepsilon(v_2) dv_1 dv_2 \\ &= \int_0^1 p_\varepsilon^{\text{Neu}}(u, w) (\Theta_u^\varepsilon(w))^2 dw - \left(\int_0^1 p_\varepsilon^{\text{Dir}}(u, v_1) \Theta_u^\varepsilon(v_1) dv_1 \right)^2, \end{aligned}$$

where the convergence is pointwise in u , with a uniform bound in δ , so in particular in $L_u^2([0, 1])$. Since $-\nabla\Theta_u^\varepsilon(w) = p_\varepsilon^{\text{Neu}}(u, w) - 1$, the first term in the last line is

$$\begin{aligned} \int_0^1 p_\varepsilon^{\text{Neu}}(u, w)(\Theta_u^\varepsilon(w))^2 dw &= -\frac{1}{3} \int_0^1 \nabla(\Theta_u^\varepsilon(w))^3 dw + \int_0^1 (\Theta_u^\varepsilon(w))^2 dw \\ &= \int_0^1 (\Theta_u^\varepsilon(w))^2 dw, \end{aligned}$$

where we used that $\Theta_u^\varepsilon(0) = \Theta_u^\varepsilon(1) = 0$ because $\Theta_u^\varepsilon \in \mathcal{S}_{\text{Dir}}$. For the remaining term we get

$$\int_0^1 p_\varepsilon^{\text{Dir}}(u, v_1) \Theta_u^\varepsilon(v_1) dv_1 = \sum_{\ell=1}^{\infty} (e^{-\varepsilon \pi^2 \ell^2})^2 (-\ell \pi)^{-1} 2 \sin(\ell \pi u) \cos(\ell \pi u) = \Theta_u^{2\varepsilon}(u).$$

In other words, we obtain $\lim_{\delta \rightarrow 0} E^{-2} K_u^{\varepsilon, \delta} = \|\Theta_u^\varepsilon\|_{L^2}^2 - |\Theta_u^{2\varepsilon}(u)|^2$. Now note that $\Theta_u^\varepsilon(v)$ also writes as

$$\Theta_u^\varepsilon(v) = \int_0^1 \Theta_u(w) p_\varepsilon^{\text{Dir}}(v, w) dw,$$

where $\Theta_u(w)$ is the unique kernel satisfying $\nabla_u \langle \Theta_u, f \rangle_{L^2} = f(u)$ and $\int_0^1 \langle \Theta_u, f \rangle_{L^2} du = 0$ for all $f \in \mathcal{C}([0, 1])$. From the first condition we get $\Theta_u(w) = \mathbf{1}_{[0, u]}(w) + g(w)$ for some g , the second condition gives for almost all w

$$0 = \int_0^1 \Theta_u(w) du = \int_0^1 \mathbf{1}_{[w, 1]}(u) du + g(w) = 1 - w + g(w),$$

so $g(w) = w - 1$. Therefore,

$$\Theta_u^\varepsilon(u) = \int_0^1 \Theta_u(w) p_\varepsilon^{\text{Dir}}(u, w) dw = \mathbb{E}_u^{\text{Dir}}[\mathbf{1}_{[0, u]}(B_{2\varepsilon}) + B_{2\varepsilon} - 1],$$

where $\mathbb{E}_u^{\text{Dir}}$ is the measure under which B is a Brownian motion started at u , killed when reaching 0 or 1, and consequently $\lim_{\varepsilon \rightarrow 0} \Theta_u^\varepsilon(u) = \frac{1}{2} + u - 1 = u - \frac{1}{2}$ for all $u \in (0, 1)$. Also,

$$\lim_{\varepsilon \rightarrow 0} \|\Theta_u^\varepsilon\|_{L^2}^2 = \|\Theta_u\|_{L^2}^2 = \int_0^1 (\mathbf{1}_{[0, u]}(w) + w - 1)^2 dw = u^2 - u + \frac{1}{3},$$

which leads to

$$\lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} E^{-2} K_u^{\varepsilon, \delta} = \|\Theta_u^\varepsilon\|_{L^2}^2 - |\Theta_u^{2\varepsilon}(u)|^2 = u^2 - u + \frac{1}{3} - \left(u - \frac{1}{2}\right)^2 = \frac{1}{12}.$$

□

APPENDIX A. AUXILIARY COMPUTATIONS INVOLVING THE GENERATOR

A.1. Martingale decomposition for the density fluctuation field \mathcal{Y}_t^n . In this section we provide the technical computations that we need to prove (4.3), assuming $\rho = \frac{1}{2}$. We consider the general case $\gamma \geq \frac{1}{2}$, and we aim at computing $\int_0^t n^2 \mathcal{L}_n \mathcal{Y}_s^n(\varphi) ds$ when $\varphi \in \mathcal{S}_{\text{Dir}}$.

First, we note that for any $x \in \Lambda_n$, $\mathcal{L}_n \eta(x) = j_{x-1, x}(\eta) - j_{x, x+1}(\eta)$, where the local function $j_{x, x+1}$ is the microscopic current of the system, which can be decomposed into its symmetric and antisymmetric parts as

$$j_{x, x+1}(\eta) = j_{x, x+1}^s(\eta) + j_{x, x+1}^a(\eta)$$

with

$$j_{x,x+1}^s(\eta) = \eta(x) - \eta(x+1), \quad j_{x,x+1}^a(\eta) = \frac{E}{n^\gamma} \eta(x)(1 - \eta(x+1)).$$

Recall that above, as well as everywhere it appears, we assume by convention $\eta(0) = \eta(n) = \rho$.

We start by looking at the action of the symmetric part of the current in the density field. In the following, for $\varphi \in \mathcal{S}_{\text{Dir}}$, and $x \in \Lambda_n$, we denote $\varphi_x = \varphi(\frac{x}{n})$. A simple computation, which makes use of the property of the test function φ at the boundary, namely $\varphi_0 = \varphi_n = 0$, shows that

$$\begin{aligned} \frac{n^2}{\sqrt{n}} \sum_{x=1}^{n-1} \varphi_x \{j_{x-1,x}^s(\eta) - j_{x,x+1}^s(\eta)\} \\ = \frac{1}{\sqrt{n}} \sum_{x=1}^{n-1} \Delta_n \varphi_x \eta(x) + \frac{n^2}{\sqrt{n}} \varphi_1 \eta(0) + \frac{n^2}{\sqrt{n}} \varphi_{n-1} \eta(n). \end{aligned} \quad (\text{A.1})$$

Since we need to close the previous equation in terms of the density field, we note that the identity $\sum_{x=1}^{n-1} \Delta_n \varphi_x = -n^2(\varphi_1 + \varphi_{n-1})$ implies that:

$$\frac{1}{\sqrt{n}} \sum_{x=1}^{n-1} \Delta_n \varphi_x \bar{\eta}(x) = \frac{1}{\sqrt{n}} \sum_{x=1}^{n-1} \Delta_n \varphi_x (\eta(x) - \rho) = \mathcal{Y}^n(\Delta_n \varphi). \quad (\text{A.2})$$

Now we look at the action of the antisymmetric part of the current in the density field. We use the fact that, when $\rho = \frac{1}{2}$, for any $x \in \Lambda_n$ we have the identity

$$\eta(x-1)(1 - \eta(x)) = -\bar{\eta}(x)\bar{\eta}(x-1) + \frac{1}{2}(\eta(x-1) - \eta(x)) + \frac{1}{4}. \quad (\text{A.3})$$

A simple computation shows that

$$\begin{aligned} \frac{n^2}{\sqrt{n}} \sum_{x=1}^{n-1} \varphi_x \{j_{x-1,x}^a(\eta) - j_{x,x+1}^a(\eta)\} \\ = -\frac{1\sqrt{n}}{n^\gamma} \sum_{x=1}^{n-2} \nabla_n^+ \varphi_x \bar{\eta}(x)\bar{\eta}(x+1) + \frac{E\sqrt{n}}{2n^\gamma} \sum_{x=0}^{n-1} \nabla_n^+ \varphi_x (\eta(x) - \eta(x+1)) + \text{BC}_n(\varphi), \end{aligned} \quad (\text{A.4})$$

where $\text{BC}_n(\varphi)$ is a boundary term that is given by

$$\text{BC}_n(\varphi) = \frac{En^{3/2}}{2n^\gamma} \left(-\varphi_n \bar{\eta}(n-1) - \varphi_0 \bar{\eta}(1) \right)$$

Recall that $\varphi_n = \varphi_0 = 0$. Therefore, (A.4) rewrites as

$$\begin{aligned} -E \frac{\sqrt{n}}{n^\gamma} \sum_{x=1}^{n-2} \nabla_n^+ \varphi_x \bar{\eta}(x)\bar{\eta}(x+1) + \frac{E}{2n^{\gamma+1/2}} \sum_{x=1}^{n-1} \Delta_n \varphi_x \bar{\eta}(x) \\ = -E \frac{\sqrt{n}}{n^\gamma} \sum_{x=1}^{n-2} \nabla_n^+ \varphi_x \bar{\eta}(x)\bar{\eta}(x+1) + \frac{E}{2n^\gamma} \mathcal{Y}^n(\Delta_n \varphi). \end{aligned} \quad (\text{A.5})$$

From (A.2) and (A.5) we easily deduce (4.3).

A.2. Martingale decomposition for the height fluctuation field \mathcal{Z}_t^n . Now let us show (4.23) in a very similar way as we did in Appendix A.1. Recall that $h(x) = h(1) + \sum_{y=1}^{x-1} \bar{\eta}(y)$, as explained in Section 2.4.2. As before we always use the convention $\eta(0) = \eta(n) = \rho$ and we also conveniently set $h(0) = h(1)$ and $h(n+1) = h(n)$. One easily obtains, for any $x \in \{1, \dots, n\}$:

$$\begin{aligned} \mathcal{L}_n^\otimes h(x) &= -j_{x-1,x}(\eta) = \eta(x)(1 - \eta(x-1)) - \left(1 + \frac{E}{n^\gamma}\right) \eta(x-1)(1 - \eta(x)) \\ &= \Delta h(x) - \frac{E}{n^\gamma} \eta(x-1)(1 - \eta(x)), \end{aligned} \quad (\text{A.6})$$

where we wrote $\Delta h(x) = h(x+1) + h(x-1) - 2h(x)$. Note that with our convention $\Delta h(1) = \bar{\eta}(1)$ and $\Delta h(n) = -\bar{\eta}(n-1)$. In (A.6) we want to rewrite the second term in terms of the height configuration values. This can be easily done for $\rho = \frac{1}{2}$ using the identity (A.3), which we rewrite here as:

$$-\eta(x-1)(1 - \eta(x)) = \nabla^- h(x) \nabla^+ h(x) + \frac{1}{2} \Delta h(x) - \frac{1}{4},$$

where $\nabla^- h(x) = h(x) - h(x-1)$ and $\nabla^+ h(x) = h(x+1) - h(x)$. Note that the previous identity does hold for $x = 1$ and $x = n$ since with our convention, $\nabla^- h(1) = 0$ and $\nabla^+ h(n) = 0$. This implies

$$\mathcal{L}_n^\otimes h(x) = \left(1 + \frac{E}{2n^\gamma}\right) \Delta h(x) + \frac{E}{n^\gamma} \left(\nabla^- h(x) \nabla^+ h(x) - \frac{1}{4}\right), \quad \text{for any } x \in \{1, \dots, n\}. \quad (\text{A.7})$$

Now, let us take $\varphi \in \mathcal{S}_{\text{Neu}}$ and as before denote $\varphi_x = \varphi(\frac{x}{n})$. We start by treating the first term in (A.7), and more precisely $\Delta h(x)$. Since our goal is to see the height fluctuation field appear, as long as $h(x)$ is replaced with the time dependent configuration $h_{sn^2}^n(x)$, we have to recenter the heights as follows: using the fact that $\Delta c = 0$ for any constant c , a simple computation (discrete integration by parts) shows that

$$\frac{n^2}{n^{3/2}} \sum_{x=1}^n \varphi_x \Delta(h_{sn^2}^n(x) - c_n s) = \frac{1}{n^{3/2}} \sum_{x=1}^n \tilde{\Delta}_n \varphi_x (h_{sn^2}^n(x) - c_n s) \quad (\text{A.8})$$

$$+ \sqrt{n}(\varphi_1 - \varphi_0)(h_{sn^2}^n(1) - c_n s) \quad (\text{A.9})$$

$$+ \sqrt{n}(\varphi_n - \varphi_{n-1})(h_{sn^2}^n(n) - c_n s), \quad (\text{A.10})$$

where $\tilde{\Delta}_n \varphi_x$ has been defined in (4.25) as:

$$\tilde{\Delta}_n \varphi_x = \begin{cases} n^2(\varphi_{x+1} + \varphi_{x-1} - 2\varphi_x) = \Delta_n \varphi_x & \text{if } x \in \{1, \dots, n-1\}, \\ 2n^2(\varphi_{n-1} - \varphi_n) & \text{if } x = n. \end{cases}$$

Note that the term (A.8) corresponds to $\mathcal{Z}_s^n(\tilde{\Delta}_n \varphi)$, while the term (A.9) and (A.10), when integrated in time between 0 and t , will give the contribution $\mathcal{R}_t^n(\varphi)$ (see (4.24)). We are left with the second term in (A.7), that we put directly into the martingale decomposition as follows: from the computation above, and from Dynkin's formula, the following quantity is a martingale:

$$\begin{aligned} \mathcal{Z}_t^n(\varphi) - \mathcal{Z}_0^n(\varphi) - \int_0^t \left(n^2 \mathcal{L}_n^\otimes \mathcal{Z}_s^n(\varphi) - \frac{c_n}{n^{3/2}} \sum_{x=1}^n \varphi_x \right) ds \\ = \mathcal{Z}_t^n(\varphi) - \mathcal{Z}_0^n(\varphi) - \left(1 + \frac{E}{2n^\gamma}\right) \int_0^t \mathcal{Z}_s^n(\tilde{\Delta}_n \varphi) ds + \mathcal{R}_t^n(\varphi) + o_t^n(1) \end{aligned}$$

$$-\frac{E}{n^\gamma} \frac{n^2}{n^{3/2}} \int_0^t \sum_{x=1}^n \left(\nabla^- h_{sn^2}^n(x) \nabla^+ h_{sn^2}^n(x) - \frac{1}{4} \right) \varphi_x ds \quad (\text{A.11})$$

$$+ \frac{c_n t}{n^{3/2}} \sum_{x=1}^n \varphi_x, \quad (\text{A.12})$$

where the last term (A.12) cancels out with the term coming with $(-\frac{1}{4})$ in (A.11) (since $c_n = -En^{2-\gamma}/4$), and then (4.23) follows.

APPENDIX B. PROOFS OF PROPOSITION 3.14 AND PROPOSITION 3.15: BOUNDARY BEHAVIOR

As explained in Section 6, we may assume without loss of generality that $A = 1$ and $D = 2$, and recall that we denote $\bar{E} = E$. Therefore, in Proposition 3.15 we have $\frac{D(\bar{E})^2}{4A^3} = \frac{E^2}{2}$.

Proof of Proposition 3.14. Let \mathcal{V} and ρ_ε be as in the assumptions of Proposition 3.14. The map $\mathcal{V}(\rho_\varepsilon) = W_1(\rho_\varepsilon)(\mathcal{V})$ is in the first chaos, so Lemma 6.4 and Corollary 6.9 give

$$\begin{aligned} \mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t \mathcal{V}_s(\rho_\varepsilon) ds \right|^p \right] &\lesssim T^{p/2} \|\rho_\varepsilon\|_{\mathcal{H}_{\text{Dir}}^{p-1}}^p \\ &\lesssim T^{p/2} \left(\sup_{\substack{f \in \mathcal{H}_{\text{Dir}}^1 \\ \|f\|_{\mathcal{H}_{\text{Dir}}^1} = 1}} \int_0^1 \rho^\varepsilon(u) (f(u) - f(0)) du \right)^p \\ &\leq T^{p/2} \left(\sup_{\substack{f \in \mathcal{H}_{\text{Dir}}^1 \\ \|f\|_{\mathcal{H}_{\text{Dir}}^1} = 1}} \int_0^1 \rho^\varepsilon(u) \sqrt{u} \|f\|_{\mathcal{H}_{\text{Dir}}^1} du \right)^p \\ &= T^{p/2} \left(\int_0^1 \rho^\varepsilon(u) \sqrt{u} du \right)^p, \end{aligned}$$

where in the second step we used that $f(0) = 0$ and in the third step we applied the fundamental theorem of calculus and the Cauchy-Schwarz inequality. By assumption, the right hand side converges to zero as $\varepsilon \rightarrow 0$. \square

Remark B.1. In the case $p = 2$ and for $\rho_\varepsilon = \iota_\varepsilon(0)$ (resp. $\rho_\varepsilon = \iota_\varepsilon(1)$), with ι_ε as in Definition 3.2, the previous result can also be obtained from the microscopic dynamics by a two-steps procedure. We just sketch the idea and we leave the details to the reader. Note that $\mathcal{V}_s^n(\iota_\varepsilon(0)) = \sqrt{n} \vec{\eta}_{sn^2}^{\varepsilon n}(0)$. Adapting [12, Lemma 3] we know how to control the variance of $\int_0^t \sqrt{n}(\eta_{sn^2}(1) - \rho) ds$ and adapting [11, Lemma 7.1] we can control the variance of $\int_0^t \sqrt{n}(\eta_{sn^2}(1) - \vec{\eta}_{sn^2}^{\varepsilon n}(0)) ds$. This together with the convergence of the fluctuation field is enough to conclude.

Proof of Proposition 3.15. Let us use the same mapping argument as in Section 6.3. Let $\Psi_s^\delta(u) = e^{E\mathcal{V}_s(\Theta_u^\delta)}$ with $\Theta_u^\delta(v) = \int_0^1 \Theta_u(w) p_\delta^{\text{Dir}}(v, w) dw$ and where p_δ^{Dir} denotes the Dirichlet heat kernel as defined in (6.15). Then

$$\mathbb{E} \left[\left| \int_0^t \left(\frac{\Psi_s(\varepsilon) - \Psi_s(0)}{\varepsilon} + \frac{E^2}{2} \Psi_s(0) \right) ds \right|^2 \right]$$

$$\begin{aligned}
&= \lim_{\delta \rightarrow 0} \mathbb{E} \left[\left| \int_0^t \left(\frac{\Psi_s^\delta(\varepsilon) - \Psi_s^\delta(0)}{\varepsilon} + \frac{E^2}{2} \Psi_s(0) \right) ds \right|^2 \right] \\
&\leq 2 \lim_{\delta \rightarrow 0} \mathbb{E} \left[\left| \int_0^t \left(\frac{\Psi_s^\delta(\varepsilon) - \Psi_s^\delta(0)}{\varepsilon} - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_s^\delta(u) \Theta_u^{2\delta}(u) du \right) ds \right|^2 \right] \\
&\quad + 2 \lim_{\delta \rightarrow 0} \mathbb{E} \left[\left| \int_0^t \left(\frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_s^\delta(u) \Theta_u^{2\delta}(u) du + \frac{E^2}{2} \Psi_s(0) \right) ds \right|^2 \right]. \tag{B.1}
\end{aligned}$$

The Kipnis-Varadhan inequality (Corollary 6.6), yields for the first term

$$\begin{aligned}
&\mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t \left(\frac{\Psi_s^\delta(\varepsilon) - \Psi_s^\delta(0)}{\varepsilon} - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_s^\delta(u) \Theta_u^{2\delta}(u) du \right) ds \right|^2 \right] \\
&\lesssim T \left\| \frac{\Psi_0^\delta(\varepsilon) - \Psi_0^\delta(0)}{\varepsilon} - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_0^\delta(u) \Theta_u^{2\delta}(u) du \right\|_{-1,0}^2 \\
&\lesssim T \left\| \frac{1}{\varepsilon} \int_0^\varepsilon \Psi_0^\delta(u) E \mathcal{Y}_0(p_\delta^{\text{Dir}}(u, \cdot)) du - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_0^\delta(u) \Theta_u^{2\delta}(u) du \right\|_{-1,0}^2,
\end{aligned}$$

where in the last step we used the fundamental theorem of calculus and the identity $\nabla_u \mathcal{Y}_0(\Theta_u^\delta) = \mathcal{Y}_0(p_\delta^{\text{Dir}}(u, \cdot))$. Let now $F \in \mathcal{C}$ be a cylinder function. Then Gaussian integration by parts [35, Lemma 1.2.1] yields

$$\begin{aligned}
&\mathbb{E} \left[e^{E \mathcal{Y}_0(\Theta_u^\delta)} \mathcal{Y}_0(p_\delta^{\text{Dir}}(u, \cdot)) F(\mathcal{Y}_0) \right] \\
&= \int_0^1 \mathbb{E} \left[p_\delta^{\text{Dir}}(u, v) (e^{E \mathcal{Y}_0(\Theta_u^\delta)} D_v F(\mathcal{Y}_0) + F(\mathcal{Y}_0) D_v e^{E \mathcal{Y}_0(\Theta_u^\delta)}) \right] dv \\
&= \int_0^1 \mathbb{E} \left[p_\delta^{\text{Dir}}(u, v) (e^{E \mathcal{Y}_0(\Theta_u^\delta)} D_v F(\mathcal{Y}_0) + F(\mathcal{Y}_0) e^{E \mathcal{Y}_0(\Theta_u^\delta)} E \Theta_u^\delta(v)) \right] dv \\
&= \mathbb{E} \left[\int_0^1 p_\delta^{\text{Dir}}(u, v) e^{E \mathcal{Y}_0(\Theta_u^\delta)} D_v F(\mathcal{Y}_0) dv + F(\mathcal{Y}_0) e^{E \mathcal{Y}_0(\Theta_u^\delta)} E \Theta_u^{2\delta}(u) \right],
\end{aligned}$$

and therefore

$$\begin{aligned}
&\mathbb{E} \left[\left(\frac{1}{\varepsilon} \int_0^\varepsilon \Psi_0^\delta(u) E \mathcal{Y}_0(p_\delta^{\text{Dir}}(u, \cdot)) du - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_0^\delta(u) \Theta_u^{2\delta}(u) du \right) F(\mathcal{Y}_0) \right] \\
&= \mathbb{E} \left[\int_0^1 \frac{E}{\varepsilon} \int_0^\varepsilon p_\delta^{\text{Dir}}(u, v) e^{E \mathcal{Y}_0(\Theta_u^\delta)} du D_v F(\mathcal{Y}_0) dv \right] \\
&\stackrel{\text{duality}}{\leq} \mathbb{E} \left[\left\| \frac{E}{\varepsilon} \int_0^\varepsilon p_\delta^{\text{Dir}}(u, \cdot) e^{E \mathcal{Y}_0(\Theta_u^\delta)} du \right\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 \right]^{1/2} \|F\|_{1,0} \\
&\stackrel{(C-S)}{\leq} \mathbb{E} \left[\frac{E^2}{\varepsilon} \int_0^\varepsilon \|p_\delta^{\text{Dir}}(u, \cdot)\|_{\mathcal{H}_{\text{Dir}}^{-1}}^2 e^{2E \mathcal{Y}_0(\Theta_u^\delta)} du \right]^{1/2} \|F\|_{1,0} \lesssim \sup_{u \in (0, \varepsilon]} \|p_\delta^{\text{Dir}}(u, \cdot)\|_{\mathcal{H}_{\text{Dir}}^{-1}} \|F\|_{1,0},
\end{aligned}$$

which leads to the estimate

$$\lim_{\delta \rightarrow 0} \left\| \frac{\Psi^\delta(\varepsilon) - \Psi^\delta(0)}{\varepsilon} - \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi^\delta(u) \Theta_u^{2\delta}(u) du \right\|_{-1,0} \lesssim \limsup_{\delta \rightarrow 0} \sup_{u \in (0, \varepsilon]} \|p_\delta^{\text{Dir}}(u, \cdot)\|_{\mathcal{H}_{\text{Dir}}^{-1}},$$

and as in the proof of Proposition 3.14 we see that the right hand side converges to zero for $\varepsilon \rightarrow 0$. To treat the second term in (B.1) recall that we showed in the proof of

Lemma 6.24 that $\lim_{\delta \rightarrow 0} \Theta_u^{2\delta}(u) = u - \frac{1}{2}$ for all $u \in (0, 1)$. Therefore,

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \mathbb{E} \left[\sup_{t \in [0, T]} \left| \int_0^t \left(\frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_s^\delta(u) \Theta_u^{2\delta}(u) du + \frac{E^2}{2} \Psi_s(0) \right) ds \right|^2 \right] \\ \leq \lim_{\varepsilon \rightarrow 0} \mathbb{E} \left[\left(\int_0^T \left| \frac{E^2}{\varepsilon} \int_0^\varepsilon \Psi_s(u) (u - \tfrac{1}{2}) du + \frac{E^2}{2} \Psi_s(0) \right| ds \right)^2 \right] = 0 \end{aligned}$$

by Lebesgue's differentiation Theorem. The same arguments give the boundary behavior at $u = 1$. \square

APPENDIX C. HEAT KERNEL ESTIMATES

Here we collect basic estimates for the Dirichlet and Neumann heat kernels on $[0, 1]$, namely p^{Dir} and p^{Neu} , that we already defined respectively in (6.15) and (6.16).

Lemma C.1. *For all $\lambda > -1$ and $t \in (0, 1]$*

$$\begin{aligned} \sup_{u, v \in [0, 1]} |p_t^{\text{Dir}}(u, v)| &\lesssim t^{-1/2}, \quad \sup_{u \in [0, 1]} \|p_t^{\text{Dir}}(u, \cdot) \times |u - \cdot|^\lambda\|_{L^1} \lesssim t^{\lambda/2}, \\ \|\langle p_t^{\text{Dir}}(u, \cdot), 1 \rangle - 1\|_{L_u^2} &\lesssim t^{1/4}. \end{aligned}$$

The same bounds also hold for the Neumann heat kernel p^{Neu} .

Proof of Lemma C.1. The Dirichlet heat kernel is the transition density of $\{B_{2t} : t \geq 0\}$ where B is a Brownian motion that is killed when it reaches 0 or 1. In particular it is bounded from above by the transition density of the Brownian motion, i.e.

$$0 \leq p_t^{\text{Dir}}(u, v) \lesssim t^{-1/2} e^{-(u-v)^2/4t},$$

from where the first estimate follows. This also gives for $\lambda > -1$

$$\begin{aligned} \int_0^1 p_t^{\text{Dir}}(u, v) |u - v|^\lambda dv &\lesssim t^{\lambda/2} \int_0^1 t^{-1/2} e^{-\frac{1}{4} \left(\frac{u-v}{\sqrt{t}} \right)^2} \left| \frac{u-v}{\sqrt{t}} \right|^\lambda dv \\ &= t^{\lambda/2} \int_{\mathbb{R}} e^{-\frac{1}{4} u^2} |u|^\lambda du \lesssim t^{\lambda/2}. \end{aligned}$$

To estimate the L^2 -norm, note that if B is a standard Brownian motion we have

$$\begin{aligned} 1 &\geq \langle p_t^{\text{Dir}}(u, \cdot), 1 \rangle = \mathbb{P}(u + B_{2s} \in [0, 1], \forall s \in [0, t]) \\ &\geq \mathbb{P}(\{ \sup_{s \in [0, 2t]} B_s < 1 - u \} \cap \{ \inf_{s \in [0, 2t]} B_s > -u \}) \\ &\geq \mathbb{P}(\sup_{s \in [0, 2t]} B_s < 1 - u) + \mathbb{P}(\inf_{s \in [0, 2t]} B_s > -u) - 1 \geq 2\mathbb{P}(\sup_{s \in [0, 2t]} B_s < u \wedge (1 - u)) - 1 \\ &= 2\mathbb{P}(|B_{2t}| < u \wedge (1 - u)) - 1 \geq 2(1 - 2\mathbb{P}(B_{2t} \geq u \wedge (1 - u))) - 1 \\ &= 1 - 4\mathbb{P}(B_{2t} \geq u \wedge (1 - u)) \geq 1 - 4 \frac{e^{-\frac{(u \wedge (1-u))^2}{4t}}}{\sqrt{2\pi} \frac{u \wedge (1-u)}{\sqrt{2t}}}, \end{aligned}$$

where we used the reflection principle for the Brownian motion and standard tail estimates for the normal distribution. From here we get $|\langle p_t^{\text{Dir}}(u, \cdot), 1 \rangle - 1| \lesssim \frac{t^{1/2}}{u(1-u)} \wedge 1$, which leads to

$$\|\langle p_t^{\text{Dir}}(u, \cdot), 1 \rangle - 1\|_{L_u^2}^2 \lesssim \int_0^{t^{1/2}} du + \int_{t^{1/2}}^\infty \frac{t}{u^2} du \lesssim t^{1/2}.$$

For the Neumann heat kernel we have $\langle p_t^{\text{Neu}}(u, \cdot) - 1 \rangle \equiv 1$, so the last bound is trivial. The remaining bounds for the Neumann kernel follow once we know that $p_t^{\text{Neu}}(u, v) \lesssim t^{-1/2} e^{-(u-v)^2/4t}$ uniformly in $t \in (0, 1]$, which is basically (3.7) in [41]. \square

Lemma C.2. *The difference between Neumann and Dirichlet heat kernel is bounded in L^1 :*

$$\|p_\varepsilon^{\text{Neu}}(u, \cdot) - p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^1} \lesssim \frac{\varepsilon^{1/2}}{u(1-u)} \wedge 1,$$

uniformly in $\varepsilon \in (0, 1]$ and $u \in [0, 1]$.

Proof of Lemma C.2. Note that $\|p_\varepsilon^{\text{Neu}}(u, \cdot) - p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{L^1} \leq 2\|p_\varepsilon^{\text{Neu}}(u, \cdot) - p_\varepsilon^{\text{Dir}}(u, \cdot)\|_{\text{TV}}$, where

$$\|\nu\|_{\text{TV}} = \sup_{A \in \mathcal{B}([0,1])} |\nu(A)|$$

denotes the total variation norm of the signed measure ν on the Borel sets $\mathcal{B}([0, 1])$ of $[0, 1]$. We know that p^{Neu} (resp. p^{Dir}) is the transition density of a Brownian motion that is reflected (resp. killed) in 0 and 1 – both with speed 2. We write $\mathbb{P}_u^{\text{Neu}}$ respectively $\mathbb{P}_u^{\text{Dir}}$ for the law of the reflected respectively killed Brownian motion with speed 2, both started in u , while \mathbb{P}_u is the law of the (usual) Brownian motion with speed 2, started in u . Then we have for all Borel sets $A \in \mathcal{B}(\mathbb{R})$

$$\mathbb{P}_u^{\text{Dir}}(\{B_\varepsilon \in A\} \cap \{B_s \in [0, 1], \forall s \in [0, \varepsilon]\}) = \mathbb{P}_u^{\text{Neu}}(\{B_\varepsilon \in A\} \cap \{B_s \in [0, 1], \forall s \in [0, \varepsilon]\}),$$

and therefore

$$|\mathbb{P}_u^{\text{Dir}}(B_\varepsilon \in A) - \mathbb{P}_u^{\text{Neu}}(B_\varepsilon \in A)| \leq 2\mathbb{P}_u(\{\sup_{s \leq \varepsilon} B_s \geq 1\} \cup \{\inf_{s \leq \varepsilon} B_s \leq 0\}) \lesssim \frac{\varepsilon^{1/2}}{u(1-u)},$$

where the last step follows as in the proof of Lemma C.1. We now take the supremum in $A \in \mathcal{B}([0, 1])$ and get a bound for the total variation norm and thus for the L^1 -norm of the difference of the densities. \square

APPENDIX D. PROOF OF PROPOSITION 3.10

Let Φ satisfy the conditions of Proposition 3.10. As in [1, Chapter 5] we see that there exists a standard $\mathcal{S}'_{\text{Neu}}$ -valued Brownian motion with covariance (3.2) (possibly on an extended probability space) such that for all $\varphi \in \mathcal{S}_{\text{Neu}}$

$$\begin{aligned} \Phi_t(\varphi) &= \Phi_0(\varphi) + A \int_0^t \Phi_s(\Delta \varphi) ds + A \int_0^t (-\alpha \Phi_s(0) \varphi(0) + \beta \Phi_s(1) \varphi(1)) ds \\ &\quad + \sqrt{D} \int_0^1 \int_0^t \Phi_s(u) \varphi(u) d\mathcal{W}_s(u) du, \end{aligned} \tag{D.1}$$

and also that it suffices to show the *strong uniqueness* of solutions to the equation driven by this given \mathcal{W} . The proof for the strong uniqueness is essentially the same as in [41, Theorem 3.2], the only difference is that we have to deal with the additional terms

coming from the Robin boundary condition. If Φ^i , $i = 1, 2$ are solutions to (D.1), then we obtain easily (see also [41, Exercise 3.1]) that for all $t \in [0, T]$, $u \in [0, 1]$, and $\varepsilon > 0$

$$\begin{aligned}\Phi_t^i(p_\varepsilon^{\text{Neu}}(u)) &= A \int_0^1 \Phi_0(v) p_{t+\varepsilon}^{\text{Neu}}(u, v) dv \\ &\quad + A \int_0^t (-\alpha \Phi_s^i(0) p_{t-s+\varepsilon}^{\text{Neu}}(u, 0) + \beta \Phi_s^i(1) p_{t-s+\varepsilon}^{\text{Neu}}(u, 1)) ds \\ &\quad + \sqrt{D} \int_0^1 \int_0^t p_{t-s+\varepsilon}^{\text{Neu}}(u, v) \Phi_s^i(v) d\mathcal{W}_s(v) dv,\end{aligned}$$

and with the L^2 -continuity of Φ_t^i this extends to $\varepsilon = 0$ (with $p_0^{\text{Neu}}(u, \cdot) = \delta_u(\cdot)$). Then the difference $\mathcal{U} = \Phi^1 - \Phi^2$ satisfies

$$\begin{aligned}\mathbb{E}[\mathcal{U}_t(u)^2] &\lesssim \mathbb{E}\left[\left|\int_0^t (-\alpha \mathcal{U}_s(0) p_{t-s}^{\text{Neu}}(u, 0) + \beta \mathcal{U}_s(1) p_{t-s}^{\text{Neu}}(u, 1)) ds\right|^2\right] \\ &\quad + \int_0^1 \int_0^t p_{t-s}^{\text{Neu}}(u, v)^2 \mathbb{E}[\mathcal{U}_s(v)^2] ds dv.\end{aligned}$$

Set now $\mathcal{V}_t = \sup_{u \in [0, 1]} \mathbb{E}[\mathcal{U}_t(u)^2]$. Lemma C.1 gives

$$\sup_{u \in [0, 1]} \|p_t^{\text{Neu}}(u, \cdot)\|_{L^\infty} \lesssim t^{-1/2}, \quad \sup_{u \in [0, 1]} \|p_t^{\text{Neu}}(u, \cdot)\|_{L^1} \lesssim 1,$$

and we apply the Cauchy-Schwarz inequality for the measure $A \mapsto \int_0^t \mathbf{1}_A(s) (t-s)^{-1/2} ds$ to obtain, uniformly in $t \in [0, T]$,

$$\mathcal{V}_t \lesssim \int_0^t \mathcal{V}_s (t-s)^{-1/2} ds.$$

Now we can simply iterate this inequality to see that $\mathcal{V}_t = 0$ for all $t \in [0, T]$, for details see [41, Theorem 3.2]. This concludes the proof of Proposition 3.10. \square

Having proved uniqueness, let us prove an additional moment bound for the solution to (D.1), which has been used in Section 6.3.

Lemma D.1. *Assume that Φ_0 satisfies $\sup_{u \in [0, 1]} \mathbb{E}[|\Phi_0(u)|^p] = M < \infty$ for some $p \in (6, \infty)$ and let $\{\Phi_t; t \in [0, T]\}$ solve (D.1) with initial condition Φ_0 . Then there exists a constant $C > 0$ that only depends on T and α, β, A, D but not on M , such that*

$$\sup_{t \in [0, T], u \in [0, 1]} \mathbb{E}[|\Phi_t(u)|^p] \leq C \times M.$$

Remark D.2. *Of course, the same bound holds for $p \in [1, \infty)$. But for $p > 6$ the proof slightly simplifies.*

Proof of Lemma D.1. The following argument is also essentially contained in [41, Theorem 3.2]. By the Burkholder-Davis-Gundy inequality and Jensen's inequality we have uniformly in $t \in [0, T]$ and locally uniformly in α, β

$$\begin{aligned}\mathbb{E}[|\Phi_t(u)|^p] &\lesssim \int_0^1 \mathbb{E}[|\Phi_0(v)|^p] p_t^{\text{Neu}}(u, v) dv \\ &\quad + \int_0^t (p_{t-s}^{\text{Neu}}(u, 0) \mathbb{E}[|\Phi_s(0)|^p] + p_{t-s}^{\text{Neu}}(u, 1) \mathbb{E}[|\Phi_s(1)|^p]) ds\end{aligned}$$

$$+ \mathbb{E} \left[\left| \int_0^1 \int_0^t p_{t-s}^{\text{Neu}}(u, v)^2 \Phi_s(v)^2 ds dv \right|^{p/2} \right],$$

and we bound the integral inside the expectation with Hölder's inequality by

$$\begin{aligned} \left| \int_0^1 \int_0^t p_{t-s}^{\text{Neu}}(u, v)^2 \Phi_s(v)^2 ds dv \right|^{p/2} \\ \leq \left| \int_0^1 \int_0^t p_{t-s}^{\text{Neu}}(u, v)^{2q} ds dv \right|^{p/(2q)} \int_0^1 \int_0^t |\Phi_s(v)|^p ds dv \end{aligned}$$

for $q = p/(p-2)$, the conjugate exponent of $p/2$. To control the integral over the Neumann heat kernel we first integrate in the space variable and bound

$$\|p_{t-s}^{\text{Neu}}(u, \cdot)^{2q}\|_{L^1} \leq \|p_{t-s}^{\text{Neu}}(u, \cdot)^{2q-1}\|_{L^\infty} \|p_{t-s}^{\text{Neu}}(u, \cdot)\|_{L^1} \leq C(t-s)^{1/2-q},$$

where the last step follows from Lemma C.1. Since $p > 6$ we have $q = p/(p-2) < \frac{3}{2}$, and therefore $(t-s)^{1/2-q}$ is integrable on $[0, t]$ and the integral is uniformly bounded in $t \in [0, T]$. Thus, we have shown that $\mathcal{V}_t = \sup_{u \in [0, 1]} \mathbb{E}[|\Phi_t(u)|^p]$ satisfies on $[0, T]$

$$\begin{aligned} \mathcal{V}_t &\lesssim M + \int_0^t (1 + (t-s)^{-1/2}) \mathcal{V}_s ds \\ &\lesssim M + \int_0^t (1 + (t-s)^{-1/2}) \left(M + \int_0^s (1 + (s-r)^{-1/2}) \mathcal{V}_r dr \right) ds \\ &\lesssim M + \int_0^t \int_r^t (1 + (t-s)^{-1/2}) (1 + (s-r)^{-1/2}) \mathcal{V}_r ds dr \\ &\lesssim M + \int_0^t \mathcal{V}_r dr, \end{aligned}$$

where the last step is a simple computation: after expanding the product the most complicated integrand is $(t-r)^{-1/2}(s-r)^{-1/2}$, for which the integral is computed in [41, p.315]. Now the claim follows from Gronwall's Lemma. \square

APPENDIX E. SOME CONSIDERATIONS ABOUT THE MICROSCOPIC COLE-HOPF PROCESS

Since the major breakthrough made by [1], the KPZ behavior for interacting particle systems has often been investigated through the so-called Cole-Hopf process, and it has been – until very recently – the only way to prove the convergence of the microscopic height function. We note however, that this approach has only been possible for very particular microscopic dynamics which allow the Cole-Hopf transformation, such as the one we consider in this paper. Even if here we do not need to use this transformation in order to prove the KPZ-type macroscopic fluctuations, we want to highlight in this section the fact that the specific Robin boundary conditions that we obtain for the stochastic heat equation also emerge from the microscopic Cole-Hopf transformation.

In what follows, we assume $\gamma = \frac{1}{2}$ because we just want to recover the Robin boundary conditions obtained in Proposition 3.13. For that purpose, we define the *microscopic Cole-Hopf transformation* as follows: we set

$$\xi_t^n(x) = \exp\left(\frac{\theta_n}{\sqrt{n}} h_t^n(x) + \lambda_n t\right), \quad \text{for any } x \in \{1, \dots, n\}, \quad (\text{E.1})$$

where θ_n is such that:

$$\exp\left(\frac{\theta_n}{\sqrt{n}}\right) = 1 + \frac{E}{\sqrt{n}}, \quad (\text{E.2})$$

and λ_n will be chosen ahead. Note that $\theta_n \geq 0$. We will see below in Section E.1 the reason for this choice of θ_n .

Definition E.1. For any $t > 0$ let \mathcal{J}_t^n be the current fluctuation field which acts on functions $\varphi \in \mathcal{S}_{\text{Neu}}$ as:

$$\mathcal{J}_t^n(\varphi) = \frac{1}{n} \sum_{x=1}^n \varphi\left(\frac{x}{n}\right) \xi_{tn^2}^n(x).$$

Remark E.2. Before proceeding we note that by the conservation law, for any $x \in \Lambda_n$, we have that

$$\eta_t(x) - \eta_0(x) = J_{x-1,x}(t) - J_{x,x+1}(t) + \mathbf{1}_{\{1,n-1\}}(x) C_t^n(x), \quad (\text{E.3})$$

where $J_{x,x+1}(t)$ is the counting process for the net number of particles at the bond $\{x, x+1\}$ during the time interval $[0, t]$, with, by convention, $J_{0,1}(t) = 0 = J_{n-1,n}(t)$. For $x = 1$ and $x = n$, the process $C_t^n(x)$ counts the number of particles created at the site x minus the number of particles destroyed at the site x during the time interval $[0, t]$. Also note that $C_t^n(1) = -h_t^n(1)$, where $h_t^n(1)$ was introduced in (2.8). From (E.3) we have an equivalent way of defining the microscopic Cole-Hopf transform as

$$\xi_t^n(x) = \exp\left(\frac{-\theta_n}{\sqrt{n}} \left(J_{x-1,x}^n(t) - \sum_{y=1}^{x-1} \eta_0(y)\right)\right) \exp\left(-\frac{\theta_n(x-1)}{2\sqrt{n}} + \lambda_n t\right). \quad (\text{E.4})$$

The first exponential on the right hand side of the last equality corresponds to the Cole-Hopf transformation used in [21], except that our scaling is different, since the strength of asymmetry is taken here with $\gamma = \frac{1}{2}$, and not $\gamma = 1$ as in [21] (this explains the factor $\frac{1}{\sqrt{n}}$ in (E.4) instead of $\frac{1}{n}$).

Note also that, up to the level of the current fluctuation field \mathcal{J}^n , both definitions of the Cole-Hopf process basically coincide: in [21] the current field is defined on the discrete gradient of the test function φ which is assumed to vanish at the boundary, see (3.1) and the equation above (3.11) in [21]. Here, instead we take test functions in \mathcal{S}_{Neu} (Definition E.1). Moreover, the extra factor at the right hand side of (E.4) corresponds to the average of the microscopic Cole-Hopf variables and also appears in the definition of the current field in [21] (formula above (3.11)).

Remark E.3. Note that, when $t = 0$, one can compute the average of the current fluctuation field with respect to \mathbb{E}_ρ as follows:

$$\begin{aligned} \mathbb{E}_\rho[\mathcal{J}_0^n(\varphi)] &= \frac{1}{n} \sum_{x=1}^n \varphi\left(\frac{x}{n}\right) \mathbb{E}[\xi_0^n(x)] \\ &= \frac{1}{n} \sum_{x=1}^n \varphi\left(\frac{x}{n}\right) \left(\frac{1}{2} \left(e^{\theta_n/(2\sqrt{n})} + e^{-\theta_n/(2\sqrt{n})}\right)\right)^{x-1} \\ &\xrightarrow{n \rightarrow \infty} \int_0^1 \varphi(u) e^{uE^2/2} du. \end{aligned}$$

Without entering too much into details, let us give in this last section some properties of the microscopic Cole-Hopf transformation that one could prove using some well-known past works.

E.1. Martingale decomposition for the current fluctuation field \mathcal{J}_t^n . Let us first write the martingale decomposition that the current field satisfies, in the case $\gamma = \frac{1}{2}$. For that purpose we take φ as a test function, which will be chosen ahead, and denote as before $\varphi_x = \varphi(\frac{x}{n})$. Recall also from (E.1)

$$\xi_t^n(x) = \exp\left(\frac{\theta_n}{\sqrt{n}}\left(h_t^n(1) + \sum_{y=1}^{x-1} \bar{\eta}_t(y)\right) + \lambda_n t\right), \quad \text{for any } x \in \{1, \dots, n\}.$$

First, we are going to explain the choice for θ_n and λ_n . By using the usual convention $\eta_t(0) = \eta_t(n) = \frac{1}{2}$, since ξ_t^n is a function of η and $h_t^n(1)$, recalling (4.22), we have, for $x \in \{1, \dots, n\}$, that

$$\begin{aligned} \xi_{tn^2}^n(x) - \xi_0^n(x) - \int_0^t n^2 \xi_{sn^2}^n(x) \left[\lambda_n + (e^{\theta_n/\sqrt{n}} - 1)(1 - \eta_{sn^2}(x-1))\eta_{sn^2}(x) \right] ds \\ - \int_0^t n^2 \xi_{sn^2}^n(x) \left[(e^{-\theta_n/\sqrt{n}} - 1)\left(1 + \frac{E}{\sqrt{n}}\right)\eta_{sn^2}(x-1)(1 - \eta_{sn^2}(x)) \right] ds, \end{aligned}$$

is a martingale. From the choice (E.2) we see that the integral term above simplifies to

$$\int_0^t n^2 \xi_{tn^2}^n(x) \lambda_n + En^{3/2} \xi_{sn^2}^n(x) [\eta_{sn^2}(x) - \eta_{sn^2}(x-1)] ds. \quad (\text{E.5})$$

In order to close the equation above in terms of $\xi_{sn^2}^n(x)$ we note that from the trivial identities:

$$\begin{aligned} \xi_{sn^2}^n(x) \eta_{sn^2}(x-1) a_n + \xi_{sn^2}^n(x) b_n &= \xi_{sn^2}^n(x-1) - \xi_{sn^2}^n(x), \quad \text{for } x \in \{2, \dots, n\} \\ \xi_{sn^2}^n(x) \eta_{sn^2}(x) a_n + \xi_{sn^2}^n(x) b_n e^{-\theta_n/(2\sqrt{n})} &= \xi_{sn^2}^n(x) - \xi_{sn^2}^n(x+1), \quad \text{for } x \in \{1, \dots, n-1\} \end{aligned} \quad (\text{E.6})$$

with $a_n = (e^{-\theta_n/(2\sqrt{n})} - e^{\theta_n/(2\sqrt{n})})$ and $b_n = e^{\theta_n/(2\sqrt{n})} - 1$, the expression (E.5), for $x \in \{2, \dots, n-1\}$, is reduced to

$$\int_0^t n^2 \xi_{sn^2}^n(x) \lambda_n - \frac{En^{3/2}}{a_n} \Delta \xi_{sn^2}^n(x) + \frac{En^{3/2}}{a_n} \xi_{sn^2}^n(x) b_n^2 e^{-\theta_n/(2\sqrt{n})} ds.$$

Now note that for the choice $\sqrt{n} \lambda_n = -\frac{E}{a_n} b_n^2 e^{-\theta_n/(2\sqrt{n})}$ the last identity becomes

$$\int_0^t -\frac{En^{3/2}}{a_n} \Delta \xi_{sn^2}^n(x) ds.$$

Also note that

$$\theta_n \xrightarrow{n \rightarrow \infty} E; \quad a_n \xrightarrow{n \rightarrow \infty} \frac{-E}{\sqrt{n}}; \quad \lambda_n \xrightarrow{n \rightarrow \infty} \frac{E^2}{4n}. \quad (\text{E.7})$$

To treat the boundary we do the following. For $x = 1$ (resp. $x = n$) we plug the second (resp. first) identity of (E.6) in (E.5) to arrive at (E.5) for $x = 1$ (resp. $x = n$) written as

$$\begin{aligned} \int_0^t n^2 \xi_{sn^2}^n(1) \lambda_n - \frac{En^{3/2}}{a_n} \nabla^+ \xi_{sn^2}^n(1) - En^{3/2} \xi_{sn^2}^n(1) \left(\frac{b_n e^{-\theta_n/(2\sqrt{n})}}{a_n} + \frac{1}{2} \right) ds, \\ \left(\text{resp. } \int_0^t n^2 \xi_{sn^2}^n(n) \lambda_n + \frac{En^{3/2}}{a_n} \nabla^- \xi_{sn^2}^n(n) + En^{3/2} \xi_{sn^2}^n(n) \left(\frac{b_n}{a_n} + \frac{1}{2} \right) ds. \right) \end{aligned} \quad (\text{E.8})$$

Collecting the previous observations we conclude that for any $x \in \{1, \dots, n\}$,

$$\xi_{tn^2}^n(x) - \xi_0^n(x) - \int_0^t (T_n \xi_{sn^2}^n)(x) ds$$

is a martingale, where T_n is given by

$$\begin{cases} T_n \xi^n(1) &= n\alpha_n \xi^n(1) + D_n n^2 \nabla^+ \xi^n(1), \\ T_n \xi^n(x) &= D_n n^2 \Delta \xi^n(x), & \text{if } x \in \{2, \dots, n-1\} \\ T_n \xi^n(n) &= n\alpha_n \xi^n(n) - D_n n^2 \nabla^- \xi^n(n), \end{cases} \quad (\text{E.9})$$

where

$$D_n = -\frac{E}{a_n \sqrt{n}}, \quad \alpha_n = n\lambda_n - \frac{E\sqrt{n}}{2} \left(\frac{2 - e^{\theta_n/(2\sqrt{n})} - e^{-\theta_n/(2\sqrt{n})}}{a_n} \right)$$

and from (E.7) we get that $\alpha_n \rightarrow \frac{E^2}{8}$ and $D_n \rightarrow 1$, as $n \rightarrow \infty$. Now, from Dynkin's formula, we compute the martingale decomposition for the current fluctuation field \mathcal{J}_t^n . From the computations above and by doing a summation by parts, we get that

$$\mathcal{J}_t^n(\varphi) - \mathcal{J}_0^n(\varphi) - \int_0^t \frac{1}{n} \sum_{x=1}^n (T_n \xi_{sn^2}^n)(x) \varphi_x ds \quad (\text{E.10})$$

is a martingale, where the integral term above can be rewritten as

$$\begin{aligned} \int_0^t \frac{1}{n} \sum_{x=1}^{n-1} D_n \Delta_n \varphi_x \xi_{sn^2}^n(x) ds + \int_0^t \frac{1}{n} \left(D_n n^2 (\varphi_1 - \varphi_0) + n\alpha_n \varphi_1 \right) \xi_{sn^2}^n(1) ds \\ + \int_0^t \frac{1}{n} \left(D_n n^2 (\varphi_{n-1} - \varphi_n) + n\alpha_n \varphi_n \right) \xi_{sn^2}^n(n) ds. \end{aligned}$$

Now to close the field let us sum and subtract, inside the previous sum, the term $\Delta_n \varphi_{n-1} \xi_{sn^2}^n(n)$ so that the last expression becomes:

$$\begin{aligned} \int_0^t \frac{1}{n} \sum_{x=1}^n (T_n \xi_{sn^2}^n)(x) \varphi_x ds &= \int_0^t \frac{1}{n} \sum_{x=1}^{n-1} D_n \tilde{\Delta}_n \varphi_x \xi_{sn^2}^n(x) ds \\ &+ \int_0^t \frac{1}{n} \left(D_n n^2 (\varphi_1 - \varphi_0) + n\alpha_n \varphi_1 \right) \xi_{sn^2}^n(1) ds \\ &+ \int_0^t \frac{1}{n} \left(D_n n^2 (\varphi_{n-1} - \varphi_n - \tilde{\Delta}_n \varphi_n) + n\alpha_n \varphi_n \right) \xi_{sn^2}^n(n) ds. \end{aligned}$$

where $\tilde{\Delta}_n$ is a slightly different approximation of the discrete Laplacian defined as

$$\tilde{\Delta}_n \varphi\left(\frac{x}{n}\right) = \begin{cases} \Delta_n \varphi\left(\frac{x}{n}\right) & \text{if } x \in \{1, \dots, n-1\}, \\ \Delta_n \varphi\left(\frac{n-1}{n}\right) & \text{if } x = n. \end{cases} \quad (\text{E.11})$$

Putting everything together, we rewrite (E.10) as

$$\mathcal{J}_t^n(\varphi) - \mathcal{J}_0^n(\varphi) - D_n \int_0^t \mathcal{J}_s^n(\tilde{\Delta}_n \varphi) ds \quad (\text{E.12})$$

$$- \frac{1}{n} \left(D_n n^2 \left(\varphi\left(\frac{1}{n}\right) - \varphi(0) \right) + n\alpha_n \varphi\left(\frac{1}{n}\right) \right) \int_0^t \xi_{sn^2}^n(1) ds \quad (\text{E.13})$$

$$-\frac{1}{n} \left(D_n n^2 \left(3\varphi\left(\frac{n-1}{n}\right) - 2\varphi(1) - \varphi\left(\frac{n-2}{n}\right) \right) + n\beta_n \varphi(1) \right) \int_0^t \xi_{sn^2}^n(n) ds. \quad (\text{E.14})$$

We note that last expression is a martingale whose quadratic variation is given by

$$\int_0^t \frac{E^2}{n} \sum_{x=1}^n \left(\varphi\left(\frac{x}{n}\right) \right)^2 \left(\xi_{sn^2}^n(x) \right)^2 \left[\eta_{sn^2}(x) (1 - \eta_{sn^2}(x-1)) + \eta_{sn^2}(x-1) (1 - \eta_{sn^2}(x)) e^{-\theta_n/\sqrt{n}} \right] ds. \quad (\text{E.15})$$

Note that in (E.14)

$$\left(3\varphi\left(\frac{n-1}{n}\right) - 2\varphi(1) - \varphi\left(\frac{n-2}{n}\right) \right) = -\frac{1}{n} \nabla \varphi(1) + O(n^{-2}).$$

E.2. Asymptotic limit. From (E.12)–(E.14), one sees that there are two natural ways to close the martingale problem. First, note that if $\varphi \in \mathcal{C}^\infty([0, 1])$, then $\tilde{\Delta}_n \varphi$ is an approximation of $\Delta \varphi$. Moreover,

- if φ satisfies the Robin boundary condition
 - (i) $\nabla \varphi(0) = -\frac{E^2}{8} \varphi(0)$, then (E.13) is of order $1/n$ in $L^2(\mathbb{P}_\rho)$ and therefore vanishes as $n \rightarrow \infty$;
 - (ii) $\nabla \varphi(1) = \frac{E^2}{8} \varphi(1)$, then (E.14) is also of order $1/n$ and vanishes.
- If, however, $\varphi \in \mathcal{S}_{\text{Neu}}$, then $n(\varphi_1 - \varphi_0) \rightarrow 0$ as $n \rightarrow \infty$, and therefore only one term remains in (E.13), which reads

$$\alpha_n \varphi\left(\frac{1}{n}\right) \int_0^t \xi_s^n(1) ds \underset{n \rightarrow \infty}{\simeq} \frac{E^2}{8} \varphi\left(\frac{1}{n}\right) \int_0^t \xi_s^n(1) ds.$$

In the macroscopic limit, the last term will correspond to

$$\frac{E^2}{8} \int_0^t \Phi_s(0) \varphi(0) ds$$

in the definition (3.13) of the solution to the SHE with Robin boundary condition, as soon as one proves convergence of $\xi_{tn^2}^n$ in $\mathcal{D}([0, T], \mathcal{C}([0, 1]))$.

E.3. Exponential moments and quadratic variation. Finally, one might check that the quadratic variation (E.15) converges to

$$\frac{E^2 t}{2} \int_0^1 |\Phi_s(u)|^2 \varphi^2(u) du,$$

where Φ is the limit of the current field \mathcal{J}^n in $\mathcal{D}([0, T], \mathcal{C}([0, 1]))$. Heuristically, this is indeed the case if one is able to replace $\xi_{sn^2}^n(x)^2$ in (E.15) with $(\mathbb{E}[\xi_{sn^2}^n(x)])^2$. This could be proved by using some ideas taken from [21, Lemma 4.3] which permit to control all the exponential moments

$$\sup_{x \in \{1, \dots, n\}} \mathbb{E}_\rho \left[\left(\xi_{sn^2}^n(x) - \mathbb{E}[\xi_{sn^2}^n(x)] \right)^k \right],$$

with $k \in \mathbb{N}$.

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REFERENCES

- [1] L. Bertini and G. Giacomin. Stochastic Burgers and KPZ equations from particle systems. *Comm. Math. Phys.*, 183(3):571–607, 1997.
- [2] O. Blondel, P. Gonçalves, and M. Simon. Convergence to the stochastic Burgers equation from a degenerate microscopic dynamics. *Electron. J. Probab.*, 21:25 pp., 2016.
- [3] C. Chang, C. Landim, and S. Olla. Equilibrium Fluctuations of asymmetric simple exclusion processes in dimension $d \geq 3$. *Probability Theory and Related Fields*, 119(3):381–409, 2001.
- [4] I. Corwin. The Kardar-Parisi-Zhang equation and universality class. *Random Matrices: Theory and Applications*, 01(01):1130001, 2012.
- [5] I. Corwin and H. Shen. Open ASEP in the weakly asymmetric regime. *Communications in Pure Appl. Math.*, 2016.
- [6] I. Corwin, H. Shen, and L.-C. Tsai. Asep(q,j) converges to the KPZ equation. *Ann. Inst. H. Poincaré Probab. Statist.*, 54(2):995–1012, 02 2018.
- [7] I. Corwin and L.-C. Tsai. KPZ equation limit of higher-spin exclusion processes. *Ann. Probab.*, 45(3):1771–1798, 05 2017.
- [8] A. Dembo and L.-C. Tsai. Weakly asymmetric non-simple exclusion process and the Kardar–Parisi–Zhang equation. *Communications in Mathematical Physics*, 341(1):219–261, Jan 2016.
- [9] J. Diehl, M. Gubinelli, and N. Perkowski. The Kardar-Parisi-Zhang equation as scaling limit of weakly asymmetric interacting Brownian motions. *Comm. Math. Phys.*, 354(2):549–589, 2017.
- [10] T. Franco, P. Gonçalves, and M. Simon. Crossover to the Stochastic Burgers Equation for the WASEP with a slow bond. *Comm. Math. Phys.*, 346(3):801–838, 2016.
- [11] T. Franco, P. Gonçalves, and A. Neumann. Phase transition in equilibrium fluctuations of symmetric slowed exclusion. *Stochastic Processes and their Applications*, 123(12):4156 – 4185, 2013.
- [12] Tertuliano Franco, Patrícia Gonçalves, and Adriana Neumann. Equilibrium fluctuations for the slow boundary exclusion process. In Patrícia Gonçalves and Ana Jacinta Soares, editors, *From Particle Systems to Partial Differential Equations*, pages 177–197, Cham, 2017. Springer International Publishing.
- [13] M. Freidlin. *Functional integration and partial differential equations*, volume 109 of *Annals of Mathematics Studies*. Princeton University Press, Princeton, NJ, 1985.
- [14] J. Gärtner. Convergence towards Burger's equation and propagation of chaos for weakly asymmetric exclusion processes. *Stochastic Processes and their Applications*, 27(Supplement C):233 – 260, 1987.
- [15] Máté Gerencsér and Martin Hairer. Singular spdes in domains with boundaries. *Probability Theory and Related Fields*, Mar 2018.
- [16] P. Gonçalves. Central limit theorem for a tagged particle in asymmetric simple exclusion. *Stochastic Processes and their Applications*, 118(3):474–502, 2008.
- [17] P. Gonçalves and M. Jara. Nonlinear fluctuations of weakly asymmetric interacting particle systems. *Arch. Rational Mech. Anal.*, 212(2):597–644, 2014.

- [18] P. Gonçalves and M. Jara. Stochastic Burgers Equations from Long Range Exclusion Interactions. *to appear in Stochastic Processes and Their Applications*, 2017.
- [19] P. Gonçalves, M. Jara, and S. Sethuraman. A stochastic Burgers equation from a class of microscopic interactions. *Ann. Probab.*, 43(1):286–338, 2015.
- [20] P. Gonçalves, M. Jara, and M. Simon. Second order Boltzmann-Gibbs principle for polynomial functions and applications. *J. Stat. Phys.*, 166(1):90–113, 2017.
- [21] P. Gonçalves, C. Landim, and A. Milanés. Nonequilibrium fluctuations of one-dimensional boundary driven weakly asymmetric exclusion processes. *ArXiv preprint:1504.05771*, 2016.
- [22] M. Gubinelli, P. Imkeller, and N. Perkowski. Paracontrolled distributions and singular PDEs. *Forum of Mathematics, Pi*, 3(e6), 2015.
- [23] M. Gubinelli and M. Jara. Regularization by noise and stochastic Burgers equations. *Stochastic Partial Differential Equations: Analysis and Computations*, 1(2):325–350, 2013.
- [24] M. Gubinelli and N. Perkowski. The Hairer–Quastel universality result at stationarity. *RIMS Kôkyûroku Bessatsu*, B59, 2016.
- [25] Massimiliano Gubinelli and Nicolas Perkowski. Energy solutions of KPZ are unique. *J. Amer. Math. Soc.*, 31(2):427–471, 2018.
- [26] Massimiliano Gubinelli and Nicolas Perkowski. Probabilistic approach to the stochastic burgers equation. In Andreas Eberle, Martin Grothaus, Walter Hoh, Moritz Kassmann, Wilhelm Stannat, and Gerald Trutnau, editors, *Stochastic Partial Differential Equations and Related Fields*, pages 515–527, Cham, 2018. Springer International Publishing.
- [27] M. Hairer. Solving the KPZ equation. *Annals of Mathematics*, 178(2):559–664, 2013.
- [28] M. Hairer. A theory of regularity structures. *Invent. Math.*, 198(2):269–504, 2014.
- [29] S. Janson. *Gaussian Hilbert spaces*, volume 129 of *Cambridge Tracts in Mathematics*. Cambridge University Press, Cambridge, 1997.
- [30] M. Kardar, G. Parisi, and Y. Zhang. Dynamic scaling of growing interfaces. *Phys. Rev. Lett.*, 56:889–892, 1986.
- [31] C. Kipnis and C. Landim. *Scaling limits of interacting particle systems*, volume 320 of *Grundlehren der mathematischen Wissenschaften*. Springer-Verlag Berlin Heidelberg, 1st edition, 1999.
- [32] T. Komorowski, C. Landim, and S. Olla. *Fluctuations in Markov Processes*, volume 345 of *Grundlehren der mathematischen Wissenschaften*. Springer-Verlag Berlin Heidelberg, 1st edition, 2012.
- [33] C. Landim, A. Milanés, and S. Olla. Stationary and non-equilibrium fluctuations in boundary driven exclusion processes. *Markov Proc. Rel. Fields*, 14:165–184, 2008.
- [34] T. J. Lyons, M. Caruana, and T. Lévy. *Differential equations driven by rough paths*, volume 1908 of *Lecture Notes in Mathematics*. Springer, Berlin, 2007.
- [35] D. Nualart. *The Malliavin calculus and related topics*. Probability and its Applications (New York). Springer-Verlag, Berlin, second edition, 2006.
- [36] V. G. Papanicolaou. The probabilistic solution of the third boundary value problem for second order elliptic equations. *Probab. Theory Related Fields*, 87(1):27–77, 1990.
- [37] Shalin Parekh. The kpz limit of ase with boundary. *Communications in Mathematical Physics*, Sep 2018.
- [38] J. Quastel. Introduction to KPZ. In *Current developments in mathematics, 2011*, pages 125–194. Int. Press, Somerville, MA, 2012.
- [39] J. Quastel and H. Spohn. The one-dimensional KPZ equation and its universality class. *Journal of Statistical Physics*, 160:965–984, August 2015.
- [40] Herbert Spohn. The Kardar-Parisi-Zhang equation: a statistical physics perspective. In *Stochastic processes and random matrices*, pages 177–227. Oxford Univ. Press, Oxford, 2017.
- [41] J. B. Walsh. An introduction to stochastic partial differential equations. In *École d’été de probabilités de Saint-Flour, XIV—1984*, volume 1180 of *Lecture Notes in Math.*, pages 265–439. Springer, Berlin, 1986.
- [42] L. C. Young. An inequality of the Hölder type, connected with Stieltjes integration. *Acta Math.*, 67(1):251–282, 1936.

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